



Higgs-boson-pair production $H(\rightarrow bb)H(\rightarrow \gamma\gamma)$ from gluon fusion at the HL-LHC and HL-100 TeV hadron collider

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- Ref : 1. An exploratory study of Higgs-boson pair production (JHEP 1508(2015) 133)
2. Higgs-boson-pair production $H(\rightarrow bb)H(\rightarrow \gamma\gamma)$ from gluon fusion
at the HL-LHC and HL-100 TeV hadron collider (arxiv :1804.07130)
3. [Higgs-boson-pair production \$H\(\rightarrow bb\)H\(\rightarrow \gamma\gamma\)\$ from gluon fusion with multivariate technique](#)
(Work in Progress)

24 August @ ACADEMIA SINICA, TAIWAN
Jubin Park (Chonnam National University)

Contents I



- Motivations
- Understanding the process in the effective Lagrangian
- Outline of simulations and event selections
- Essence of analysis results at the HL-LHC and HL-100 TeV hadron collider
- Conclusion I

Contents II



- Why Higgs pair production so difficult ?
- Why Higgs pair production so interesting ?
- Machine Learning approaches to the Higgs boson self coupling
- Summary Table
- Conclusion II

Contents in the near future



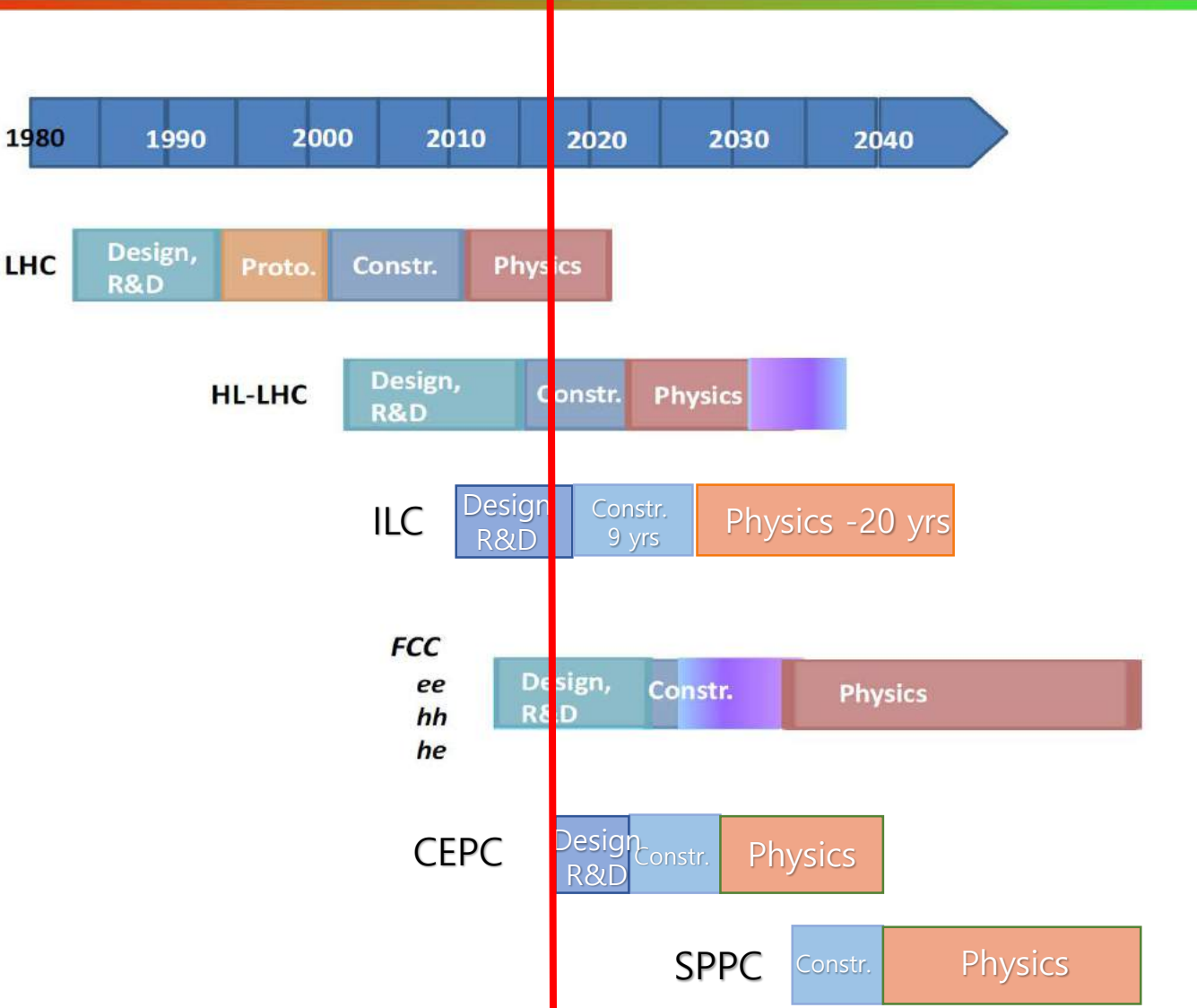
- Recent updates
 1. NLO and NNLO QCD corrections....
 2. TMVA is used. The BTD method can give a better result.
 2. Improving the efficiencies (or removing background more) using the deep neural network....

Motivations

- Self-coupling of the Higgs boson is a crucial property which depends on the dynamics of the electroweak symmetry breaking sector.
- One of the probes of Higgs self-coupling is the Higgs-boson-pair production at the LHC.
- In this work, we perform the most up-to-date comprehensive signal-background analysis for Higgs-pair production through gluon fusion and the $H(\rightarrow bb)H(\rightarrow \gamma\gamma)$ channel at the HL-LHC and HL-100 TeV hadron collider, with the goal of probing the self-coupling λ_{3H} of the Higgs boson.

Timeline

We are here



Understanding the process

in the effective Lagrangian

$$-\mathcal{L} = \frac{1}{3!} \left(\frac{3M_H^2}{v} \right) \lambda_{3H} H^3 + \frac{m_t}{v} \bar{t} (g_t^S + i\gamma_5 g_t^P) t H + \frac{1}{2} \frac{m_t}{v^2} \bar{t} (g_{tt}^S + i\gamma_5 g_{tt}^P) t H^2$$

SM Higgs self couplings

★ In the SM, $\lambda_{3H} = g_t^S = 1$ and $g_{tt}^S = g_{tt}^P = 0$

$$\mathcal{L} = -\frac{1}{2} m_H^2 H^2 - \frac{g_{HHH}}{3!} H^3 - \frac{g_{HHHH}}{4!} H^4$$

In the gluon fusion process at the hadron collide

$$g(p_1)g(p_2) \rightarrow H(p_3)H(p_4)$$

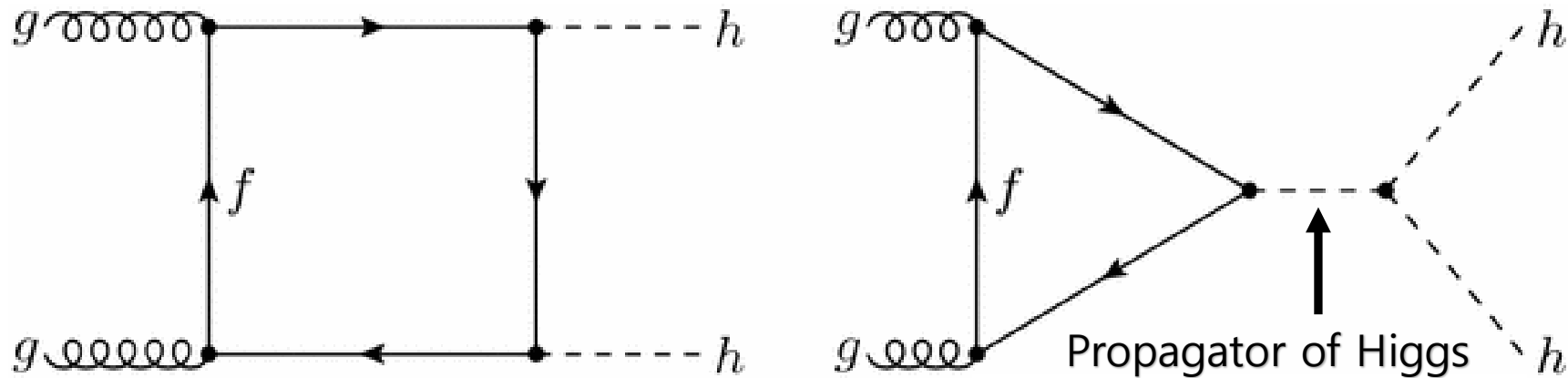
$$g_{HHH} = \frac{3m_H^2}{v}, \quad g_{HHHH} = \frac{3m_H^2}{v^2}$$

The differential cross section is given by

$$\frac{d\hat{\sigma}(gg \rightarrow HH)}{d\hat{t}} = \frac{G_F^2 \alpha_s^2}{512(2\pi)^3} \left[\left| \lambda_{3H} g_t^S D(\hat{s}) F_{\Delta}^S + (g_t^S)^2 F_{\square}^{SS} \right|^2 + \left| (g_t^S)^2 G_{\square}^{SS} \right|^2 \right]$$

Feynman diagrams

Only QCD Leading Order (LO)



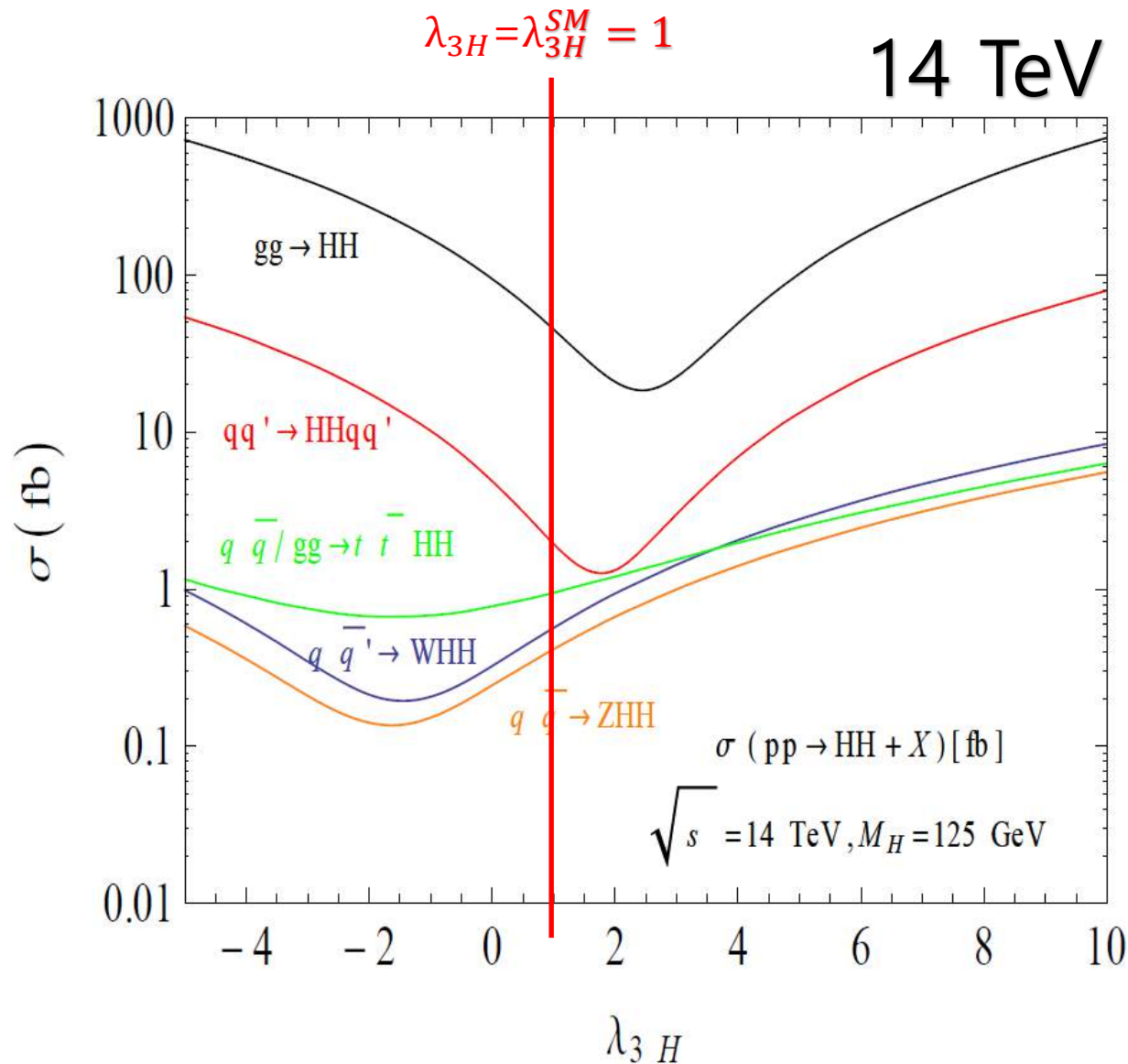
$$D(\hat{s}) = \frac{3M_H^2}{\hat{s} - M_H^2 + iM_H\Gamma_H}$$

$$\frac{d\hat{\sigma}(gg \rightarrow HH)}{d\hat{t}} = \frac{G_F^2 \alpha_s^2}{512(2\pi)^3} \left[\left| \lambda_{3H} g_t^S D(\hat{s}) F_{\Delta}^S + (g_t^S)^2 F_{\square}^{SS} \right|^2 + \left| (g_t^S)^2 G_{\square}^{SS} \right|^2 \right]$$



