

*Studies on Neutrino-Electron Elastic
Scattering in the Standard Model
and Beyond*



Muhammed Deniz^{1,2}

1: IoP, Academia Sinica, Taiwan

2: KTU & METU, Turkey

On behalf of TEXONO Collaboration



Institute of Physics, Academia Sinica

Taipei, TAIWAN, 24 Jan 2011

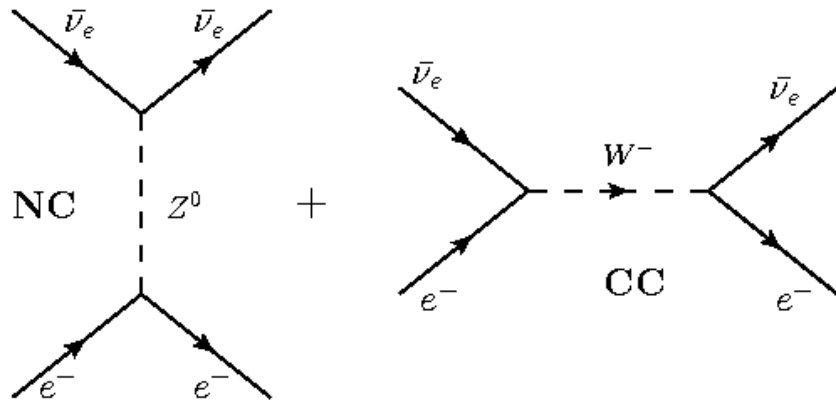
OUTLINE

- Theory overview $\bar{\nu}_e - e^-$ Scattering – Motivation
- TEXONO Physics Program
- TEXONO Experiment – CsI(Tl) Array
 - Event Selection & Data Analysis Outline
 - Background Understanding & Suppression
 - Analysis Results
- Cross Section & EW Parameters – World Status
- Probing New Physics – NSI & UP with $\bar{\nu}_e - e^-$
- Summary

$\bar{\nu}_e - e^-$ Scattering Formalism

$$\bar{\nu}_e + e^- \longrightarrow \bar{\nu}_e + e^-$$

- A basic SM process with **CC, NC & Interference**
- Not well-studied in reactor energy range \sim **MeV**



2

$$(R_{CC} : R_{NC} : R_{Int})$$

$$R_{SM}(\bar{\nu}_e e) \rightarrow (0.77 : 0.92 : -0.69)$$

$$R_{SM}(\nu_e e) \rightarrow (1.83 : 0.17 : -0.99)$$

$$\delta[\sin^2 \theta_W] \sim \begin{cases} 0.14 \cdot \delta[\xi(\bar{\nu}_e e)] \\ 0.32 \cdot \delta[\xi(\nu_e e)] \end{cases}$$

$$\xi = \frac{R_{expt}(\nu)}{R_{SM}(\nu)}$$

$$\mathcal{L}^{NC} = -\frac{G_F}{\sqrt{2}} [\bar{\nu}_e \gamma^\alpha (1 - \gamma_5) \nu_e] [\bar{e} \gamma_\alpha (g_V - g_A \gamma_5) e]$$

$$\mathcal{L}^{CC} = -\frac{G_F}{\sqrt{2}} [\bar{e} \gamma^\alpha (1 - \gamma_5) \nu_e] [\bar{\nu}_e \gamma_\alpha (1 - \gamma_5) e]$$

$$\frac{d\sigma_{SM}}{dT}(\bar{\nu}_e e) = \frac{G_F^2 m_e}{2\pi} \left[\begin{aligned} &(g_V - g_A)^2 + (g_V + g_A + 2)^2 \left(1 - \frac{T}{E_\nu}\right)^2 \\ &- (g_V - g_A)(g_V + g_A + 2) \frac{m_e T}{E_\nu^2} \end{aligned} \right]$$

$$g_A = -\frac{1}{2}$$

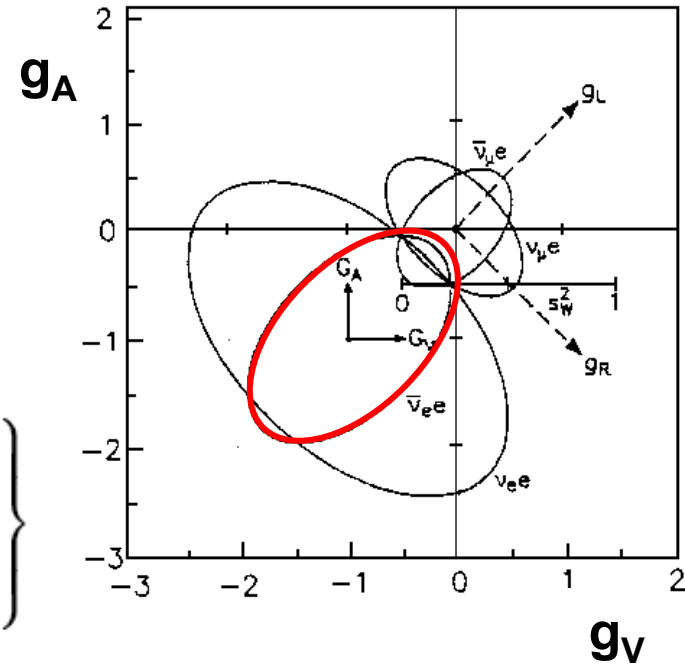
$$g_V = 2 \sin^2 \theta_W - \frac{1}{2}$$

Neutrino-Electron Scattering Cross-Section



$$\frac{d\sigma_{SM}}{dT}(\bar{\nu}_e e) = \frac{G_F^2 m_e}{2\pi} \left\{ \begin{array}{l} (g_V + 1)^2 \left[1 + \left(1 - \frac{T}{E_\nu}\right)^2 - \frac{m_e T}{E_\nu^2} \right] \\ + (g_A + 1)^2 \left[1 + \left(1 - \frac{T}{E_\nu}\right)^2 + \frac{m_e T}{E_\nu^2} \right] \\ - (g_V + 1)(g_A + 1) \left[2 - 2 \left(1 - \frac{T}{E_\nu}\right)^2 \right] \end{array} \right\}$$

$$\frac{d\sigma_{SM}}{dT}(\bar{\nu}_e e) = \frac{G_F^2 m_e}{2\pi} \left\{ \begin{array}{l} 4 \sin^4 \theta_W \left[1 + \left(1 - \frac{T}{E_\nu}\right)^2 - \frac{m_e T}{E_\nu^2} \right] \\ + 4 \sin^2 \theta_W \left[\left(1 - \frac{T}{E_\nu}\right)^2 - \frac{m_e T}{2E_\nu^2} \right] + \left(1 - \frac{T}{E_\nu}\right)^2 \end{array} \right\}$$



Expected event rate for $\bar{\nu}_e e$ scattering can be written as

$$r_{SM}(T) = \int_{E_\nu} \frac{d\sigma_{SM}}{dT} \frac{d\phi}{dE_\nu} dE_\nu$$

$$R_{SM} = \rho_e \int_T r_{SM}(T) dT = \rho_e \int_T \int_{E_\nu} \left[\frac{d\sigma}{dT} \right]_{SM} \frac{d\phi}{dE_\nu} dE_\nu dT = 0.0128 \text{ day}^{-1} \times kg^{-1} \text{ in } 3 - 8 \text{ MeV}$$

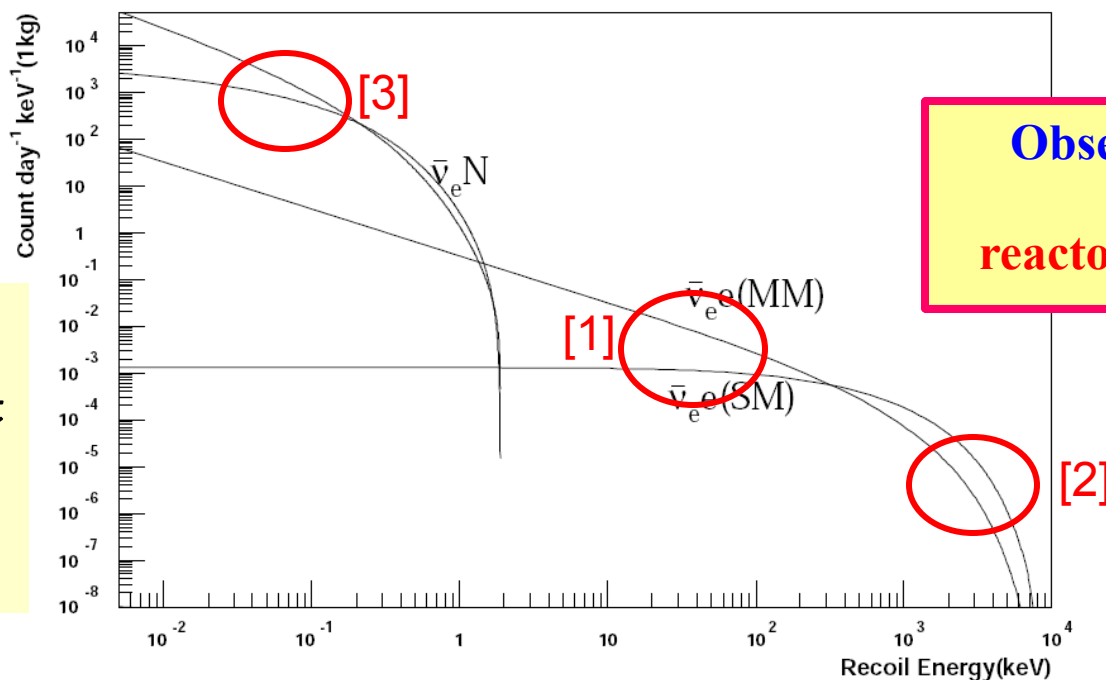
TEXONO Physics Program

TEXONO Collaboration: Taiwan (AS, INER, KSNPS, NTU);
China (IHEP, CIAE, THU, NKU, SCU, LNU); Turkey (METU, KTU); India (BHU)
Program: Low Energy Neutrino & Astroparticle Physics

quality

Detector requirements

mass



Observable Spectrum
with typical
reactor neutrino "beam"