Important Interference term !!! $\leftrightarrow \lambda_{3H}^{Non-SM}$

For the reference, there are various production modes

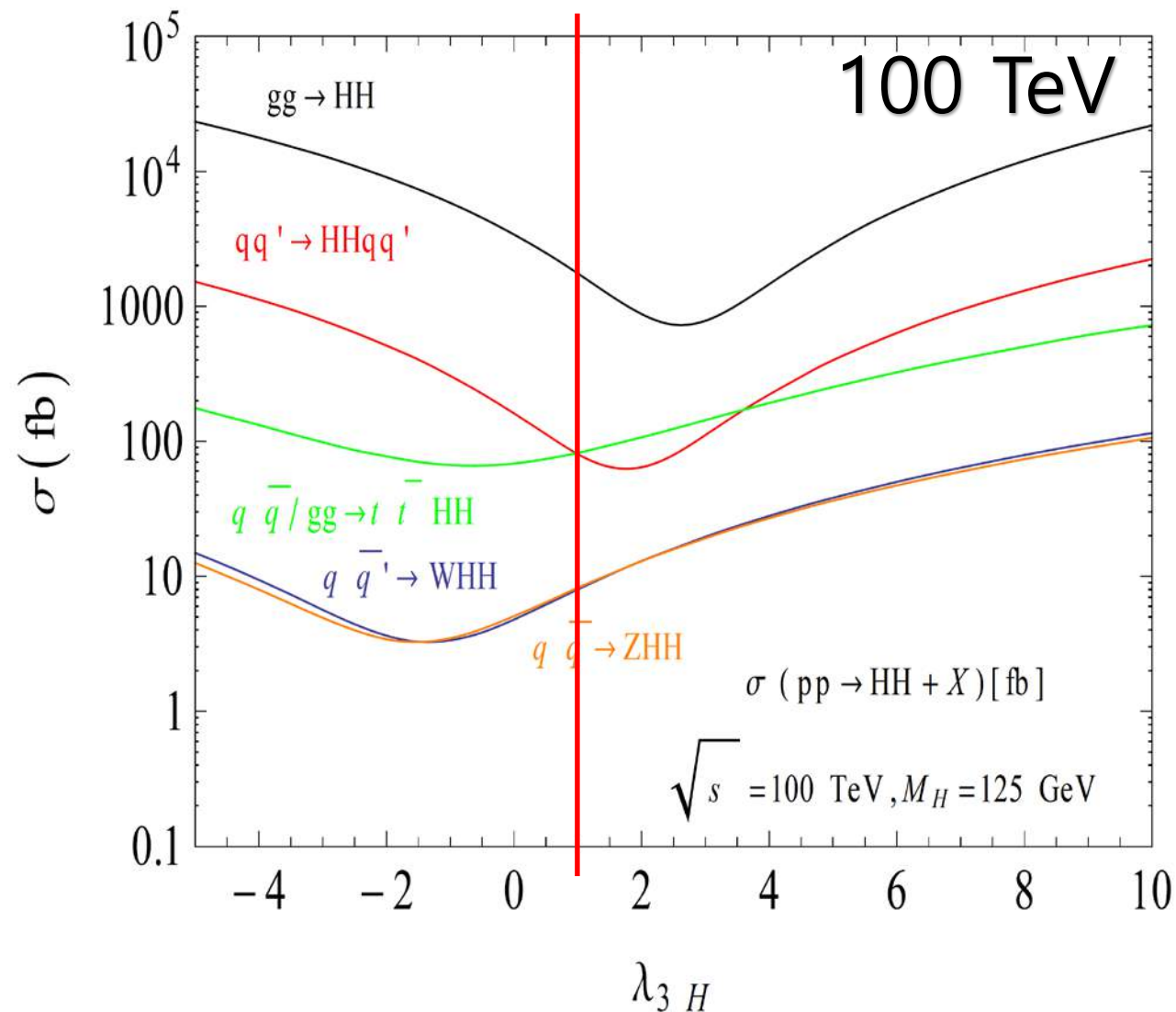


The gluon fusion production mode is dominant one !

$(gg \rightarrow HH) = 45.05 \text{ fb},$
 $(qq_0 \rightarrow HHqq_0) = 1.94 \text{ fb},$
 $(qq_0 \rightarrow V HH) = 0.567 (V = W) = 0.415 (V = Z) \text{ fb},$
 $(gg/qq \rightarrow ttHH) = 0.949 \text{ fb}$
 are calculated at
 NNLO+NNLL, NLO, NNLO, and NLO, respectively

continue

$$\lambda_{3H} = \lambda_{3H}^{SM} = 1$$



The 100 TeV cross sections

$(gg \rightarrow HH) = 1749 \text{ fb},$

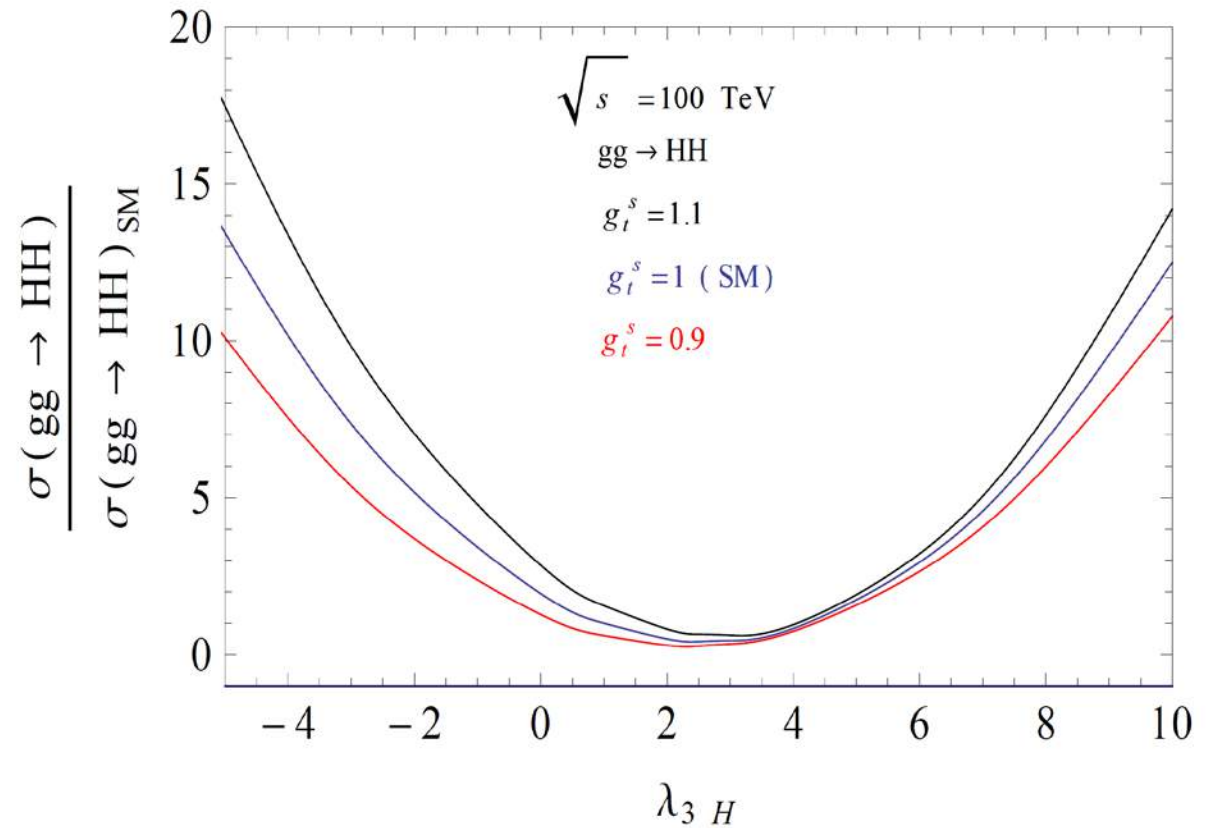
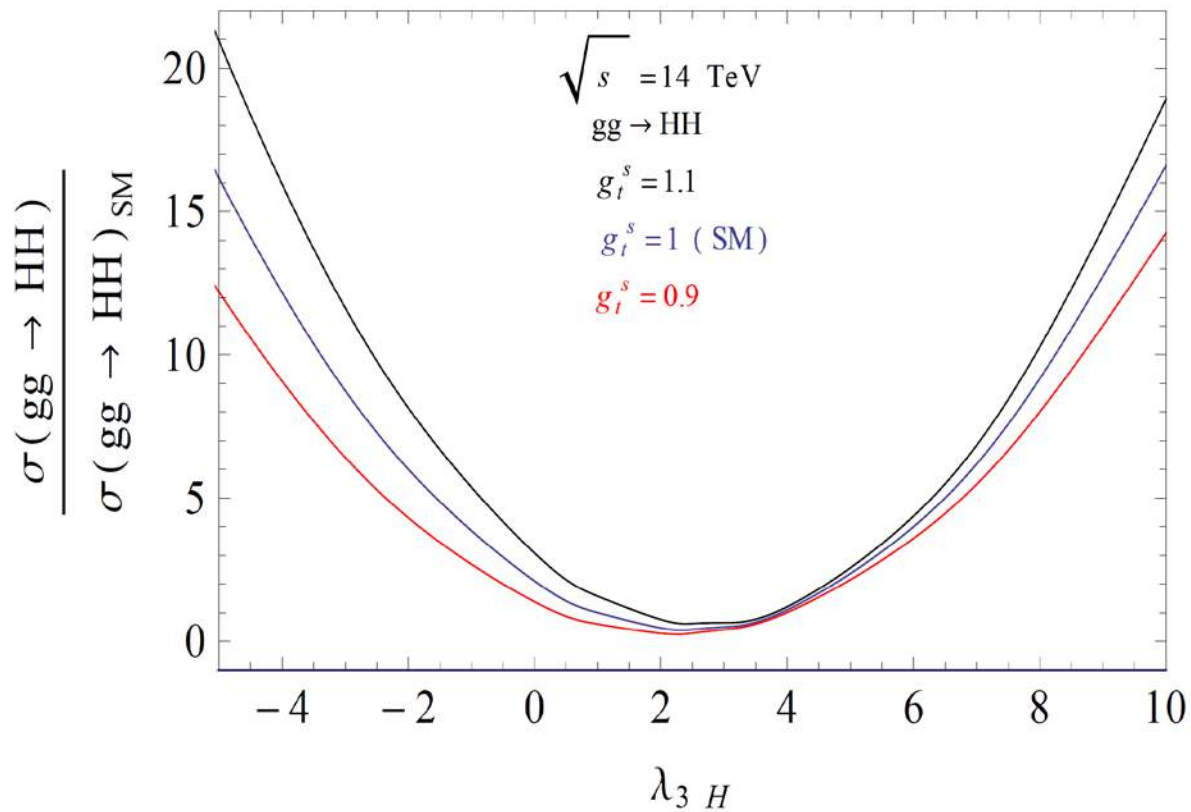
$(qq_0 \rightarrow HHqq_0) = 80.3 \text{ fb},$

$(qq_0 \rightarrow V HH) = 8.00 (V=W) = 8.23 (V=Z) \text{ fb},$

$(gg/qq \rightarrow ttHH) = 82.1 \text{ fb}$

are calculated at the same orders as at 14 TeV.

Including top Yukawa uncertainty !



Ratio of cross sections $(gg \rightarrow HH) = (gg \rightarrow HH)_{\text{SM}}$ versus λ_{3H} taking account of 10% uncertainty of the top-Yukawa coupling: $g_{t^S} = 1:1$ (black), 1 (blue), and 0:9 (red) for $\sqrt{s} = 14 \text{ TeV}$ (left) and $\sqrt{s} = 100 \text{ TeV}$ (right).

$$\frac{d\hat{\sigma}(gg \rightarrow HH)}{d\hat{t}} = \frac{G_F^2 \alpha_s^2}{512(2\pi)^3} \left[\left| \lambda_{3H} g_t^S D(\hat{s}) F_{\Delta}^S + (g_t^S)^2 F_{\square}^{SS} \right|^2 + \left| (g_t^S)^2 G_{\square}^{SS} \right|^2 \right]$$

In the heavy quark limit

$$F_{\Delta}^S = +\frac{2}{3} + \mathcal{O}(\hat{s}/m_Q^2), \quad F_{\square}^{SS} = -\frac{2}{3} + \mathcal{O}(\hat{s}/m_Q^2), \quad F_{\square}^{PP} = +\frac{2}{3} + \mathcal{O}(\hat{s}/m_Q^2)$$

There is large cancellation between the triangle and box diagrams

The production cross section normalized to the corresponding SM cross section :

$$\frac{\sigma^{\text{LO}}(gg \rightarrow HH)}{\sigma_{\text{SM}}^{\text{LO}}(gg \rightarrow HH)} = \underbrace{c_1(s)}_{0.263} \lambda_{3H}^2 (g_t^S)^2 + \overset{\star \text{Interference term}}{\underbrace{c_2(s)}_{-1.310}} \lambda_{3H} (g_t^S)^3 + \underbrace{c_3(s)}_{2.047} (g_t^S)^4$$

14 TeV

100 TeV

Outline of simulations and event selections

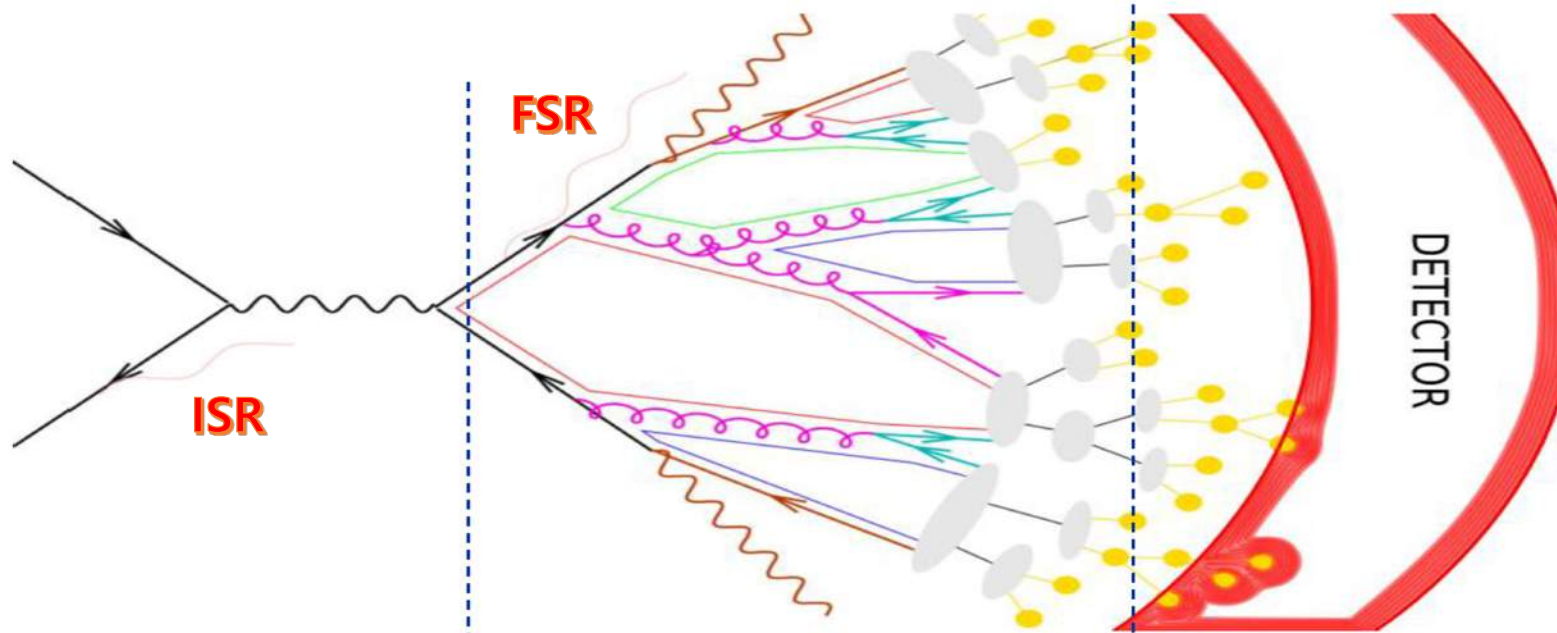
- Our goal is to disentangle the effects of the trilinear Higgs coupling, which is present in the triangle diagram, in Higgs pair production.
- We vary the value of the trilinear coupling λ_{3H} between -5 and 10 to visualize the effects of λ_{3H} .
- We consider the full set of backgrounds.

Outline of simulations

Hard Scattering Part

Hadronization and showering
Processes

Reconstruction level



Parton level

Hadron level

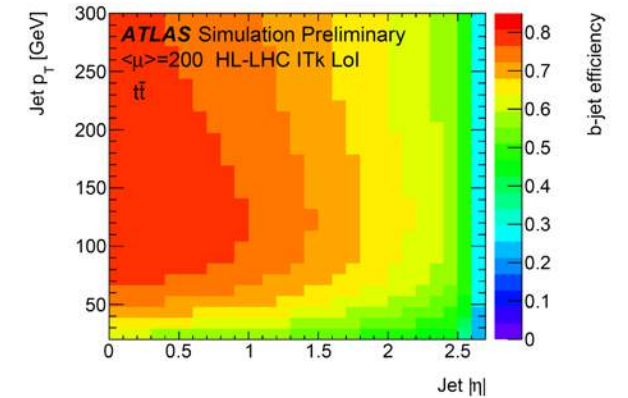
Reconstructed
objects level

LHE files

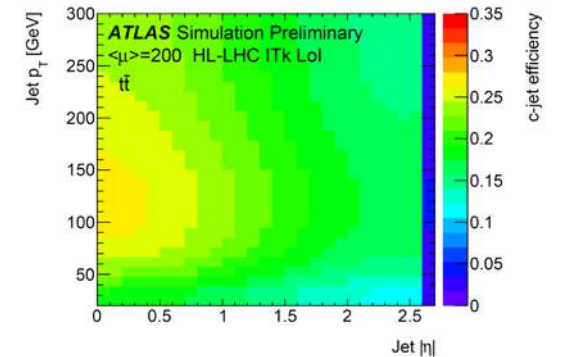
STDHEP/HEPMC files

LHCO, ROOT files

b-jet efficiency



c-jet efficiency



ATL-PHYS-PUB-2016-026

Signal

HL-LHC 14 TeV

Signal process	Generator/Parton Shower	$\sigma \cdot BR$ [fb]	Order	PDF used
$gg \rightarrow HH \rightarrow b\bar{b}\gamma\gamma$ [15]	MG5_aMC@NLO/PYTHIA8	0.119	NNLO in QCD +NNLL	NNPDF2.3LO

Backgrounds

Backgrounds					
Background(BG)	Process	Generator/Parton Shower	$\sigma \cdot BR$ [fb]	Order	PDF used
				in QCD	
Single-Higgs associated BG [15]	$ggH(\rightarrow \gamma\gamma)$	POWHEG – BOX/PYTHIA6	1.20×10^2	NNLO	CT10
	$t\bar{t}H(\rightarrow \gamma\gamma)$	PYTHIA8/PYTHIA8	1.37	NLO	
	$ZH(\rightarrow \gamma\gamma)$	PYTHIA8/PYTHIA8	2.24	NLO	
	$b\bar{b}H(\rightarrow \gamma\gamma)$	PYTHIA8/PYTHIA8	1.26	NLO	
Non-resonant BG	$b\bar{b}\gamma\gamma$	MG5_aMC@NLO/PYTHIA8	1.40×10^2	LO	CTEQ6L1
	$c\bar{c}\gamma\gamma$	MG5_aMC@NLO/PYTHIA8	1.14×10^3	LO	
	$jj\gamma\gamma$	MG5_aMC@NLO/PYTHIA8	1.62×10^4	LO	
	$b\bar{b}j\gamma$	MG5_aMC@NLO/PYTHIA8	3.67×10^5	LO	
	$c\bar{c}j\gamma$	MG5_aMC@NLO/PYTHIA8	1.05×10^6	LO	
	$b\bar{b}jj$	MG5_aMC@NLO/PYTHIA8	4.34×10^8	LO	
	$Z(\rightarrow b\bar{b})\gamma\gamma$	MG5_aMC@NLO/PYTHIA8	5.17	LO	
$t\bar{t}$ and $t\bar{t}\gamma$ BG	$t\bar{t}$ [18]	POWHEG – BOX/PYTHIA8	5.30×10^5	NNLO +NNLL	CT10
	(≥ 1 lepton) $t\bar{t}\gamma$ [19]	MG5_aMC@NLO/PYTHIA8	1.60×10^3	NLO	CTEQ6L1

Main fake processes and the corresponding rates

HL-LHC 14 TeV

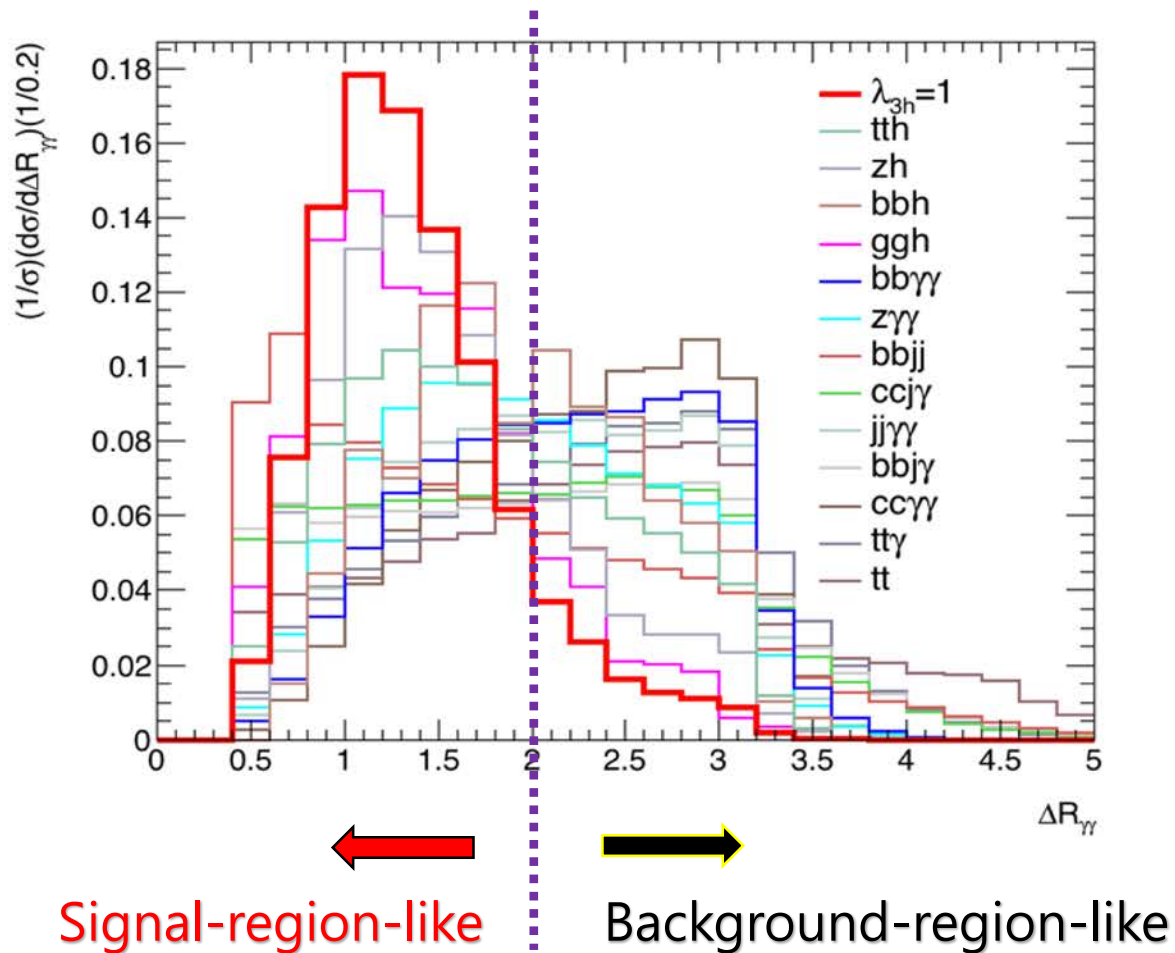
Background(BG)	Process	Fake Process	Fake rate
Non-resonant BG	$b\bar{b}\gamma\gamma$	N/A	N/A
	$c\bar{c}\gamma\gamma$	$c \rightarrow b, \bar{c} \rightarrow \bar{b}$	$(P_{c \rightarrow b})^2$
	$j\bar{j}\gamma\gamma$	$c_s \rightarrow b, \bar{c}_s \rightarrow \bar{b}$	$(P_{c_s \rightarrow b})^2$
	$b\bar{b}j\gamma$	$j \rightarrow \gamma$	5×10^{-4}
	$c\bar{c}j\gamma$	$c \rightarrow b, \bar{c} \rightarrow \bar{b}, j \rightarrow \gamma$	$(P_{c \rightarrow b})^2 \cdot (5 \times 10^{-4})$
	$b\bar{b}jj$	$j \rightarrow \gamma, j \rightarrow \gamma$	$(5 \times 10^{-4})^2$
	$Z(\rightarrow b\bar{b})\gamma\gamma$	N/A	N/A
$t\bar{t}$	Leptonic decay	$e \rightarrow \gamma, e \rightarrow \gamma$	$(0.02)^2/0.02 \cdot 0.05/(0.05)^2$
	Semi-leptonic decay	$e \rightarrow \gamma, j \rightarrow \gamma$	$(0.02) \cdot 5 \times 10^{-4}/(0.05) \cdot 5 \times 10^{-4}$
$t\bar{t}\gamma$	Leptonic decay	$e \rightarrow \gamma$	0.02/0.05
	Semi-leptonic	$e \rightarrow \gamma$	0.02/0.05

Event selection

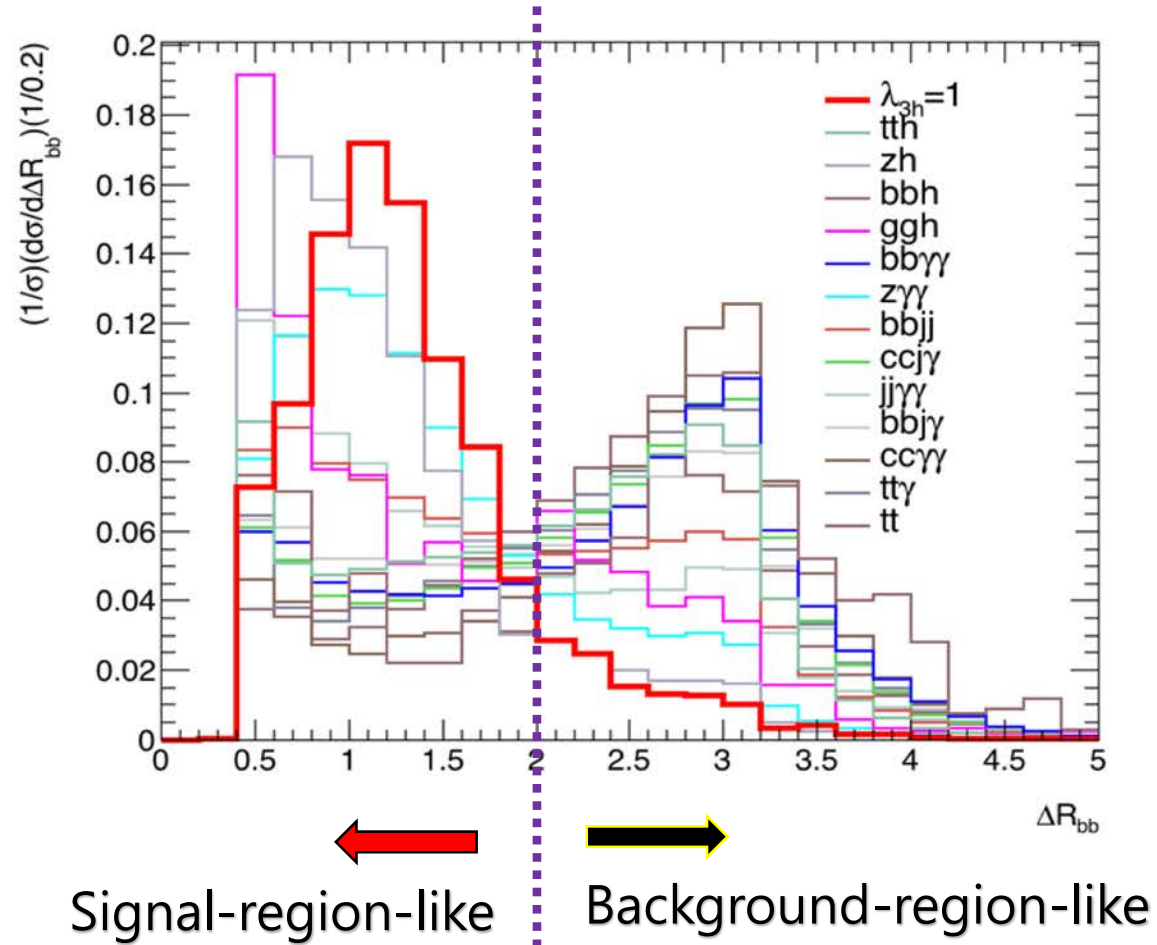
Sequence	Event Selection Criteria at the HL-LHC
1	Di-photon trigger condition, ≥ 2 isolated photons with $P_T > 25$ GeV, $ \eta < 2.5$
2	≥ 2 isolated photons with $P_T > 30$ GeV, $ \eta < 1.37$ or $1.52 < \eta < 2.37$, $\Delta R_{j\gamma} > 0.4$
3	≥ 2 jets identified as b-jets with leading(subleading) $P_T > 40(30)$ GeV, $ \eta < 2.4$
4	Events are required to contain ≤ 5 jets with $P_T > 30$ GeV within $ \eta < 2.5$
5	No isolated leptons with $P_T > 25$ GeV, $ \eta < 2.5$
6	$0.4 < \Delta R_{b\bar{b}} < 2.0$, $0.4 < \Delta R_{\gamma\gamma} < 2.0$
7	$122 < M_{\gamma\gamma}/\text{GeV} < 128$ and $100 < M_{b\bar{b}}/\text{GeV} < 150$
8	$P_T^{\gamma\gamma} > 80$ GeV, $P_T^{b\bar{b}} > 80$ GeV

These conditions of cuts are very important to distinguish signal and background !!!!!