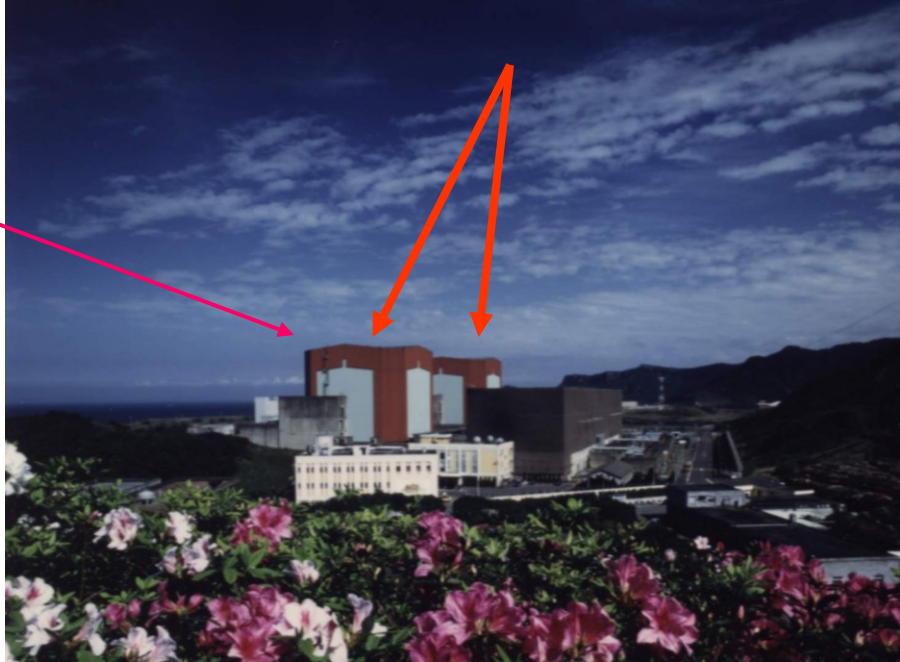
Taiwan
EXperiment
On
Neutrino

- [1] Magnetic Moment Search at ~ 10 keV \rightarrow PRL 2003, PRD 2007
- [2] Cross-Section and EW Parameters measurement at MeV range \rightarrow PRD 2010
- [3] $\bar{\nu}_e N$ Coherent Scattering & WIMP Search at sub keV range \rightarrow PRD-R 2009

Kou-Sheng Reactor Power Plant



KS NPS -II : 2 cores × 2.9 GW

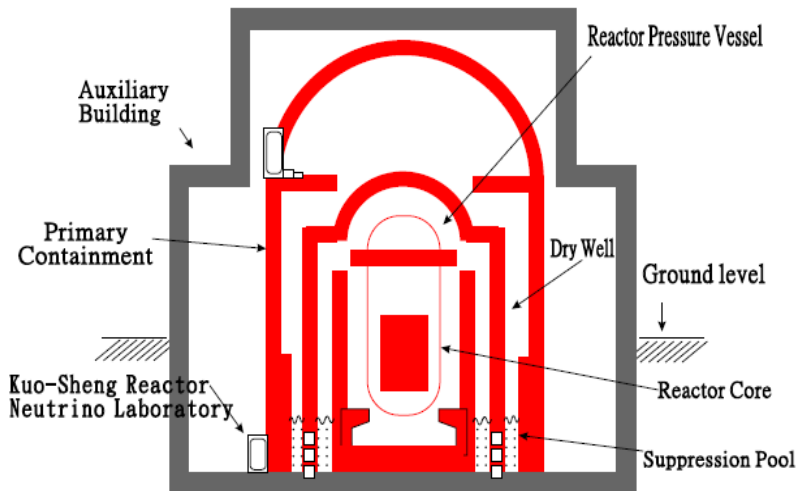


Kuo-Sheng Nuclear Power Station : Reactor Building

Total flux about $6.4 \times 10^{12} \text{ cm}^{-2} \text{ s}^{-1}$

KS v Lab: 28m from core #1

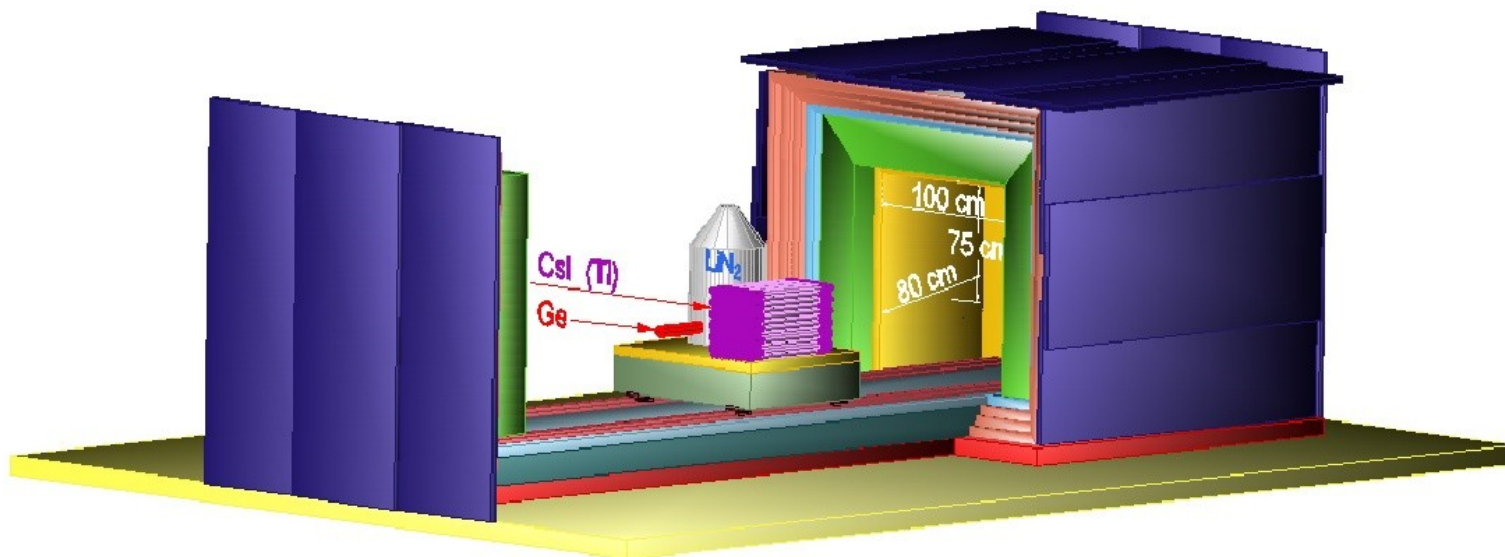
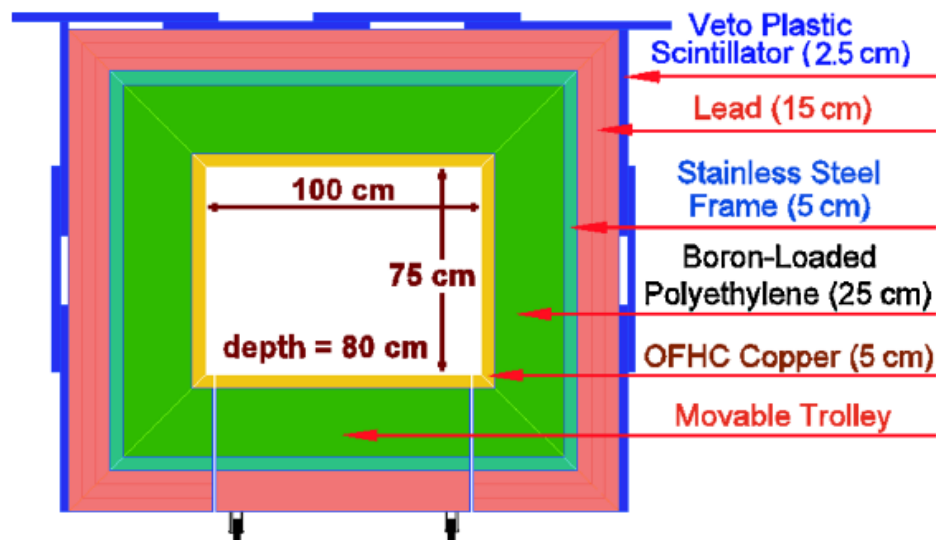
**10 m below the surface
30 mwe overburden**



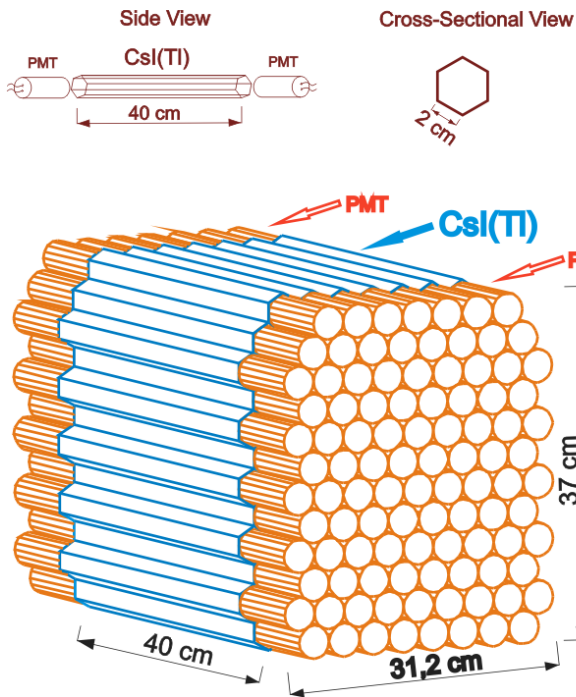
Neutrino Laboratory



Inner Target Volume & Shielding



CsI Scintillating Crystal Array



CsI(Tl) Detector
9x12 Array 200 kg



Experimental Approach; CsI(Tl) Crystal Scintillator Array:

- proton free target (suppress $\bar{\nu}_e$ -p background)
- scale to 9 (tons) design possible
- good energy resolution, **alpha & gamma**
- **Pulse Shape Discrimination (PSD)**
- allows measure **energy, position, multiplicity**
- more information for

➤ **background understanding & suppression**

❑ **DAQ Threshold: 500 keV**

❑ **Analysis Threshold: 3 MeV**

(less ambient **background** & **reactor** $\bar{\nu}_e$ spectra well known)

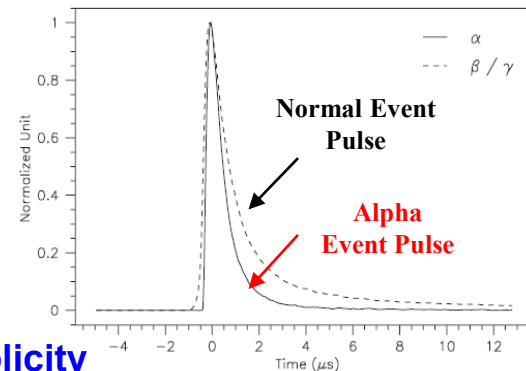
❑ **Data Volume: ~ 29883 kg-day / 7369 kg-day Reactor ON/OFF**

◆ **Energy : Total Light Collection**

◆ $\sigma(E) \sim 6\% @ E > 660 \text{ keV}$

◆ **Z-position : The variation of Ratio**

◆ $\sigma(Z) \sim 1.3 \text{ cm} @ E > 660 \text{ keV}$

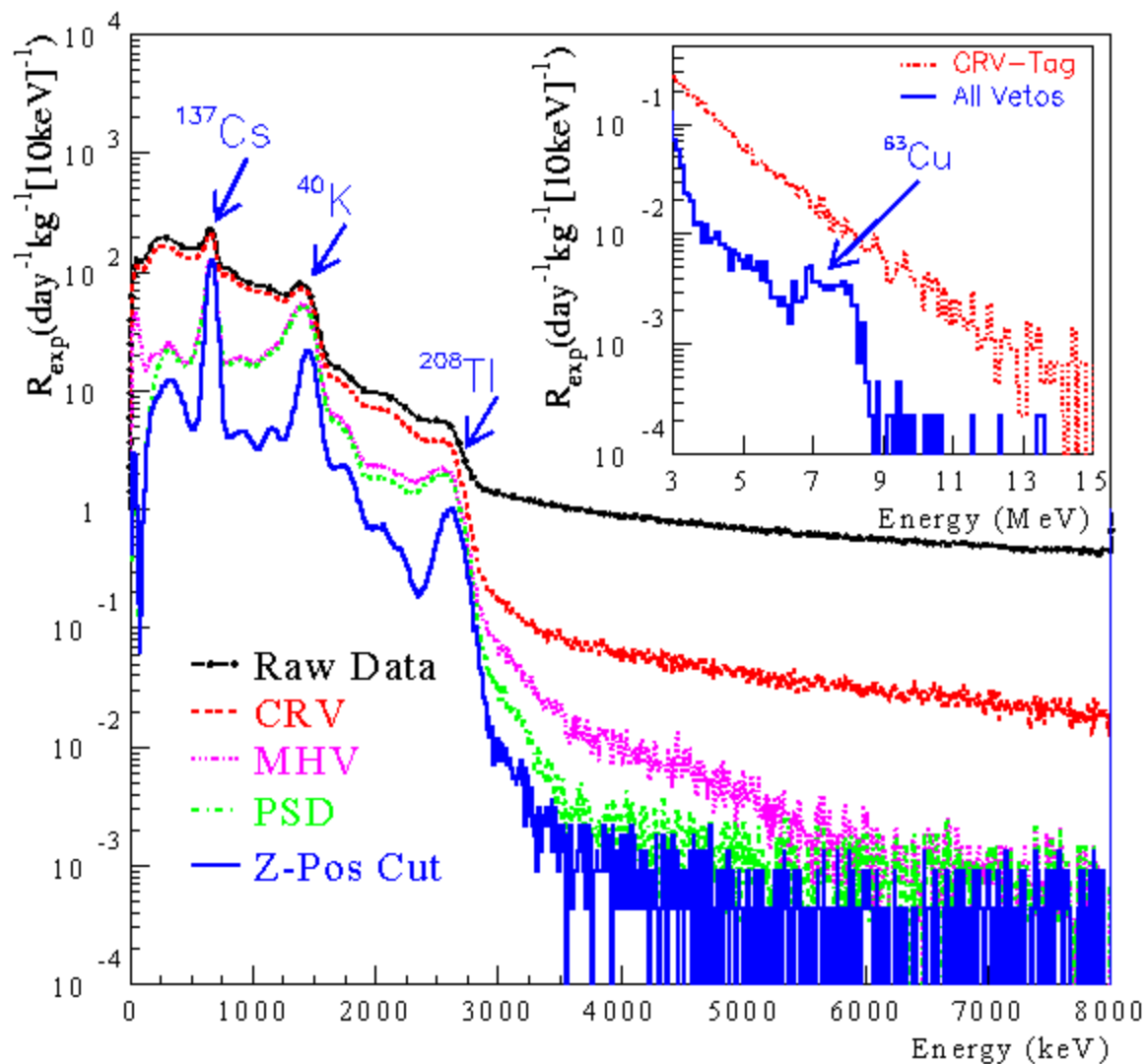


$$E \approx \sqrt{Q_L \times Q_R}$$

$$Z \approx (Q_L - Q_R) / (Q_L + Q_R)$$

Data Analysis: Event Selection

Reactor OFF



CUTS (3 - 8 MeV)	Efficiencies DAQ Live Time Eff.
CRV	92.7 %
MHV	99.9 %
PSD	~100 %
Z-pos	80%
Total	77.1 %

$$\frac{S}{B} \cong \frac{1}{30} \text{ at } 3 \text{ MeV}$$

Background Understanding

A. Radioactive Contaminants

➤ Decays of radioactive contaminants mainly ^{232}Th and ^{238}U decay chain produce background in the region of interest. Estimate the abundance of ^{137}Cs , ^{238}U and ^{232}Th inside the detector.

IDEA: By monitoring the **timing and position** information related β - α or α - α events can provide distinct signature to identify the decay process and the consistency of the isotopes involved.

B. Environmental Backgrounds

➤ Cosmic Ray muons, Products of cosmic ray muons, Spallation neutrons and High Energy γ 's from such as ^{63}Cu , ^{208}Tl

IDEA: **multiple-hit** analysis can give us very good understanding ^{208}Tl , High Energy γ and **cosmic** related background in the region of interest.

➤ **Cosmic & High Energy Gamma Ray**

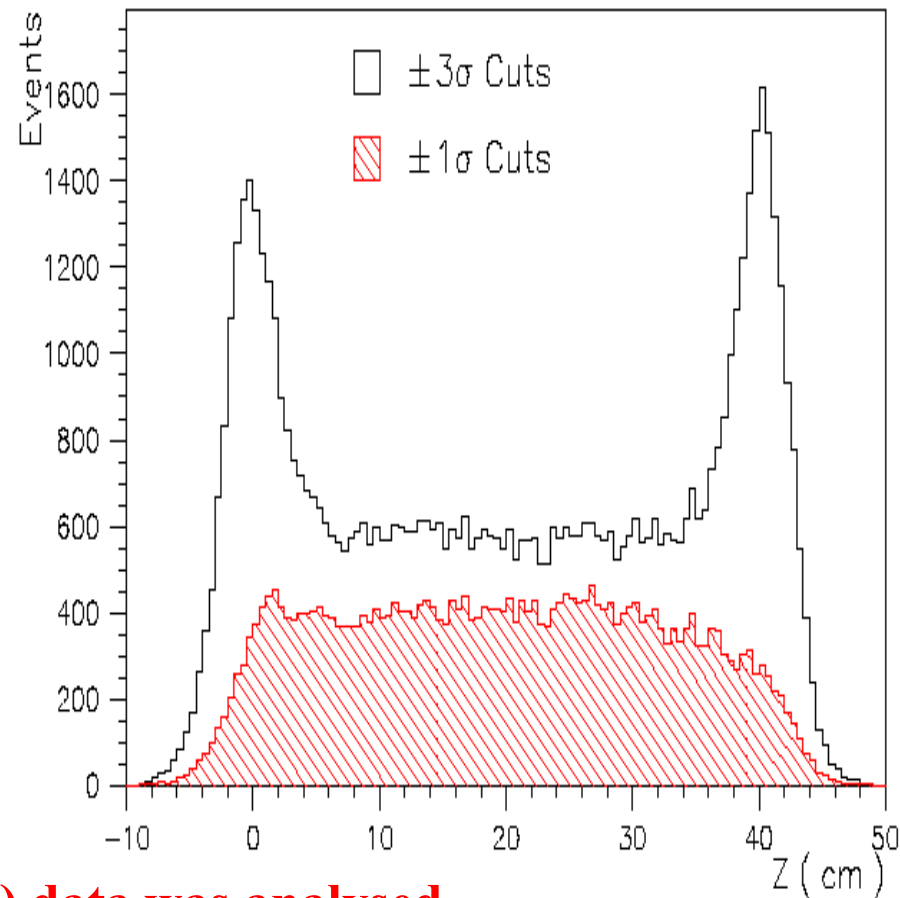
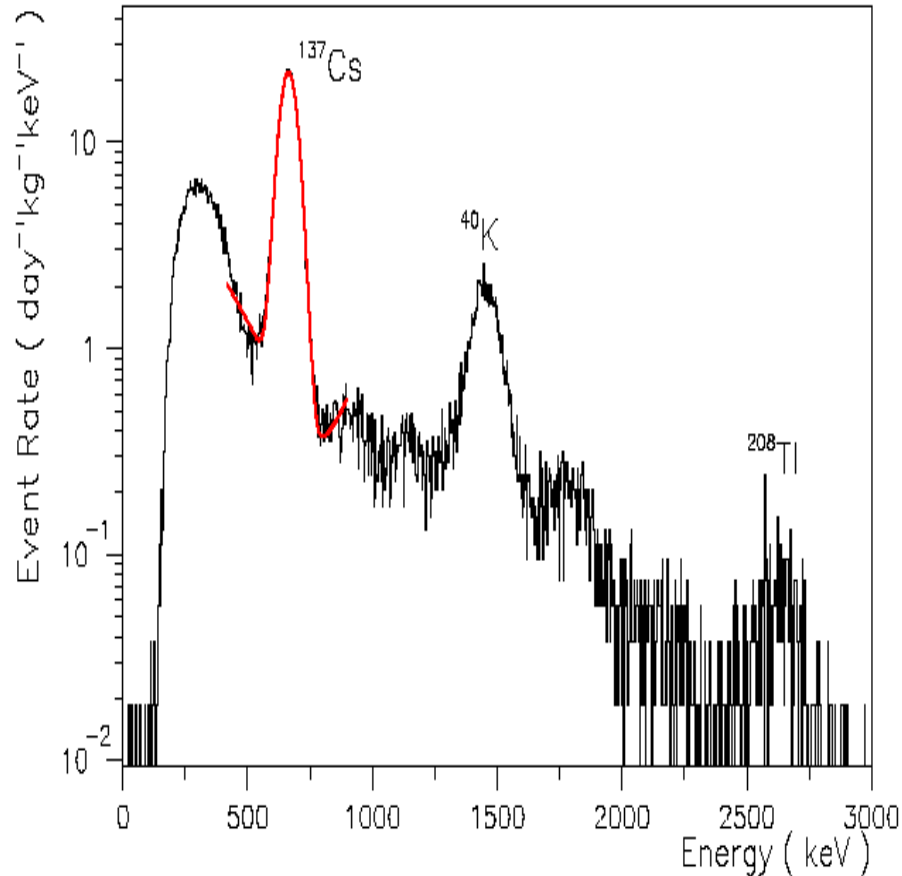
- By comparing **cosmic** and **non-cosmic** multiple-hit spectra.