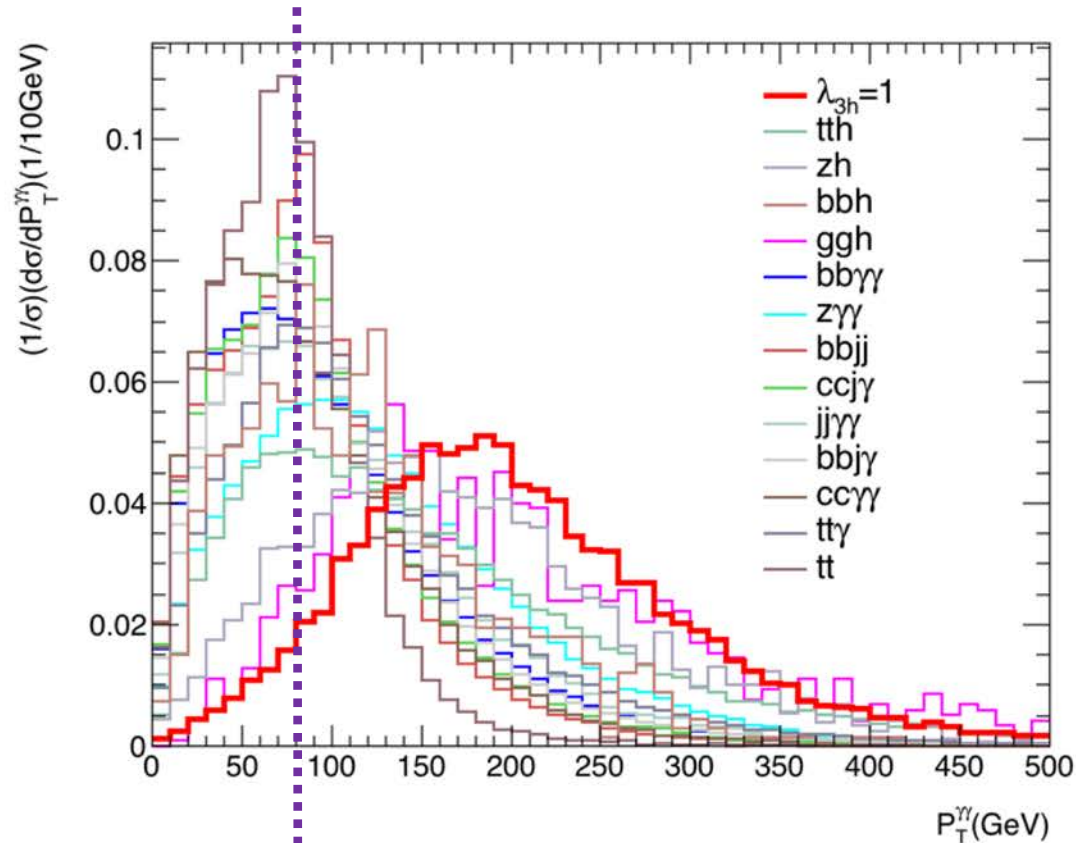
$$0.4 < \Delta R_{\gamma\gamma} < 2.0$$



$$0.4 < \Delta R_{b\bar{b}} < 2.0$$



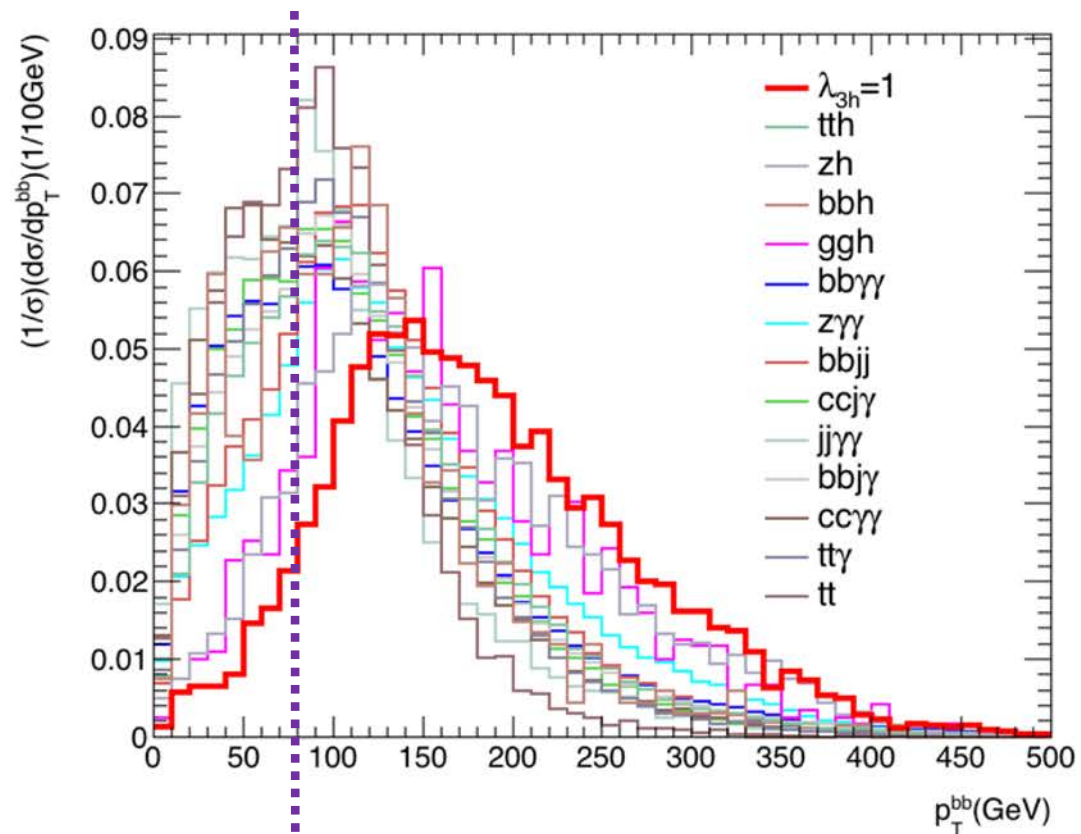
$$P_T^{\gamma\gamma} > 80 \text{ GeV}$$



Background
-region-like

Signal-region-like

$$P_T^{b\bar{b}} > 80 \text{ GeV}$$



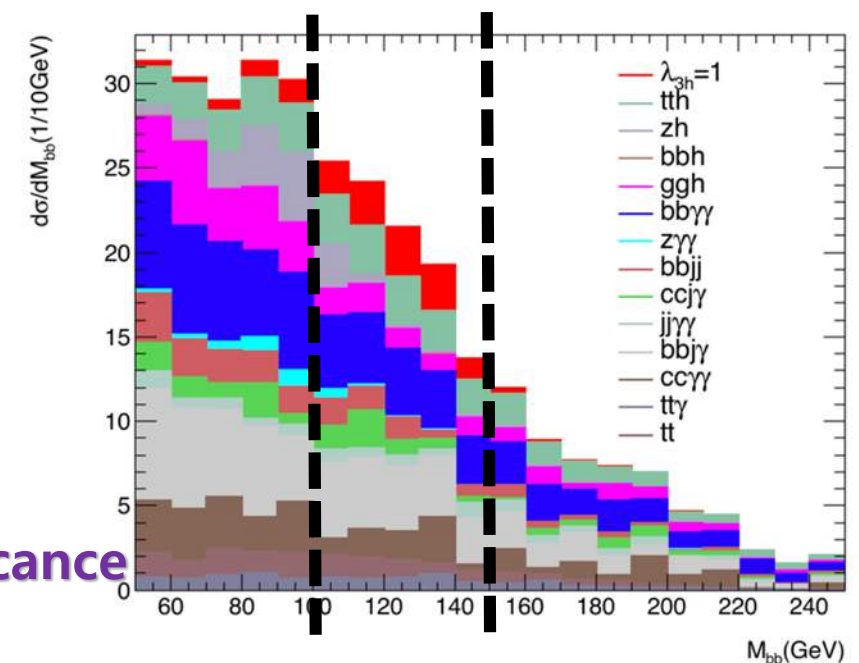
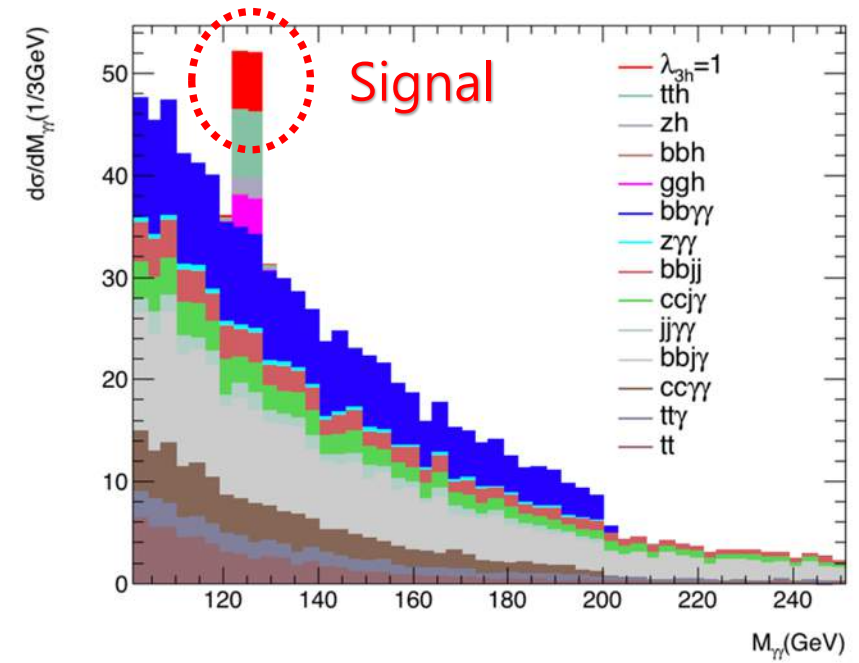
Background
-region-like

Signal-region-like

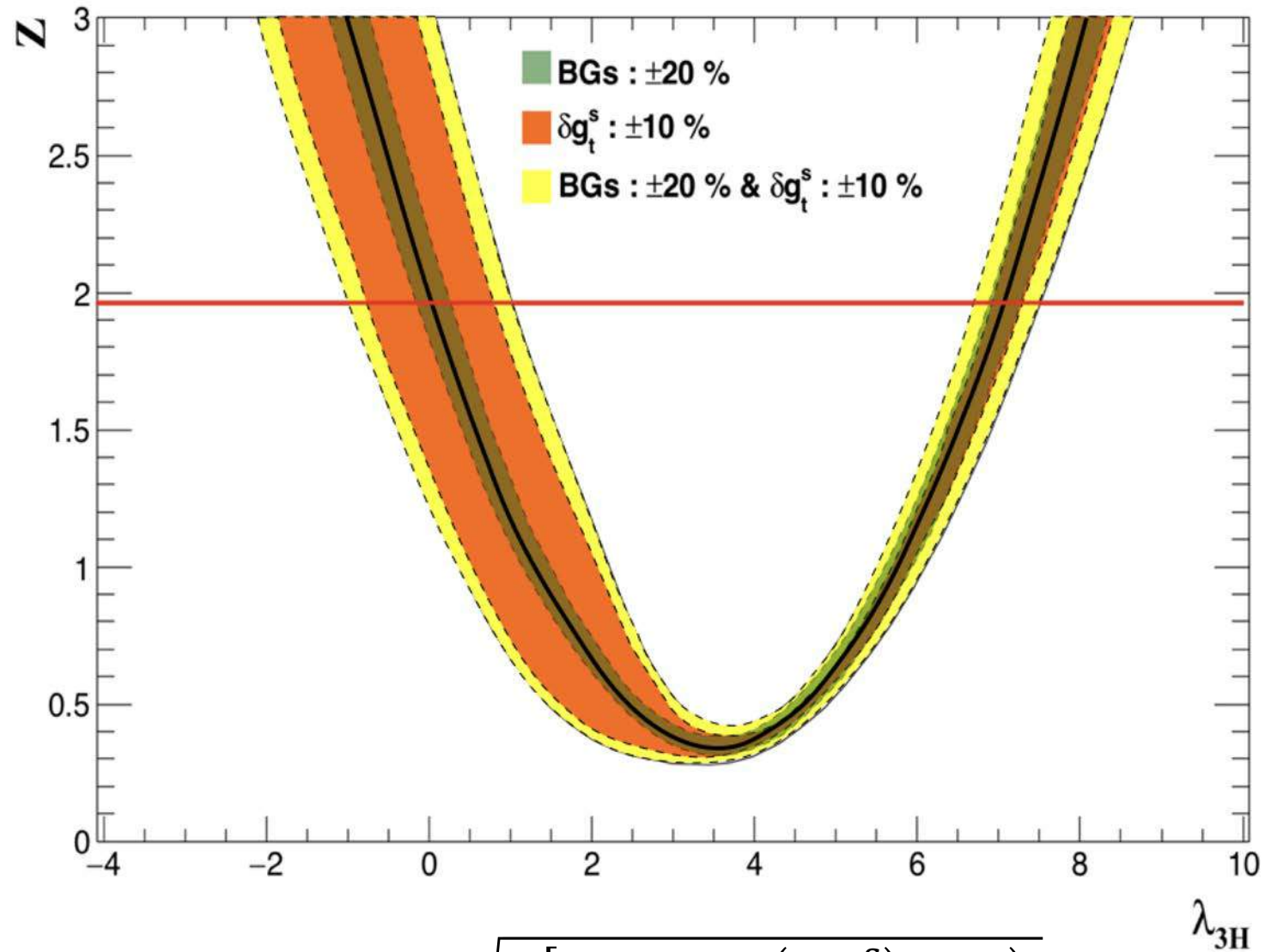
Final Results I :
Expected Yields and Kinematic distributions
of $M_{\gamma\gamma}$ (GeV) and $M_{b\bar{b}}$ (GeV)

Expected yields (3000 fb ⁻¹)	Total	Barrel-barrel	Other (End-cap)	Ratio (O/B)
<i>Samples</i>				
$H(b\bar{b})H(\gamma\gamma), \lambda_{3H} = -4$	77.14	57.03	20.11	0.35
$H(b\bar{b})H(\gamma\gamma), \lambda_{3H} = 0$	19.50	14.33	5.17	0.36
$H(b\bar{b})H(\gamma\gamma), \lambda_{3H} = 1$	11.42	8.53	2.89	0.34
$H(b\bar{b})H(\gamma\gamma), \lambda_{3H} = 2$	6.82	5.14	1.68	0.33
$H(b\bar{b})H(\gamma\gamma), \lambda_{3H} = 6$	11.03	7.91	3.12	0.39
$H(b\bar{b})H(\gamma\gamma), \lambda_{3H} = 10$	57.46	41.94	15.52	0.37
$ggH(\gamma\gamma)$	6.60	4.50	2.10	0.47
$t\bar{t}H(\gamma\gamma)$	13.21	9.82	3.39	0.35
$ZH(\gamma\gamma)$	3.62	2.44	1.18	0.48
$b\bar{b}H(\gamma\gamma)$	0.15	0.11	0.04	0.40
$b\bar{b}\gamma\gamma$	18.86	11.15	7.71	0.69
$c\bar{c}\gamma\gamma$	7.53	4.79	2.74	0.57
$jj\gamma\gamma$	3.34	1.59	1.75	1.10
$b\bar{b}j\gamma$	18.77	10.40	8.37	0.80
$c\bar{c}j\gamma$	5.52	3.94	1.58	0.40
$b\bar{b}jj$	5.54	3.81	1.73	0.45
$Z(b\bar{b})\gamma\gamma$	0.90	0.54	0.36	0.67
$t\bar{t} (\geq 1 \text{ leptons})$	4.98	3.04	1.94	0.64
$t\bar{t}\gamma (\geq 1 \text{ leptons})$	3.61	2.29	1.32	0.58
Total Background	92.63	58.42	34.21	0.59
Significance Z	1.163	1.090	0.487	
Combined significance		1.194		

Combined significance = 1.194

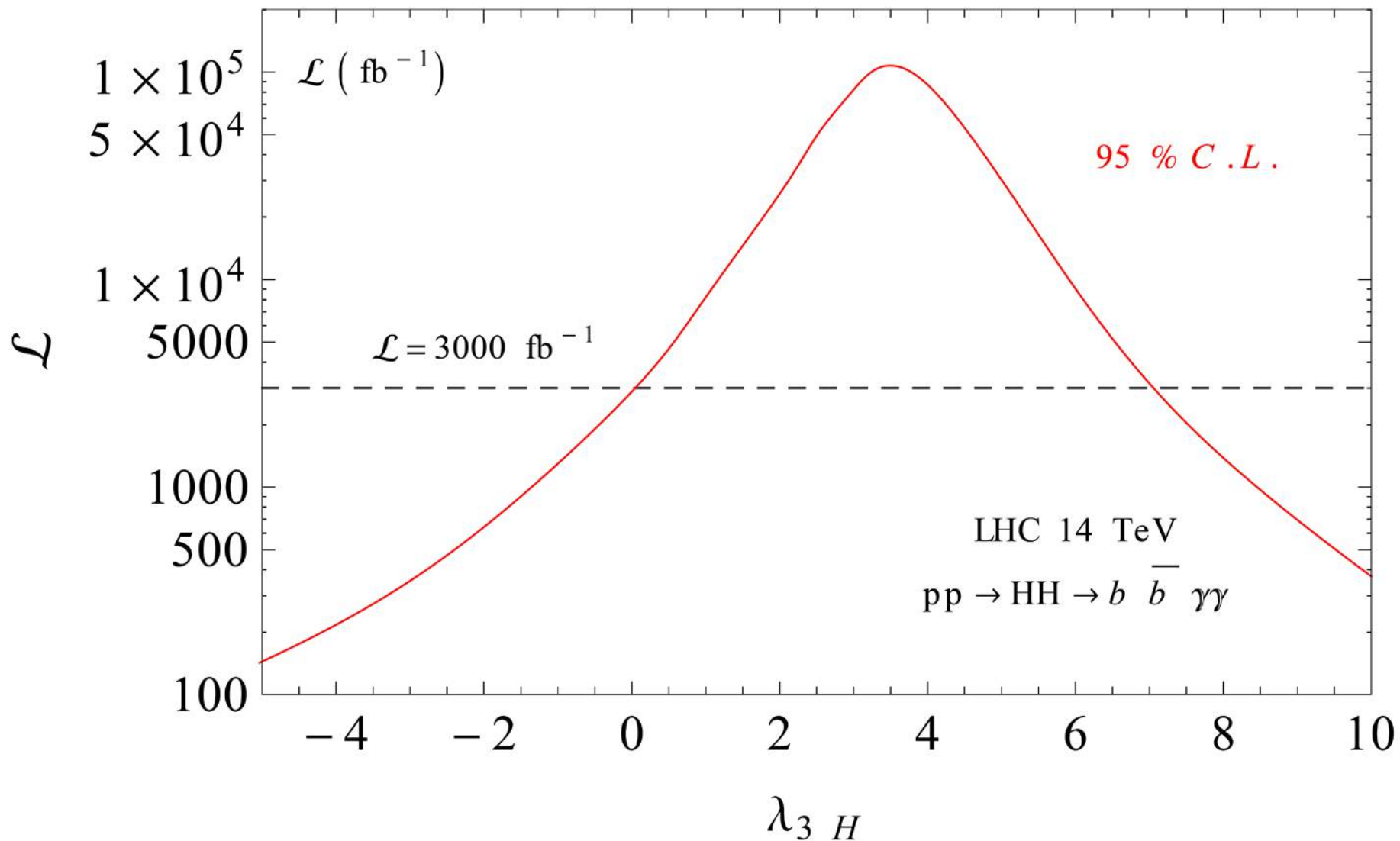


Essence of analysis results at the HL-LHC



$$Z = \sqrt{2 \left[(s + b) \ln \left(1 + \frac{s}{b} \right) - s \right]}$$

Required luminosity for 95% confidence level sensitivity at the 14 TeV HL-LHC v.s λ_{3H} .



Signal

HL-100 TeV hadron collider

TABLE V. Same as Table I but for a 100 TeV hadron collider. In the row for $t\bar{t}H(\rightarrow \gamma\gamma)$, 5FS stands for the 5-flavor scheme.

Signal				
Signal process	Generator/Parton Shower	$\sigma \cdot BR$ [fb]	Order in QCD	PDF used
$gg \rightarrow HH \rightarrow b\bar{b}\gamma\gamma$ [16]	MG5_aMC@NLO/PYTHIA8	4.62	NNLO +NNLL	NNPDF2.3LO

Backgrounds

Backgrounds					
Background(BG)	Process	Generator/Parton Shower	$\sigma \cdot BR$ [fb]	Order in QCD	PDF used
Single-Higgs associated BG	$ggH(\rightarrow \gamma\gamma)$ [16]	POWHEG – BOX/PYTHIA8	1.82×10^3	NNNLO	CT10
	$t\bar{t}H(\rightarrow \gamma\gamma)$ [16]	PYTHIA8/PYTHIA8	7.29×10^1	NLO	
	$ZH(\rightarrow \gamma\gamma)$ [16]	PYTHIA8/PYTHIA8	2.54×10^1	NNLO	
	$b\bar{b}H(\rightarrow \gamma\gamma)$ [30]	PYTHIA8/PYTHIA8	1.96×10^1	NNLO(5FS)	
Non-resonant BG	$b\bar{b}\gamma\gamma$	MG5_aMC@NLO/PYTHIA8	4.93×10^3	LO	CTEQ6L1
	$c\bar{c}\gamma\gamma$	MG5_aMC@NLO/PYTHIA8	4.54×10^4	LO	
	$jj\gamma\gamma$	MG5_aMC@NLO/PYTHIA8	5.38×10^5	LO	
	$b\bar{b}j\gamma$	MG5_aMC@NLO/PYTHIA8	1.44×10^7	LO	
	$c\bar{c}j\gamma$	MG5_aMC@NLO/PYTHIA8	4.20×10^7	LO	
	$b\bar{b}jj$	MG5_aMC@NLO/PYTHIA8	1.60×10^{10}	LO	
	$Z(\rightarrow b\bar{b})\gamma\gamma$	MG5_aMC@NLO/PYTHIA8	9.53×10^1	LO	
$t\bar{t}$ and $t\bar{t}\gamma$ BG [16] (≥ 1 lepton)	$t\bar{t}$	MG5_aMC@NLO/PYTHIA8	1.76×10^7	NLO	CT10
	$t\bar{t}\gamma$	MG5_aMC@NLO/PYTHIA8	4.18×10^4	NLO	CTEQ6L1

Main fake processes and the corresponding rates

HL 100 TeV hadron collider

Background(BG)	Process	Fake Process	Fake rate
Non-resonant BG	$b\bar{b}\gamma\gamma$	N/A	N/A
	$c\bar{c}\gamma\gamma$	$c \rightarrow b, \bar{c} \rightarrow \bar{b}$	$(0.1)^2$
	$j\bar{j}\gamma\gamma$	$c_s \rightarrow b, \bar{c}_s \rightarrow \bar{b}$	$(0.1)^2$
	$b\bar{b}j\gamma$	$j \rightarrow \gamma$	1.35×10^{-3}
	$c\bar{c}j\gamma$	$c \rightarrow b, \bar{c} \rightarrow \bar{b}, j \rightarrow \gamma$	$(0.1)^2 \cdot (1.35 \times 10^{-3})$
	$b\bar{b}jj$	$j \rightarrow \gamma, j \rightarrow \gamma$	$(1.35 \times 10^{-3})^2$
	$Z(\rightarrow b\bar{b})\gamma\gamma$	N/A	N/A
$t\bar{t}$	Leptonic decay	$e \rightarrow \gamma, e \rightarrow \gamma$	$(0.02)^2/0.02 \cdot 0.05/(0.05)^2$
	Semi-leptonic decay	$e \rightarrow \gamma, j \rightarrow \gamma$	$(0.02) \cdot 1.35 \times 10^{-3}/(0.05) \cdot 1.35 \times 10^{-3}$
$t\bar{t}\gamma$	Leptonic decay	$e \rightarrow \gamma$	$0.02/0.05$
	Semi-leptonic	$e \rightarrow \gamma$	$0.02/0.05$

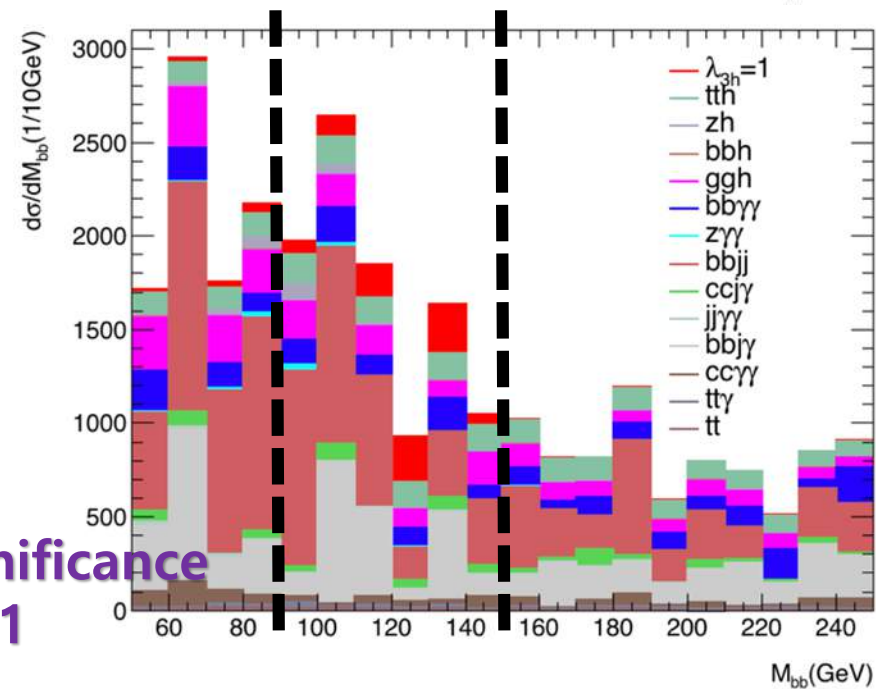
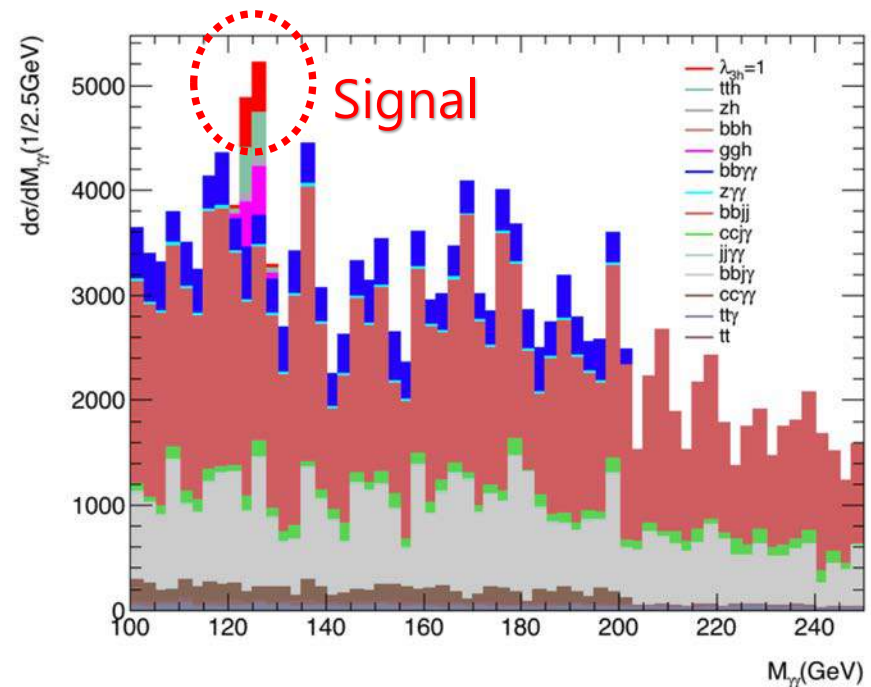
Event selection at the 100 TeV Hadron collider

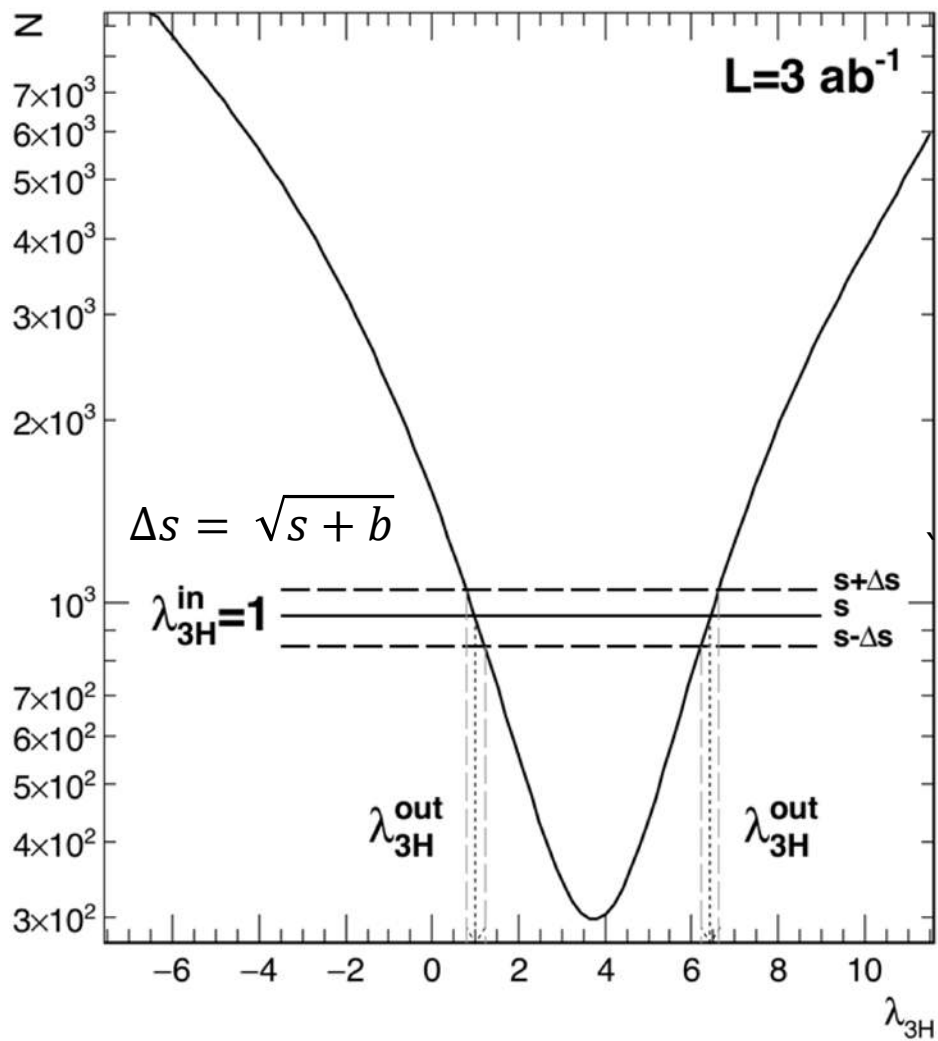
Sequence	Event Selection Criteria at the HL-100 TeV hadron collider
1	Di-photon trigger condition, ≥ 2 isolated photons with $P_T > 30$ GeV, $ \eta < 5$
2	≥ 2 isolated photons with $P_T > 40$ GeV, $ \eta < 3$, $\Delta R_{j\gamma} > 0.4$
3	≥ 2 jets identified as b-jets with leading(subleading) $P_T > 50(40)$ GeV, $ \eta < 3$
4	Events are required to contain ≤ 5 jets with $P_T > 40$ GeV within $ \eta < 5$
5	No isolated leptons with $P_T > 40$ GeV, $ \eta < 3$
6	$0.4 < \Delta R_{b\bar{b}} < 3.0$, $0.4 < \Delta R_{\gamma\gamma} < 3.0$
7	$122.5 < M_{\gamma\gamma}/\text{GeV} < 127.5$ and $90 < M_{b\bar{b}}/\text{GeV} < 150$
8	$P_T^{\gamma\gamma} > 100$ GeV, $P_T^{b\bar{b}} > 100$ GeV

We relaxed these two conditions to enhance the signal yields !