➤ **Tl-208 Decay Cascade**

- By examining multiple-hit spectra as well as simulation of **Tl-208** decay chain energies to **understand/suppress** background in the region of 3-4 MeV.

Intrinsic ^{137}Cs Level

Nucl. Instr. and Meth. A 557 (2006) 490-500.



31.3 kg-day of CsI(Tl) data was analysed.

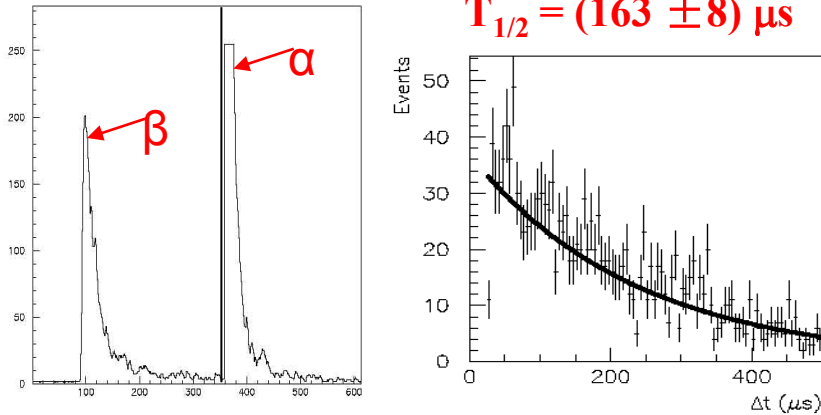
**^{137}Cs contamination level in CsI was derived ==>
 $(1.55 \pm 0.02) \times 10^{-17}$ g/g**

Intrinsic U and Th Contamination Level

Data: The total of 40 crystals with data size of 1725 kg·day was analyzed.



Selection: 1st pulse is $\gamma(\beta)$ shaped &
2nd pulse α shaped



^{238}U abundance = $0.82 \pm 0.02 \times 10^{-12}$ g/g



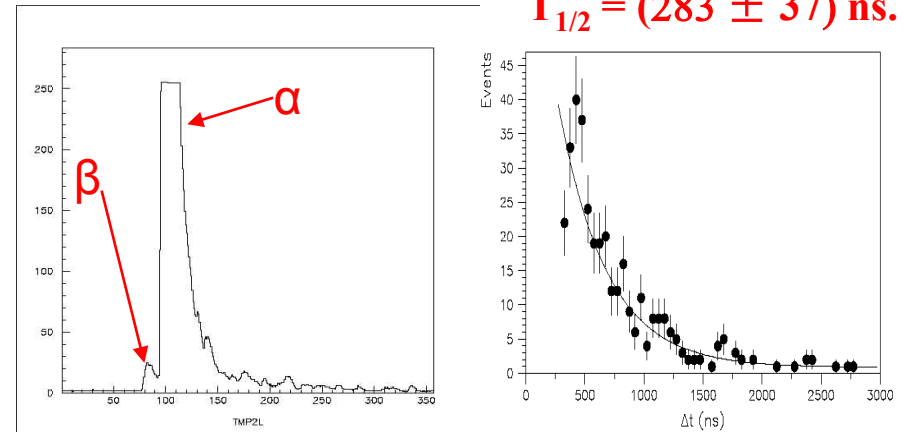
Selection: two α events with time delay less than 1s

^{232}Th abundance = $2.23 \pm 0.06 \times 10^{-12}$ g/g

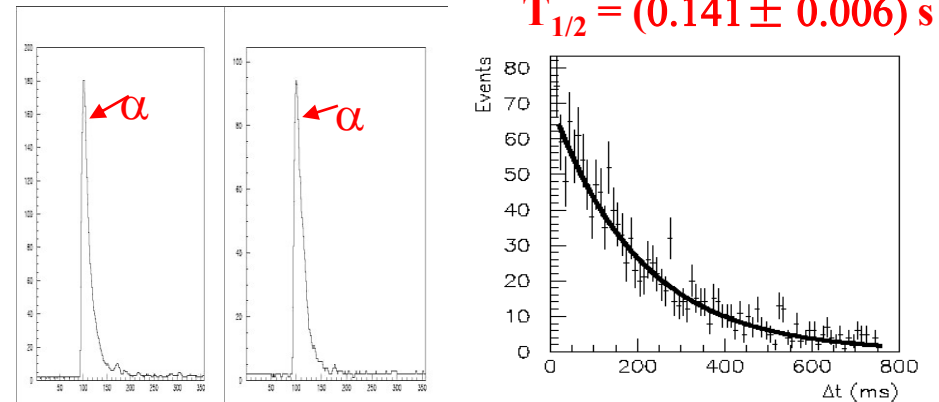


Selection: β - pulse followed by a large α pulse

$T_{1/2} = (283 \pm 37)$ ns.



^{232}Th abundance = $2.3 \pm 0.1 \times 10^{-12}$ g/g



Intrinsic Radiopurity Measurement and Contamination Level

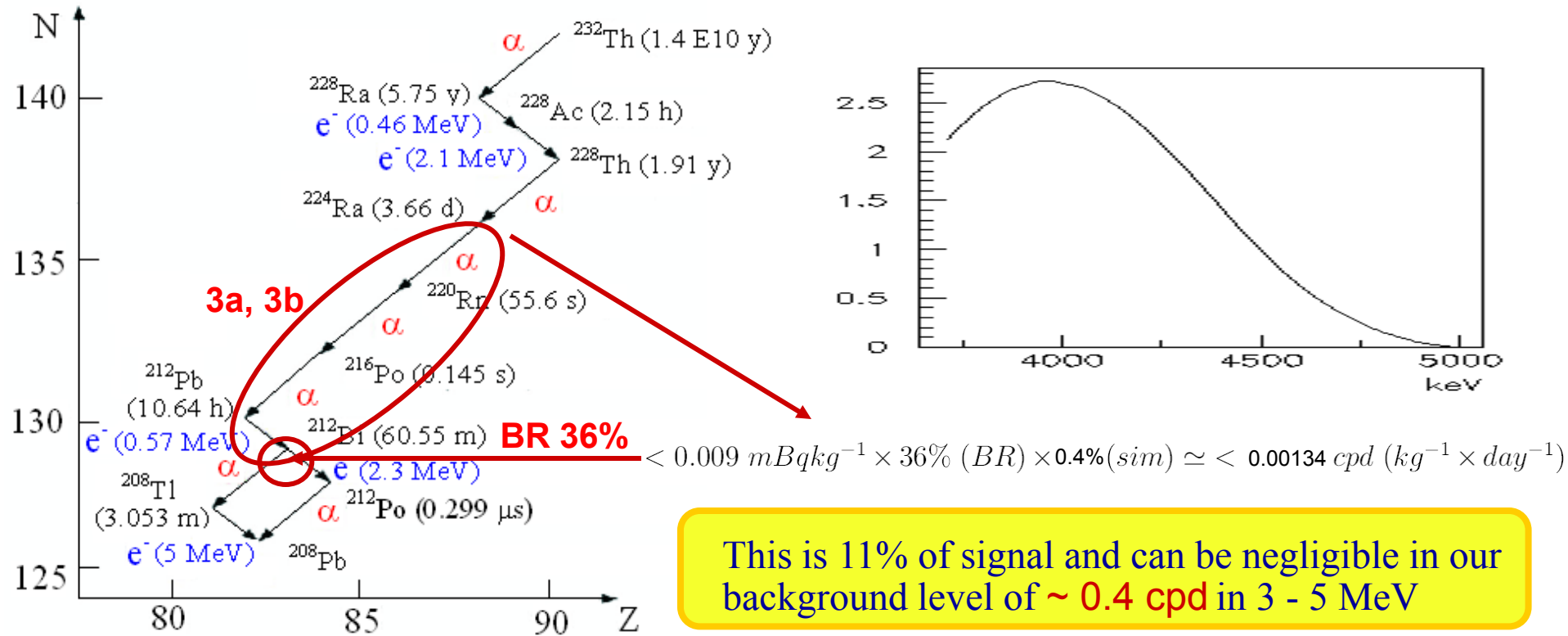
DS	Signature	Selection Efficiency	Background -to-signal	Half-life ($\tau_{1/2}$)	
				Nominal	Measured
1	$\alpha - \alpha$	0.93	0.51	3.10 min	3.2 ± 0.2 min
2	$\beta - \alpha$	0.77	3.2×10^{-3}	164 μ s	$163 \pm 8 \mu$ s
3a	$\alpha - \alpha$	0.86		55 s	54.4 ± 2.4 s
3b	α	0.97	9×10^{-5}	0.15 s	0.141 ± 0.006 s
4	$\beta - \alpha$	0.37	3×10^{-5}	299 ns	283 ± 37 ns
5a	$\alpha - \alpha$	0.78	–	3.96 s	No signal
5b	α	–	–	1.78 ms	DAQ inactive