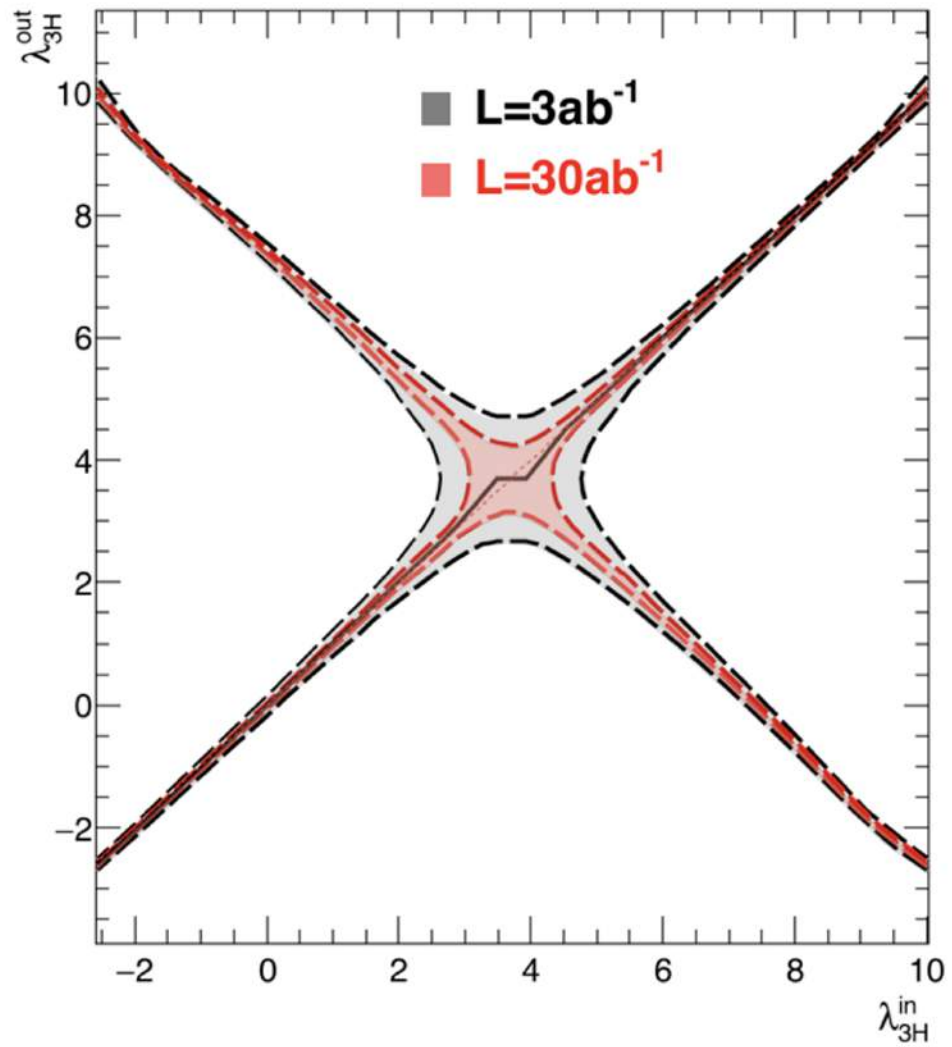
Final Results II :
Expected Yields and Kinematic distributions
of $M_{\gamma\gamma}$ (GeV) and $M_{b\bar{b}}$ (GeV)
at HL-100 TeV collider

Expected yields (3000 fb ⁻¹)	Total	Barrel-barrel	Other (End-cap)	Ratio (O/B)
<i>H</i> (<i>b</i> \bar{b}) <i>H</i> ($\gamma\gamma$), $\lambda_{3H} = -4$	5604.46	4257.36	1347.10	0.32
<i>H</i> (<i>b</i> \bar{b}) <i>H</i> ($\gamma\gamma$), $\lambda_{3H} = 0$	1513.56	1163.04	350.52	0.30
<i>H</i> (<i>b</i> \bar{b}) <i>H</i> ($\gamma\gamma$), $\lambda_{3H} = 1$	941.37	723.86	217.51	0.30
<i>H</i> (<i>b</i> \bar{b}) <i>H</i> ($\gamma\gamma$), $\lambda_{3H} = 2$	557.36	431.45	125.91	0.29
<i>H</i> (<i>b</i> \bar{b}) <i>H</i> ($\gamma\gamma$), $\lambda_{3H} = 6$	753.18	566.18	187.00	0.33
<i>H</i> (<i>b</i> \bar{b}) <i>H</i> ($\gamma\gamma$), $\lambda_{3H} = 10$	3838.33	2924.25	914.08	0.31
<i>ggH</i> ($\gamma\gamma$)	890.47	742.97	147.50	0.20
<i>t</i> \bar{t} <i>H</i> ($\gamma\gamma$)	868.73	659.33	209.40	0.32
<i>ZH</i> ($\gamma\gamma$)	168.86	122.91	45.95	0.37
<i>b</i> \bar{b} <i>H</i> ($\gamma\gamma$)	9.82	7.00	2.82	0.40
<i>b</i> \bar{b} $\gamma\gamma$	783.87	443.70	340.17	0.77
<i>c</i> \bar{c} $\gamma\gamma$	222.88	111.44	111.44	1.00
<i>jj</i> $\gamma\gamma$	32.28	20.98	11.30	0.54
<i>b</i> \bar{b} <i>jj</i>	1982.88	1516.32	466.56	0.31
<i>c</i> \bar{c} <i>jj</i>	293.81	216.49	77.32	0.36
<i>b</i> \bar{b} <i>jj</i>	3674.16	1924.56	1749.60	0.91
<i>Z</i> (<i>b</i> \bar{b}) $\gamma\gamma$	54.87	35.72	19.15	0.54
<i>t</i> \bar{t} (≥ 1 leptons)	59.32	38.32	21.00	0.55
<i>t</i> \bar{t} γ (≥ 1 leptons)	105.68	62.53	43.15	0.69
Total Background	9147.63	5902.27	3245.36	0.55
Significance <i>Z</i>	9.681	9.239	3.777	Combined significance = 9.981
Combined significance	9.981			



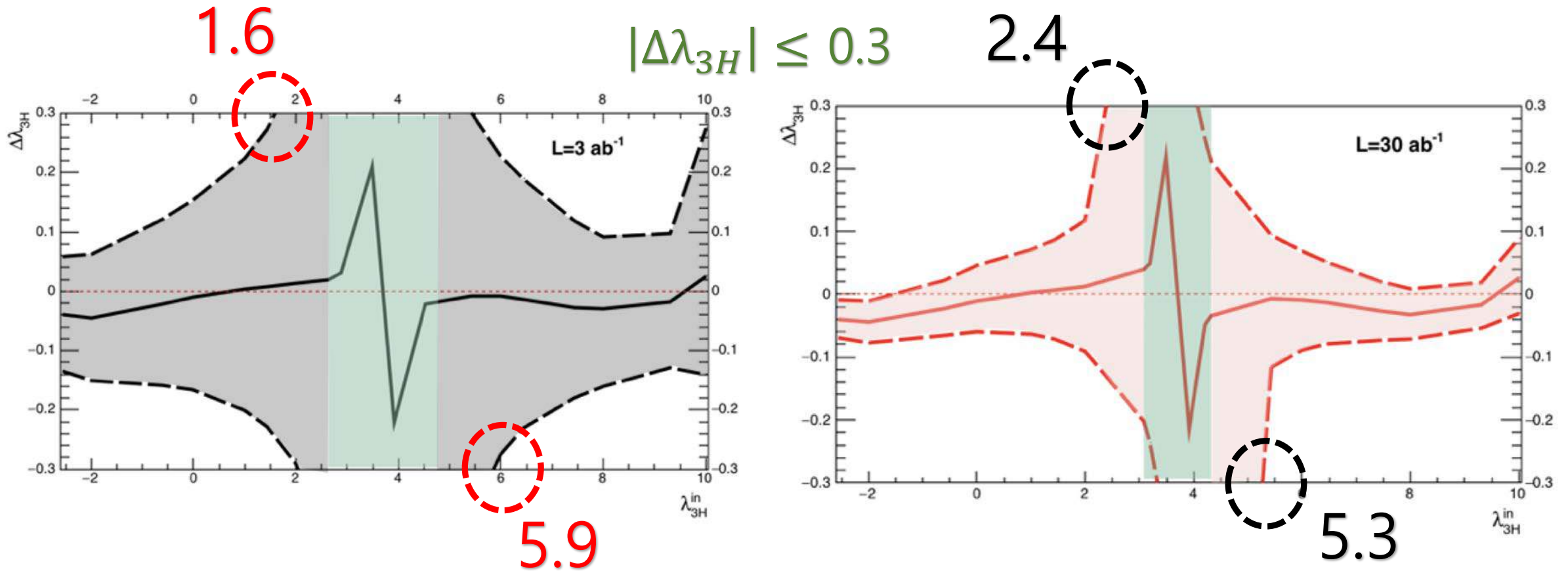


The number of signal event N
v.s λ_{3H} with 3 ab^{-1}



The $1\text{-}\sigma$ error regions v.s the input values of λ_{3H}^{in}
assuming 3 ab^{-1} (black) and 30 ab^{-1} (red)

HL-100 TeV hadron collider



$\Delta\lambda_{3H} = \lambda_{3H}^{\text{out}} - \lambda_{3H}^{\text{in}}$ v.s λ_{3H}^{in} along the $\lambda_{3H}^{\text{out}} = \lambda_{3H}^{\text{in}}$ line
 with 3 ab^{-1} (left) and 30 ab^{-1} (right)

Conclusion I [HL-LHC]

- We find that even for the most promising channel $H(\rightarrow bb)H(\rightarrow \gamma\gamma)$ at the HL-LHC with a luminosity of 3000 fb^{-1} , the significance is still not high enough to establish the Higgs self-coupling at the SM value.
- Instead, we can only constrain the self-coupling to $-1.0 < \lambda_{3H} < 7.6$ at 95% confidence level after considering the uncertainties associated with the top-Yukawa coupling and the estimation of backgrounds.

Conclusion II [HL-100 TeV hadron collider]

- With a luminosity of 3 ab^{-1} , we find there exists a bulk region of $2.6 < \lambda_{3H} < 4.8$ in which one can not pin down the trilinear coupling.
- At the SM value, we show that the coupling can be measured with about 20% accuracy.
- While assuming 30 ab^{-1} , the bulk region reduces to $3.1 < \lambda_{3H} < 4.3$ and the trilinear coupling can be measured with about 7% accuracy at the SM value.

Machine learning approaches to the Higgs boson self coupling

Contents



Why Higgs pair production so difficult ?

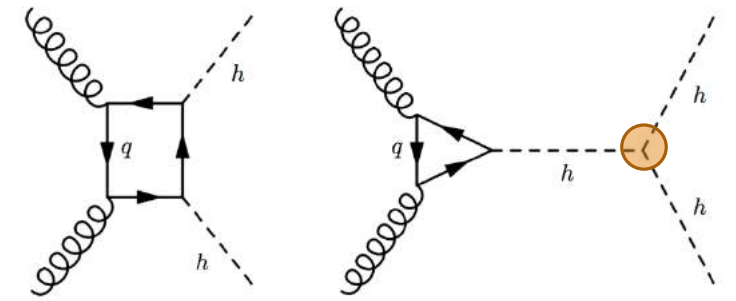
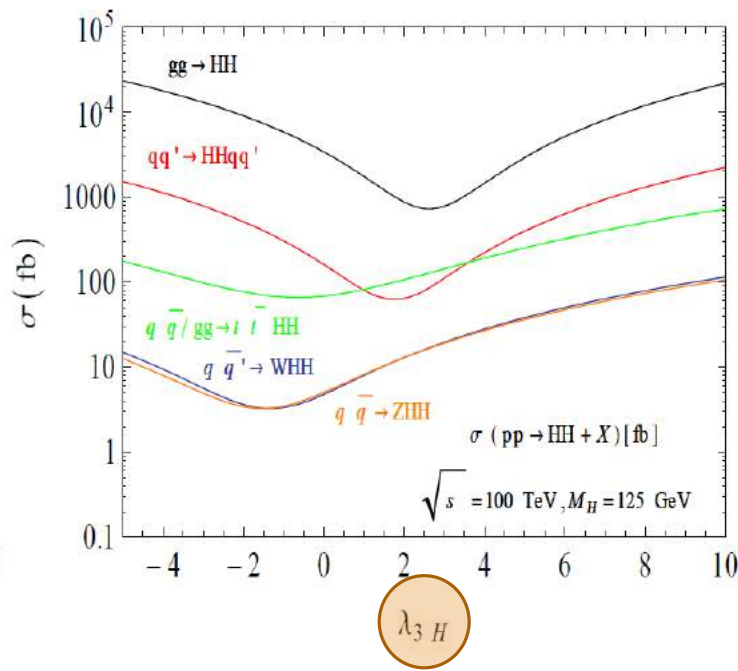
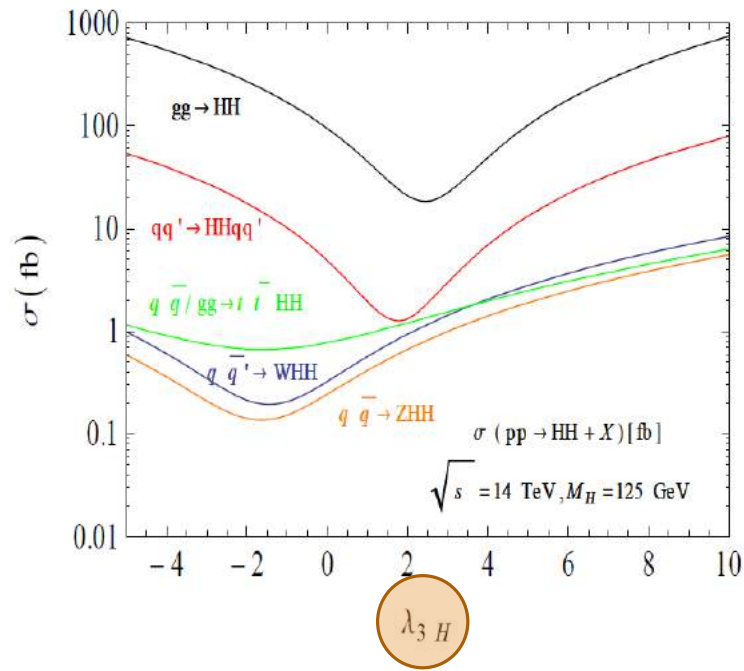
Why Higgs pair production so interesting ?