DS	Measured activity ($mBqkg^{-1}$)	Contaminations of long-lived parents (g/g)	Contaminations of series (g/g)
1	0.0107 ± 0.0004	$^{226}Ra : (2.92 \pm 0.11) \times 10^{-19}$	$^{238}U : (0.86 \pm 0.03) \times 10^{-12}$
2	0.0102 ± 0.0003	$^{226}Ra : (2.79 \pm 0.07) \times 10^{-19}$	$^{238}U : (0.82 \pm 0.02) \times 10^{-12}$
3a, 3b	0.0090 ± 0.0002	$^{228}Th : (2.97 \pm 0.08) \times 10^{-22}$	$^{232}Th : (2.23 \pm 0.06) \times 10^{-12}$
4	0.0061 ± 0.0003	$^{228}Th : (3.1 \pm 0.2) \times 10^{-22}$	$^{232}Th : (2.3 \pm 0.1) \times 10^{-12}$
5a, 5b	< 0.003	$^{227}Th : < 1.6 \times 10^{-21}$	$^{235}U : < 4.9 \times 10^{-14}$

Estimate the background due to Intrinsic ^{208}Tl

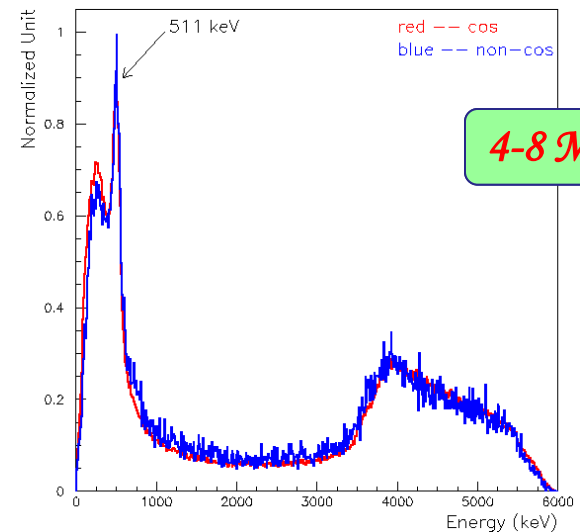
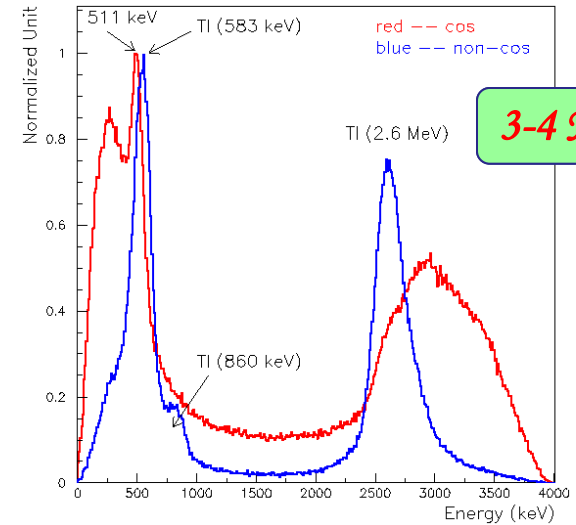
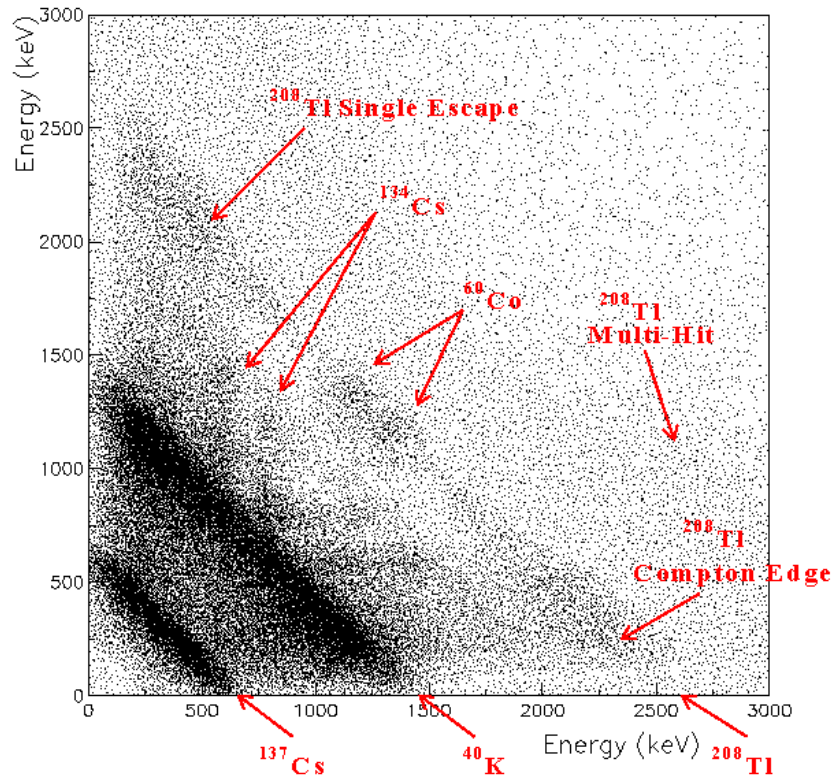
^{232}Th (decay chain)

^{208}Tl beta with associated gammas
energies deposit in one crystal.

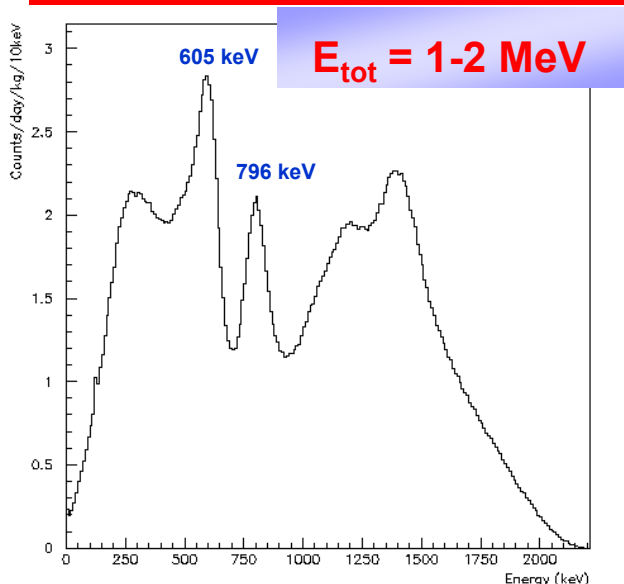


Background Understanding: via Multiple Hit Analysis

2 HIT SPECTRUM



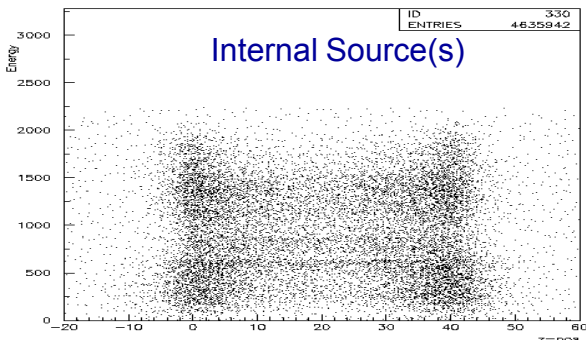
Background Understanding via Multi Hit



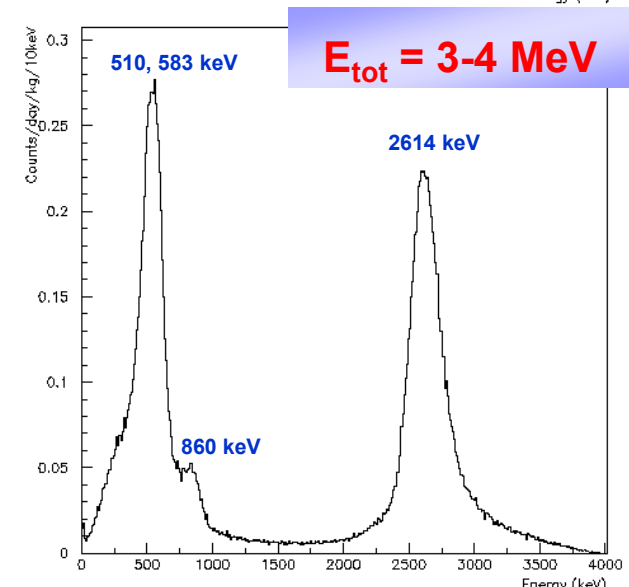
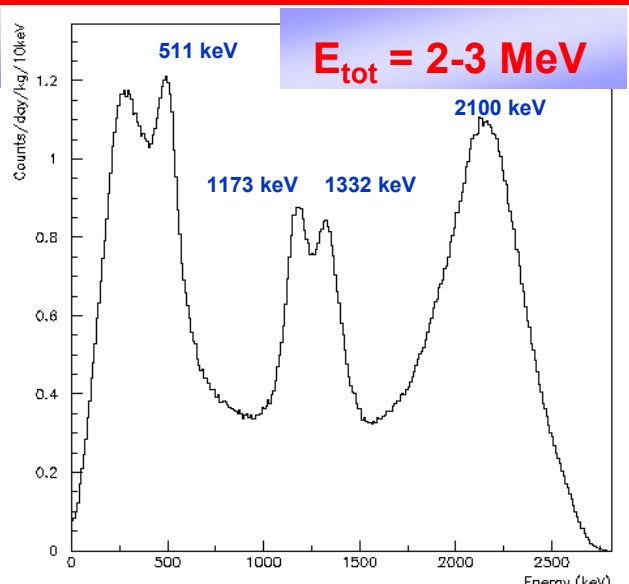
Cs-134 ($n + {}^{133}\text{Cs} \rightarrow {}^{134}\text{Cs}$)