Machine Learning approaches to the Higgs boson self coupling

Summary Table

Conclusion

Why Higgs pair production so difficult ?



Distructive interference !

In the SM, hh rates are small : In the leading gluon fusion production mode, the cross section at 14 TeV is only 40 fb, further suppressed by each decay branching fractions.



Why Higgs pair production so difficult ?

· $X_{\text{sec}}(\text{gg} \rightarrow \text{hh}) = 39.64 \cdot {}_{-6.0}^{+4.4} (\text{scale}) \pm 2.1 (\text{PDF}) \pm 2.2 (\alpha_s) \text{ fb} @ [14 \text{ TeV}, m_h = 125 \text{ GeV}]$

NNLO cross sections including top quark mass effects to NLO
Phys. Rev. Lett. 117, 012001 [S.Borowka, et al.]

· $O(10^{-3})$ smaller than the single Higgs production (SM)

$X_{\text{sec}}(\text{gg} \rightarrow \text{hh}) \sim 40 \text{ fb}$  $X_{\text{sec}}(\text{gg} \rightarrow \text{h}) \sim 50 \text{ fb}$ @ 14 TeV

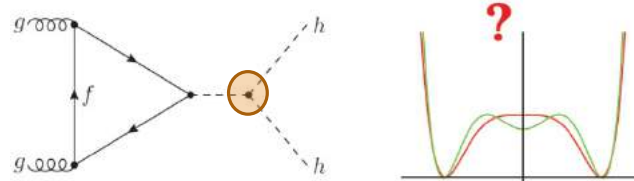
· For the reference, with $X_{\text{sec}} \sim 33 \text{ fb}$ at 13 TeV,
2017 LHC @ 13 TeV with 40 fb^{-1}  1320 Events
14 TeV with 40 fb^{-1}  1600 Events

Why Higgs pair production so interesting ?



Allows accessing crucial components of the Higgs sector !!!

can probe the Higgs self-coupling



can help to reconstruct the electroweak symmetry breaking potential

may reveal the doublet nature of the Higgs by means of the $hhVV$ coupling

Search channel for Higgs pair production

Channel	BR(%)	Events with 3 ab ⁻¹	
bbbb	~ 33	40080	Huge hadronic BG
bbWW	~ 25	30000	Huge ttbar BG
bbττ	~ 7.3	9000	
WWWW	~ 4.3	5200	
bbYY	~ 0.27	5200	
bbZZ(eemm)	~ 0.015	19	

Machine Learning approaches to the Higgs boson self coupling

① BDT(Boosted Decision Tree) : bbYY

1. Phys.Rev. D96 (2017) no.3, 035022 (Alves, Alexandre et al.) arXiv:1704.07395 [hep-ph]

BDT + kinematic cuts  **5 σ (4.6 σ) significance with 10 %(20%) systematics and 3 ab⁻¹**

② (Supervising) Deep Neural Networks (DNN) : bbWW + bb $\tau\tau$

1. "Supervising Deep Neural Networks with topological augmentation in search for di-Higgs production at the LHC (Won Sang Cho, next speaker)

5 classes by the number of leptonic taus

Optimass & its compatibility distance with dim. Of vars ~ 40



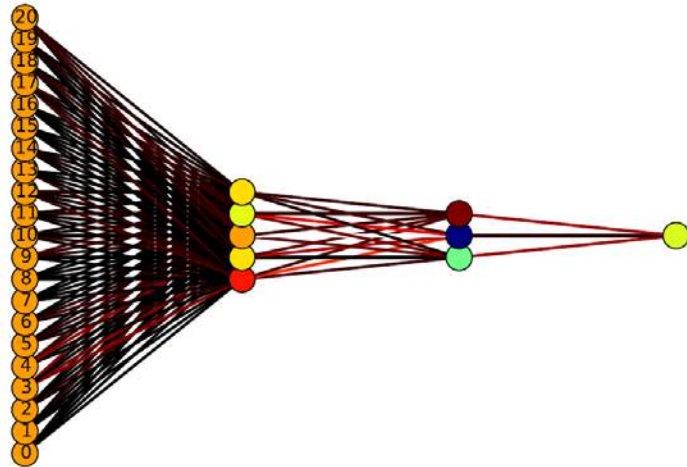
AUC of ROC = 0.991
Eff(sig)
@(Background purity=0.01) = 0.84

Machine Learning approaches to the Higgs boson self coupling

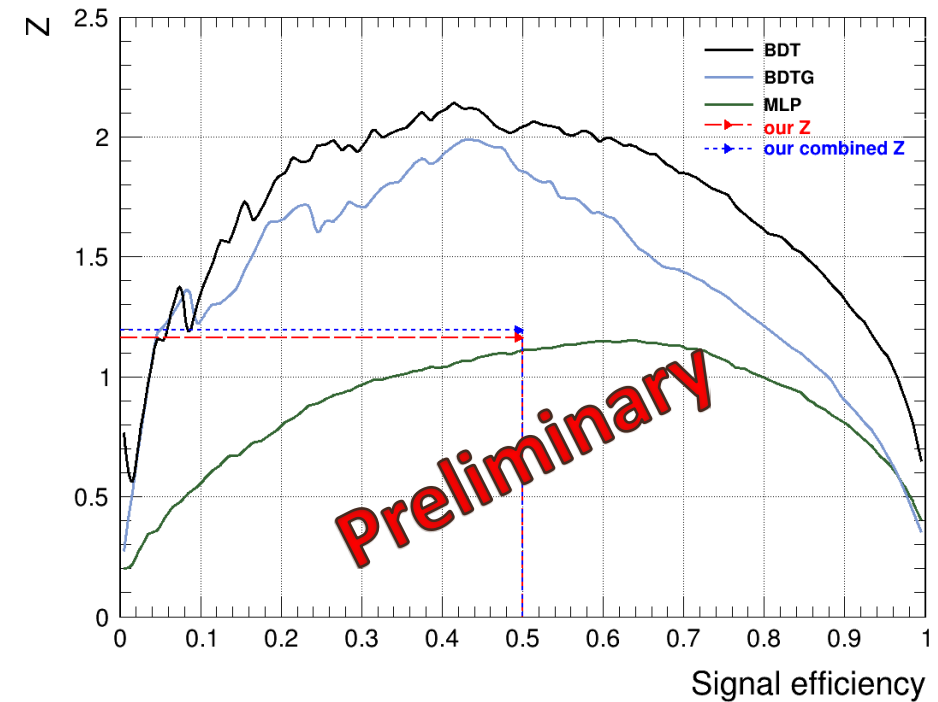
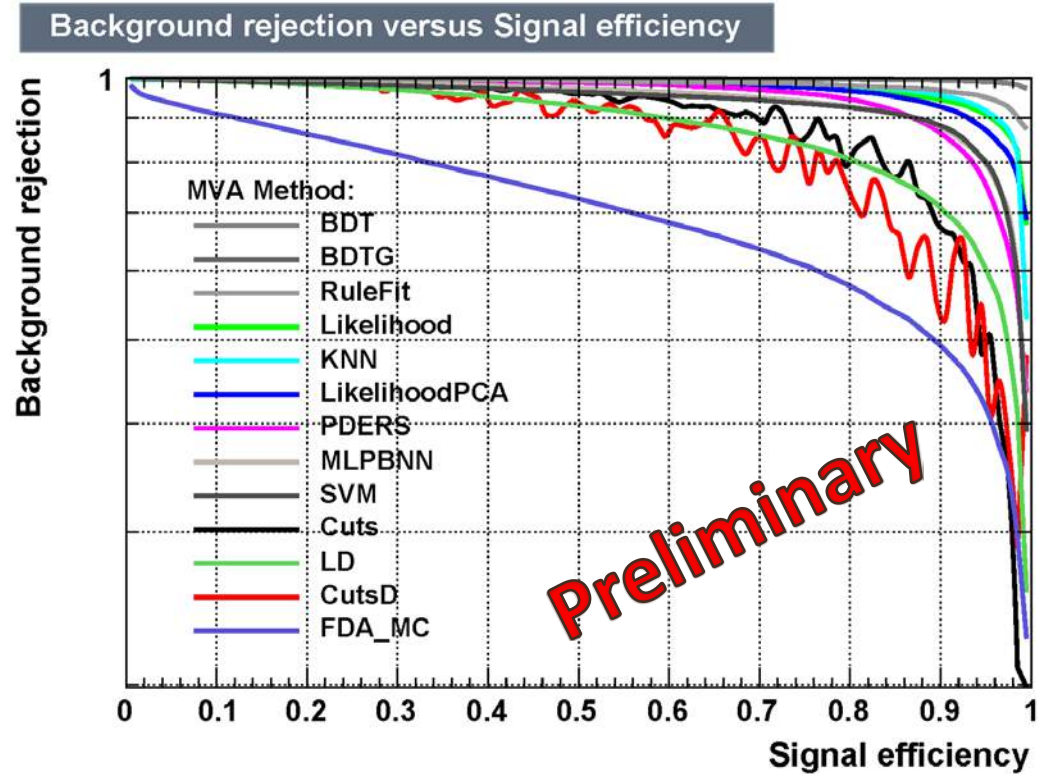
③ DNN (ANN : a multi-layer feed-forward artificial neural network) : bbbb

1. Eur. Phys. J. C (2016) 76:386 (Katharina Behr, Bortoletto et al.) arXiv:1512.08928 [hep-ph]

DNN + kinematic cuts \longrightarrow $\frac{S}{\sqrt{B}} \sim 3 \sigma$ significance with 3 ab^{-1}



Machine Learning approaches to the Higgs boson self coupling



Summary Table

Channel	Achievable Significance (σ)	Methods	Papers	Remarks
bbbb	~ 3	Kinematic Cuts+ DNN	Eur. Phys. J. C (2016) 76:386	HL-LHC (3 ab ⁻¹)
	~ (3.1 ~ 5.7)	DNN	Arxiv: 1609.002541	100 TeV FCC (10 ab ⁻¹)
bbWW		DNN	Recnt Dr. Won Sang Cho work	HL-LHC (3 ab ⁻¹)
bb $\tau\tau$				
WWWW				
bbYY	~ 5 (4.6)	Kinematic Cuts + BDT	Phys.Rev. D96 (2017) no.3, 035022	HL-LHC (3 ab ⁻¹),
	~ 2.1	Kinematic Cuts + BDT	Preriminary	With full BGs.
bbZZ(eemm)				

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bbZZ(eemm)				

Conclusion

1. Higgs pair production can allow us to reconstruct the EWSB potential and to understand the nature of the EWSB mechanism !
2. The $bbYY$ channel can offer the appropriate yields and clean(?) signal.
3. Various multivariate classification methods based on machine learning techniques are used to consider the enhancement of significance in measuring the Higgs self coupling.
4. We found that the BDT-related methods (+ cut-based analysis) can give the best results compared with other methods.
5. Presently, we are checking consistencies of our methods.

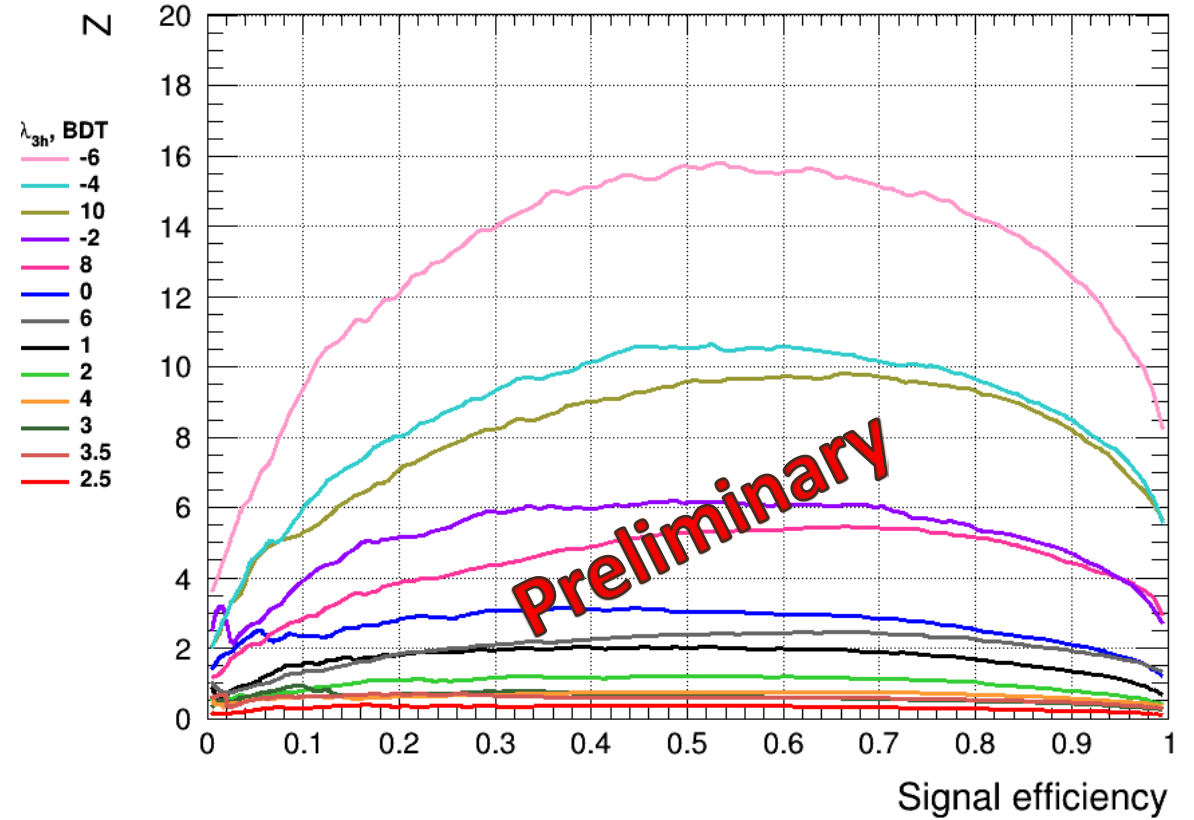
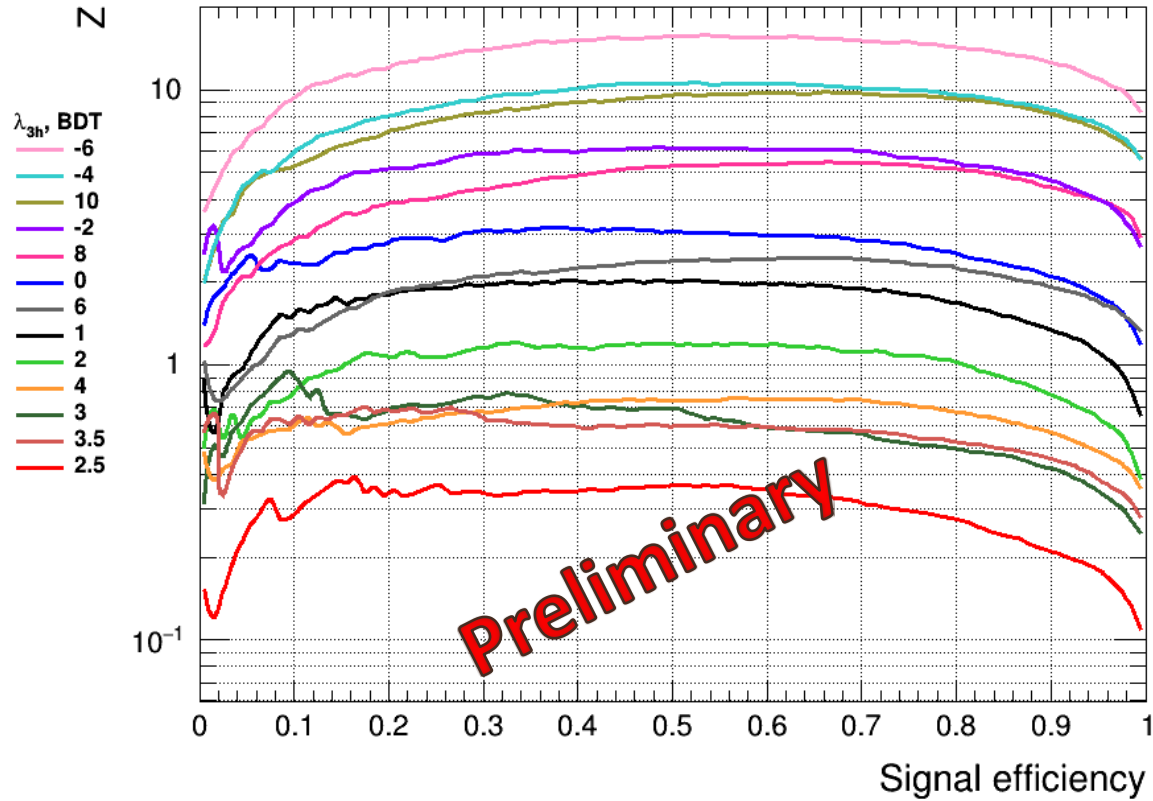
Conclusion