- 605 keV 97.6%;
- 796 keV 85.5%

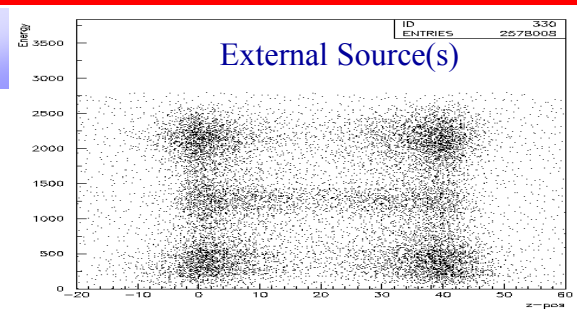
With the Q of beta decay at 2MeV



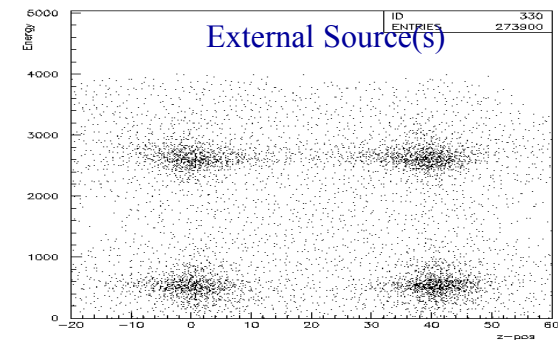
➤ Cosmic induced neutrons can be captured by the target nuclei ${}^{133}\text{Cs}$.



➤ Combination of TI gammas can affect up to around 4 MeV

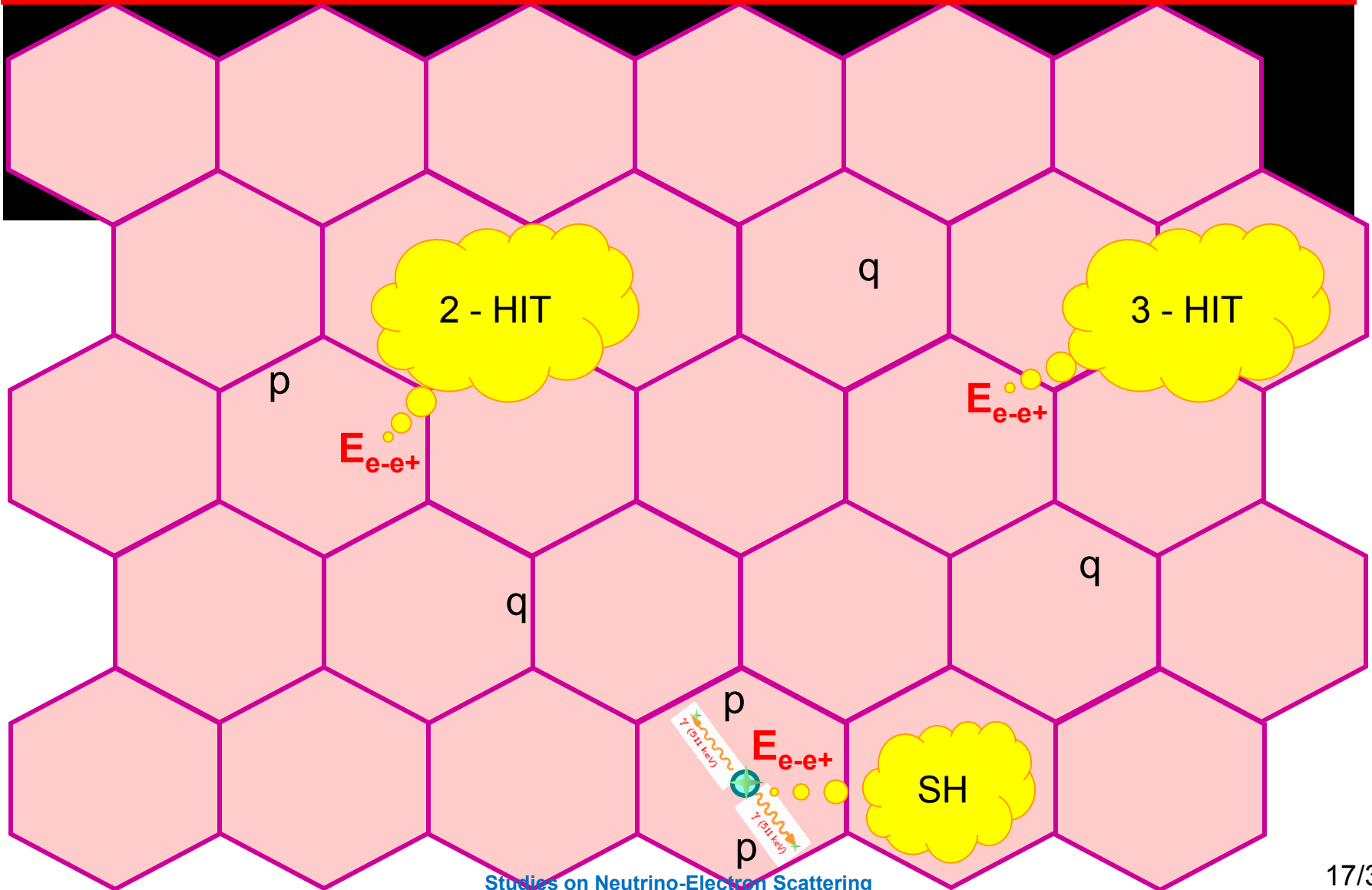


- **Co-60:** 1173.2 keV 99.86% accompanied with 1332.5 keV 99.98%
 - The background related to reactor. Mostly come from the dust.
- **TI Pair Production:** One escape peaks
 - (~ 2105 + 511 keV)



- 2614 keV 99 % accompanied with
 - 583 keV 85%
 - 510.8 keV 23%
 - 860 keV with 13%

Background Prediction via PAIR PRODUCTION



Residual Background Understanding & Suppression

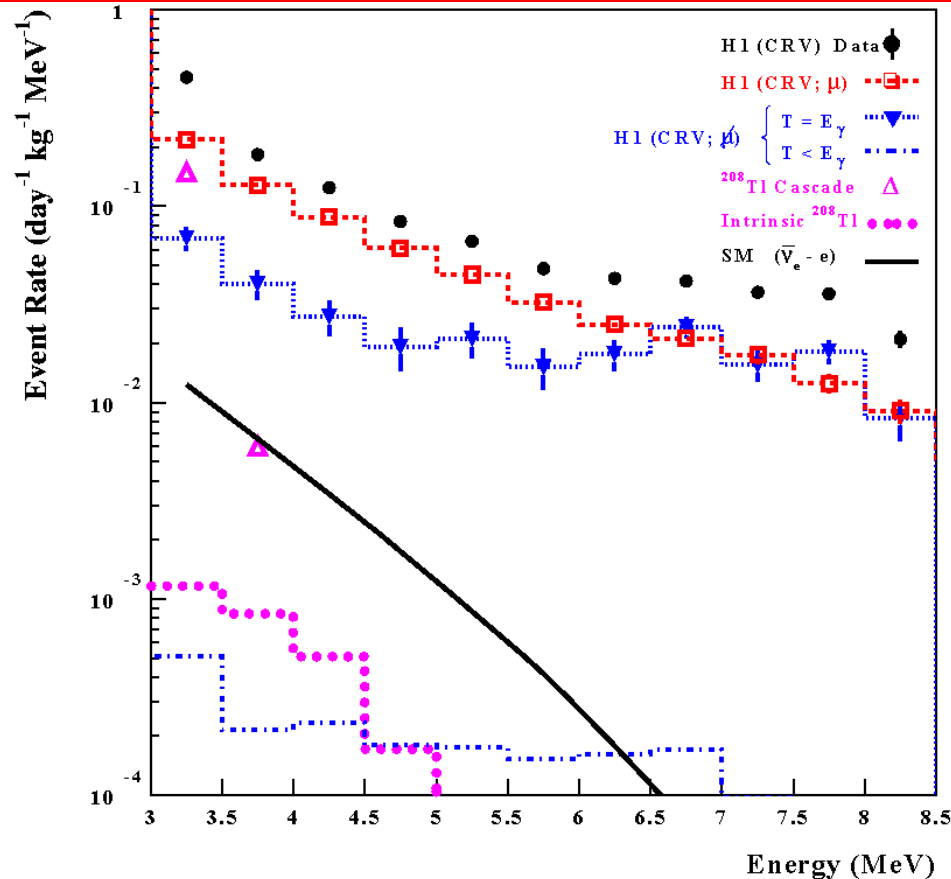
- Background Sources : **High Energy γ -rays & Cosmic Rays & ^{208}Tl**

Idea -- Use Multiple Crystal Hit (**MH**) spectra to **predict** Single Crystal Hit (**SH**) Background to the neutrino events

$$\left(\frac{MH_{non\ cos}}{MH_{tot}}\right)_{ON,OFF} = 1 - \varepsilon = \left(\frac{SH[BKG(\cos)]}{SH_{tot}}\right)_{ON,OFF}$$

$$\frac{SH[BKG(2614 + 583)]}{MH[2614 ; 583 (data)]} = \frac{SH[2614 + 583 (MC)]}{MH[2614 ; 583 (MC)]}$$

Background Understanding & Suppression



Combined **BKG(SH)** from **three measurements**:

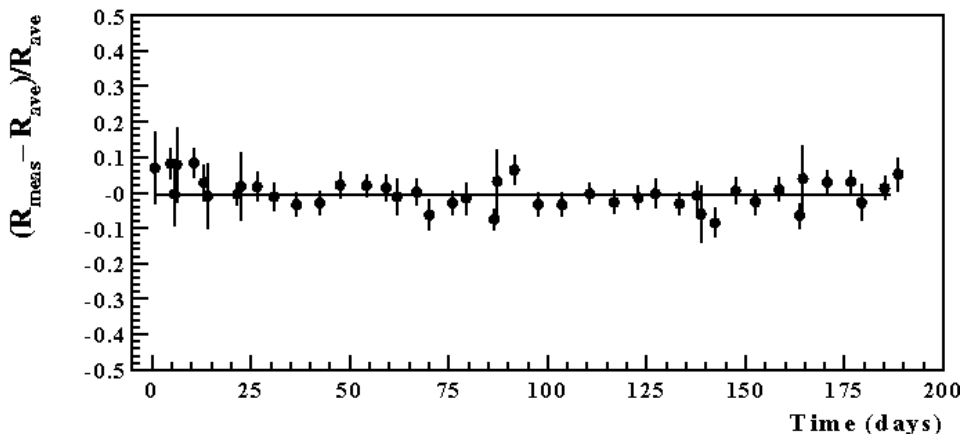
- Direct **Reactor OFF(SH)** spectra ⊕ Predicted **BKG(SH)** from **OFF(MH)**
- ⊕ Predicted **BKG(SH)** from **ON(MH)**

$$\nu = \text{ON(SH)} - \text{BKG(SH)}$$

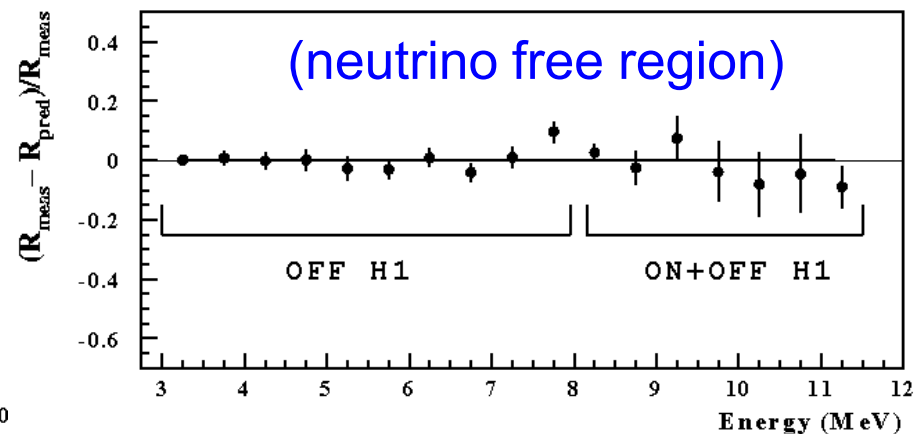
Systematic Uncertainties

Approach - Use **non- ν events** for demonstration

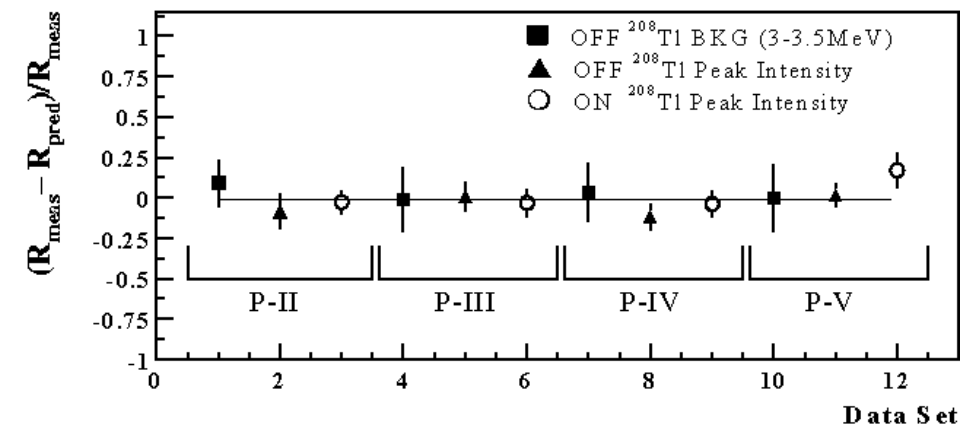
^{208}Tl Peak Events Stability



BKG - Pred.



^{208}Tl (SH) Prediction



ON-OFF Stability < $\sim 0.5\%$

Random trigger events for **DAQ & Selection Cuts**

Stability of **TI-208 (2614 keV)** peak events

Cosmic Induced BKG(SH) Prediction < $\sim 1\%$

Successfully Predict **Cosmic BKG** in **NEUTRINO FREE REGION**

TI-208 Induced BKG(SH) Prediction < $\sim 3\%$

Successfully Predict **TI-208 Induced BKG(SH) >3MeV** at Reactor **OFF** periods

Successfully Predict **TI-208** peak intensity for both Reactor **ON/OFF** with the same tools (**MC**)

Systematic Uncertainties

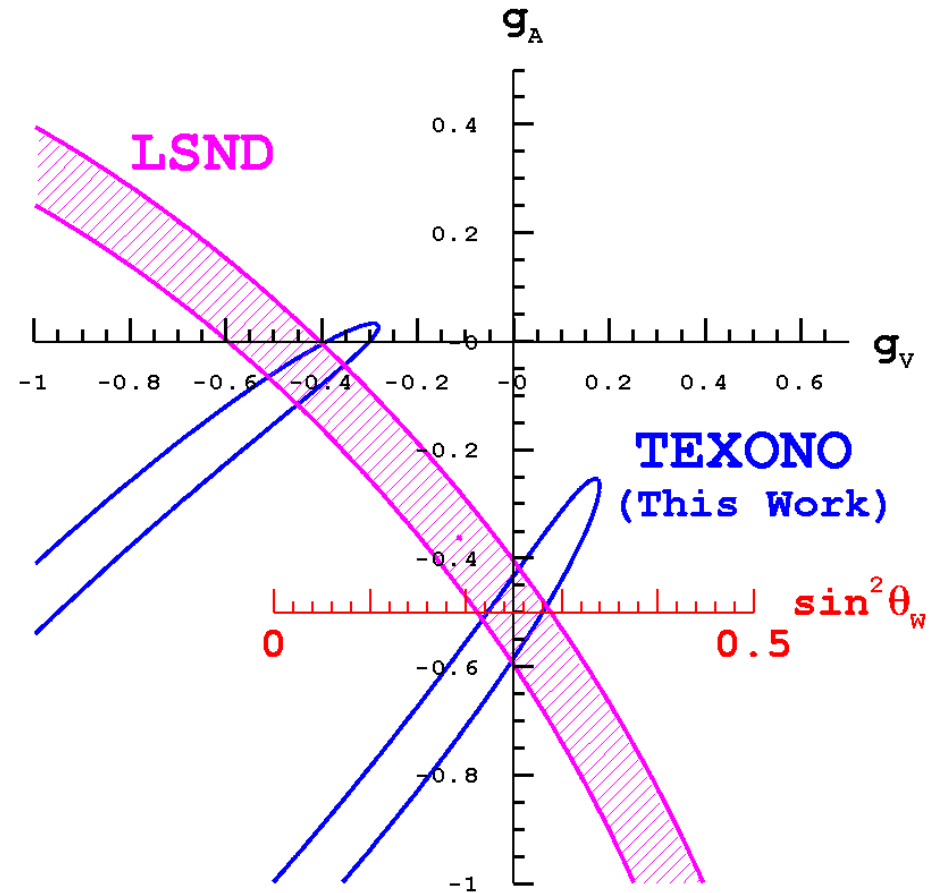
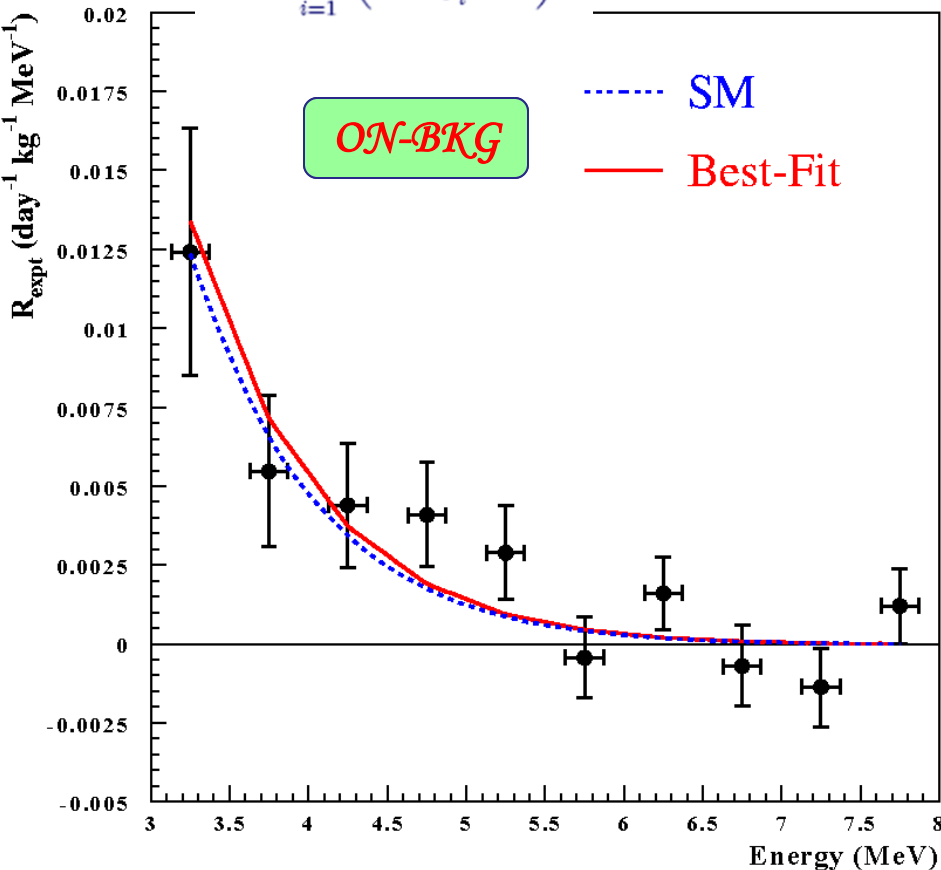
Summary of the sources of systematic errors and their contributions to the measurement uncertainties.