1. Higgs pair production can allow us to reconstruct the EWSB potential and to understand the nature of the EWSB mechanism !
2. The bbYY channel can offer the appropriate yields and clean $\gamma\gamma$ signal.
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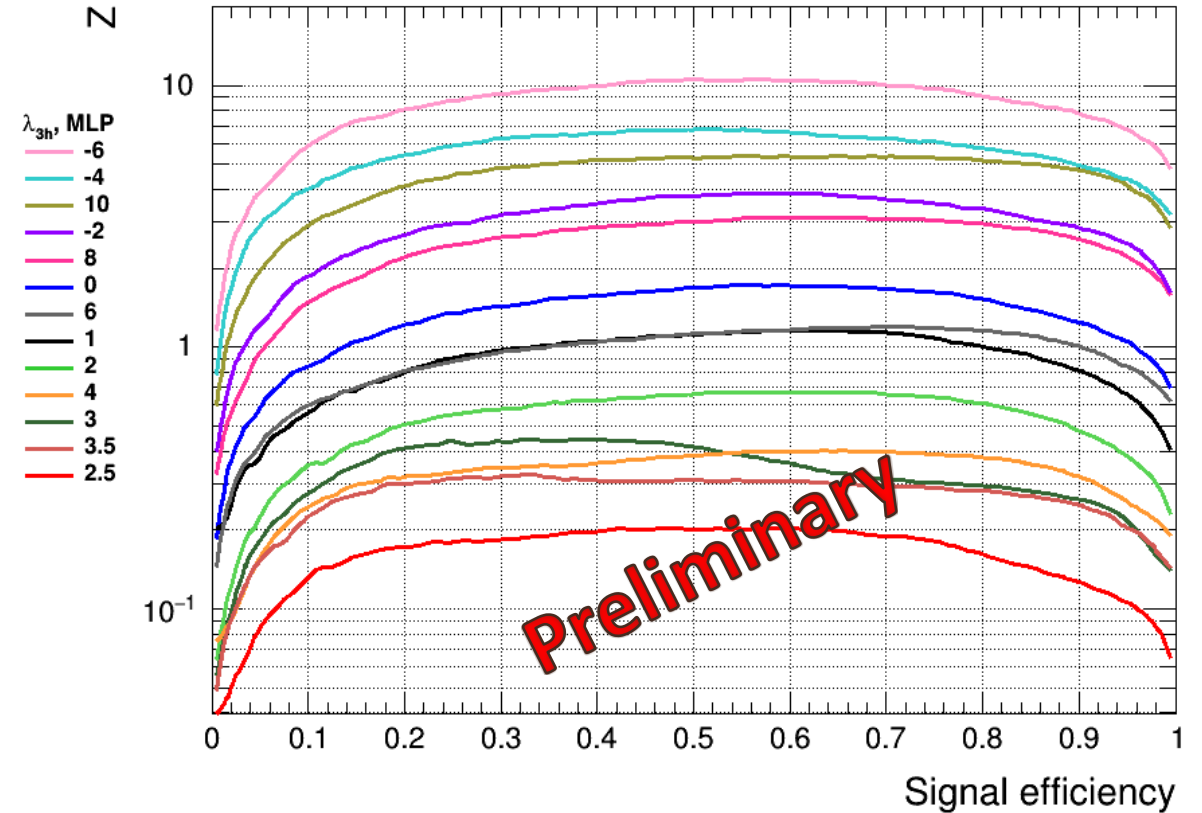
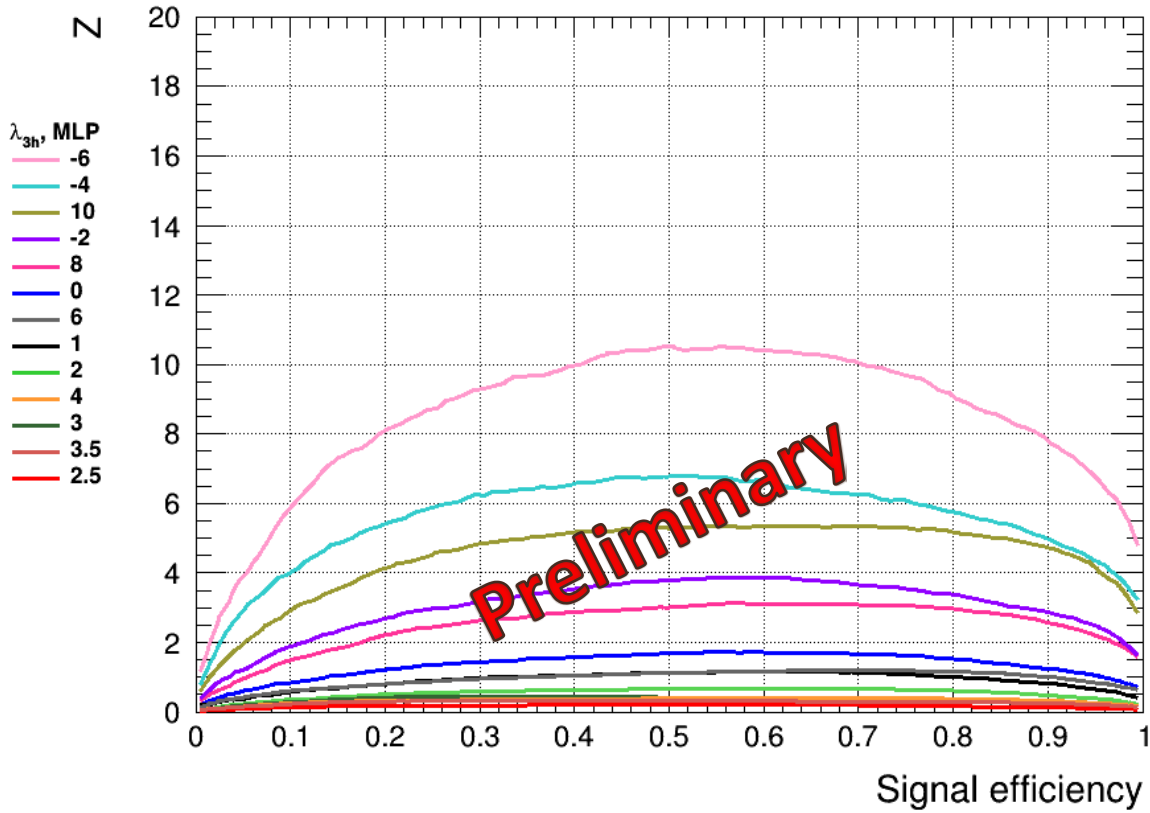
Thank you for your attention !

BACKUP SLIDES

λ dependency with BDT

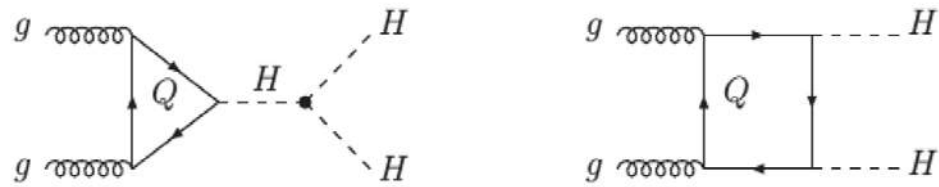


λ dependency with MLP

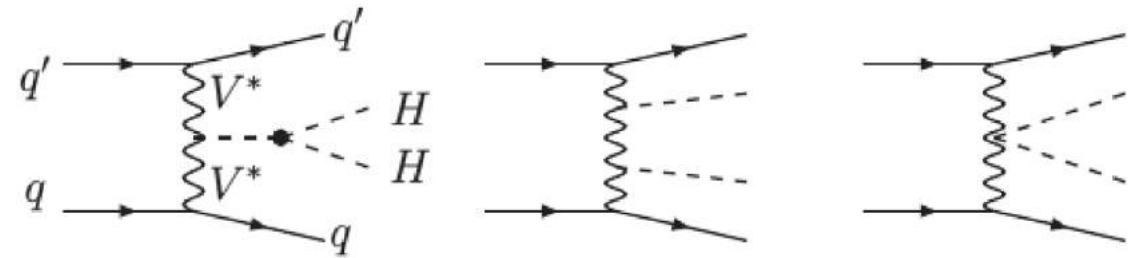


Higgs pair productions

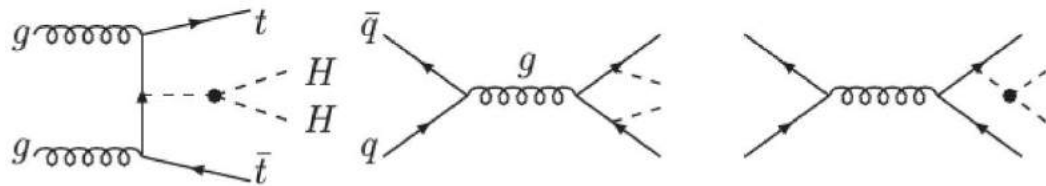
Gluon Fusion



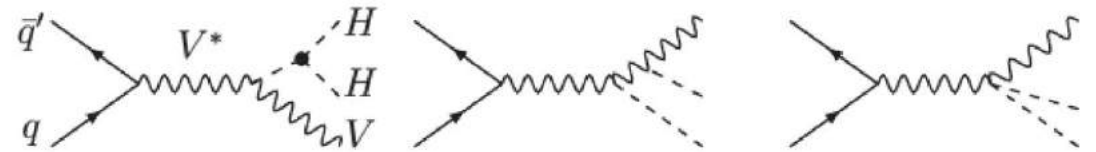
Vector Boson Fusion



Top associated productions



Higgs strahlung



Future colliders

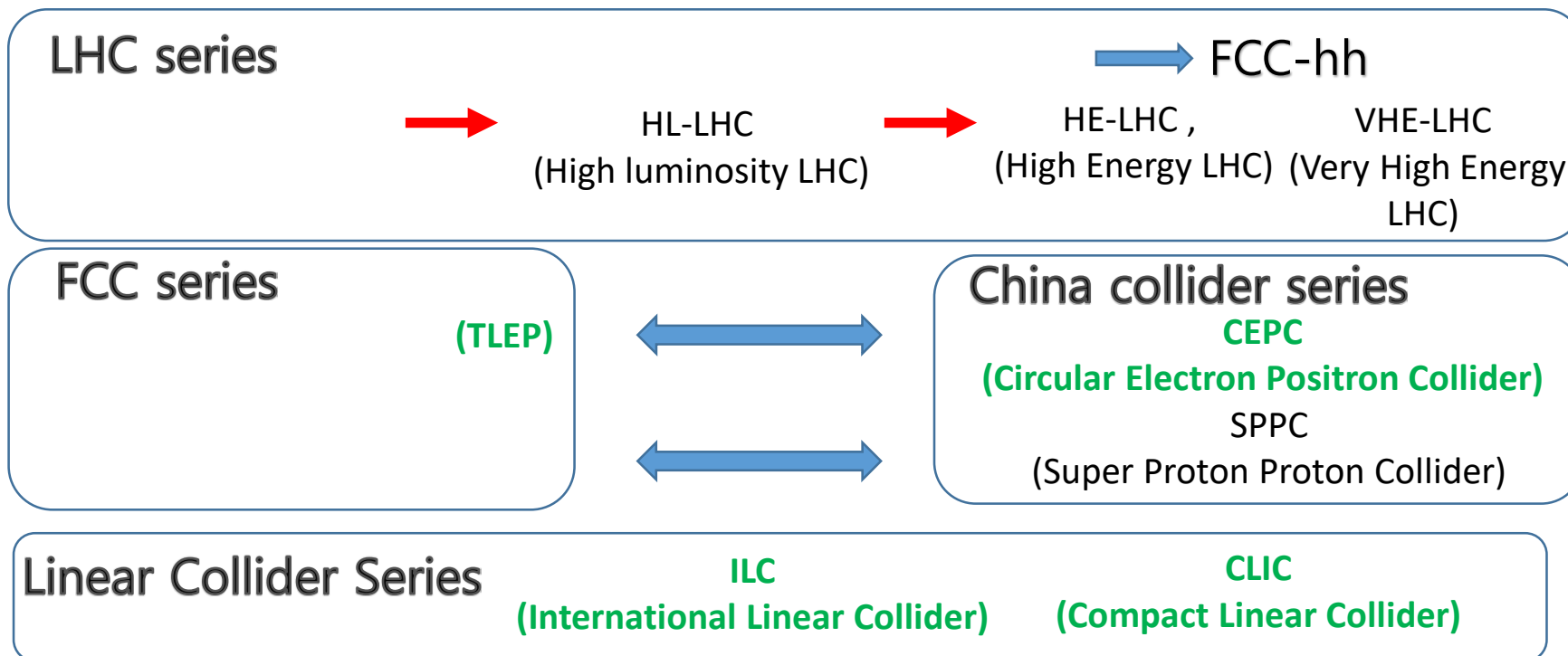


Table 1-1. Proposed running periods and integrated luminosities at each of the center-of-mass energies for each facility.

Facility	HL-LHC	ILC	ILC(LumiUp)	CLIC	FCC-ee		
					TLEP (4 IPs)	HE-LHC	VLHC
\sqrt{s} (GeV)	14,000	250/500/1000	250/500/1000	350/1400/3000	240/350	33,000	100,000
$\int \mathcal{L} dt$ (fb ⁻¹)	3000/expt	250+500+1000	1150+1600+2500	500+1500+2000	10,000+2600	3000	3000
$\int dt$ (10 ⁷ s)	6	3+3+3	(ILC 3+3+3) + 3+3+3	3.1+4+3.3	5+5	6	6