Sources	δ_{sys} (Source)	$\Delta_{\text{sys}} (\xi)$
Signal strength :		
Φ_ν Evaluation	<3%	<0.03
Efficiencies for neutrino events	<1.3%	<0.013
Fiducial target mass	<4%	<0.04
* Combined (signal)	-	<0.052
Background subtraction :		
Reactor OFF measurement	<0.4%	<0.06
Background evaluation		
$\odot\text{H1}(\text{CRV}; \text{Tl}_\gamma)$	<3%	<0.08
$\odot\text{H1}(\text{CRV}; \mu) + \text{H1}(\text{CRV}; \mu)$	<1%	<0.17
Net	-	<0.19
* Combined (background)	-	<0.15
Total		<0.16

Cross Section & Weak Mixing Angle

Phys. Rev. D 81, 072001 (2010)

$$\chi^2 = \sum_{i=1}^{10} \left(\frac{R_i - \zeta R_i^{SM}}{\sigma_i} \right)^2$$



$$R = [1.08 \pm 0.21 (\text{stat}) \pm 0.16 (\text{sys})] \times R_{SM}$$

$$\sin^2 \theta_w = 0.251 \pm 0.031 (\text{stat}) \pm 0.024 (\text{sys})$$

Better sensitivity is achieved in the measurement of weak mixing angle

World Status: Summary Table

	Experiment	Energy (MeV)	Events	Cross-Section	$\sin^2\theta_W$
$\nu_e - e^-$	LAMPF [Liquid Scin.]	7 - 60	236	$[10.0 \pm 1.5 \pm 0.9]$ $\times E_{\nu_e} 10^{-45} \text{cm}^2$	0.249 ± 0.063
	LSND [Liquid Scin.]	10 - 50	191	$[10.1 \pm 1.1 \pm 1.0]$ $\times E_{\nu_e} 10^{-45} \text{cm}^2$	0.248 ± 0.051
$\bar{\nu}_e - e^-$	Savannah-River [Plastic Scin.]	1.5 - 3.0 3.0 - 4.5	381 71	$[0.86 \pm 0.25] \times \sigma_{V-A}$ $[1.70 \pm 0.44] \times \sigma_{V-A}$	0.29 ± 0.05
	Savannah-River Re-analysed (PRD1989, Engel&Vogel)	1.5 - 3.0 3.0 - 4.5	N/A	$[1.35 \pm 0.4] \times \sigma_{SM}$ $[2.0 \pm 0.5] \times \sigma_{SM}$	N/A
	Krasnoyarsk (Fluorocarbon)	3.15 - 5.18	N/A	$[4.5 \pm 2.4]$ $\times 10^{-46} \text{cm}^2/\text{fission}$	0.22 ± 0.75
	Rovno [Si(Li)]	0.6 - 2.0	41	$[1.26 \pm 0.62]$ $\times 10^{-44} \text{cm}^2/\text{fission}$	N/A
	MUNU [CF ₄ (gas)]	0.7 - 2.0	68	1.07 ± 0.34 events day ⁻¹	N/A
	TEXONO [CsI(Tl) Scin.]	3 - 8	~ 410	$[1.08 \pm 0.21 \pm 0.16]$ $\times R_{SM}$	$0.251 \pm 0.031(\text{stat})$ $\pm 0.024(\text{sys})$

Neutrino-Electron Scattering Cross-Section

$$\bar{\nu}_e + e^- \longrightarrow \bar{\nu}_e + e^-$$

$$\sigma_{SM} = \sigma^{INT} + \sigma^{CC} + \sigma^{NC}$$

$$\sigma^{CC} = 4 \times I_2$$

$$\begin{aligned} \sigma^{NC} &= (g_V - g_A)^2 \times I_1 + (g_V + g_A)^2 \times I_2 + (g_A^2 - g_V^2)^2 \times I_3 \\ &= (4 \sin^4 \theta) \times I_1 + (1 - 2 \sin^2 \theta)^2 \times I_2 + 2 \sin^2 \theta (1 - 2 \sin^2 \theta) \times I_3 \end{aligned}$$

$$\begin{aligned} \sigma^{INT} &= 4(g_V + g_A) \times I_2 + 2(g_A - g_V) \times I_3 \\ &= 4(2 \sin^2 \theta - 1) \times I_2 - (4 \sin^2 \theta) \times I_3 \end{aligned}$$

$$I_1 = \frac{G_F^2 m_e}{2\pi} \rho_e \iint_{TE_\nu} \left(\frac{d\phi}{dE_\nu} \right) dE_\nu dT$$

$$I_2 = \frac{G_F^2 m_e}{2\pi} \rho_e \iint_{TE_\nu} \left(1 - \frac{T}{E_\nu} \right)^2 \left(\frac{d\phi}{dE_\nu} \right) dE_\nu dT$$

$$I_3 = \frac{G_F^2 m_e}{2\pi} \rho_e \iint_{TE_\nu} \left(\frac{m_e T}{E_\nu} \right) \left(\frac{d\phi}{dE_\nu} \right) dE_\nu dT$$

Interference, Neutrino Magnetic Moment and Charge Radius

Interference Term

$$R_{SM} = R^{CC} + R^{NC} + \eta \times R^I$$

Interference Term

$$\eta = -0.92 \pm 0.30(\text{stat}) \pm 0.24(\text{sys})$$

Neutrino Magnetic Moment

$$R(ON - BKG) = R(SM) + \mu_\nu^2 \times R(MM)$$

$$\mu_\nu^2 = [0.42 \pm 1.79(\text{stat}) \pm 1.49(\text{sys})] \cdot \mu_B^2$$

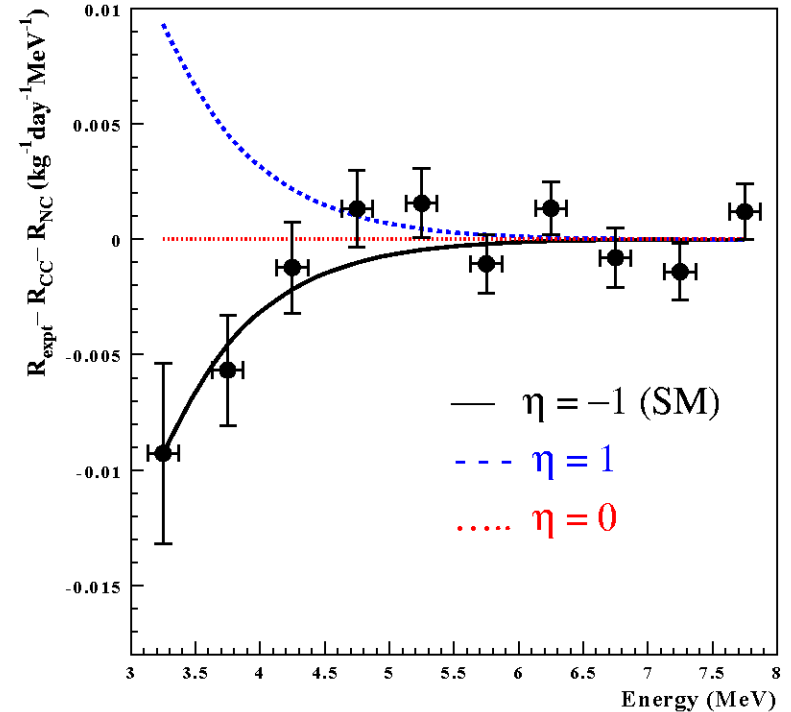
$$\mu_\nu < 2.2 \times 10^{-10} \times \mu_B$$

at 90% C. L.

Neutrino Charge Radius

$$\sin^2 \theta_W \rightarrow \sin^2 \theta_W + (\sqrt{2}\pi\alpha / 3G_F) \langle r_{\nu e}^2 \rangle$$

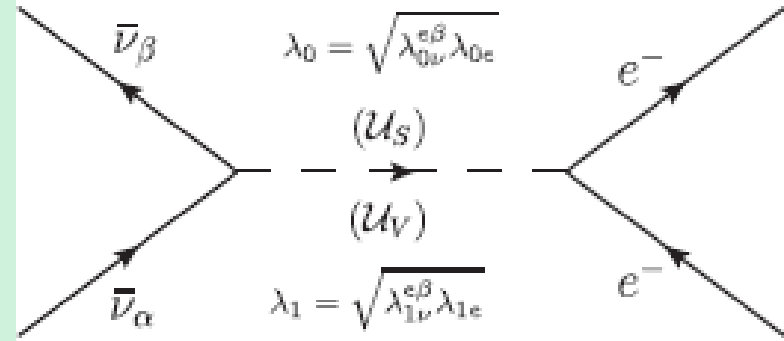
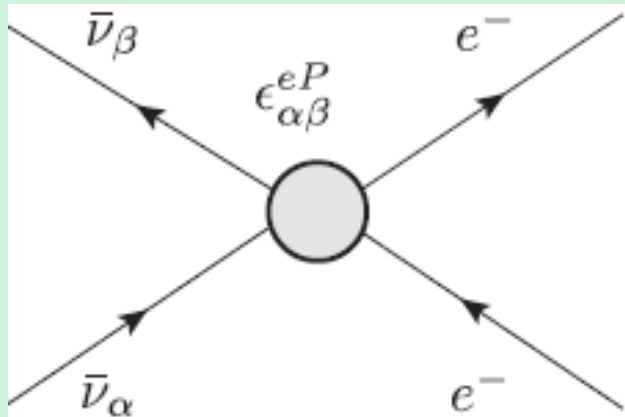
$$-2.1 \times 10^{-32} < \langle r_{\nu e}^2 \rangle < 3.3 \times 10^{-32} \text{ cm}^2$$



Constraints of New Physics with $\bar{\nu}_e - e^-$ Scattering

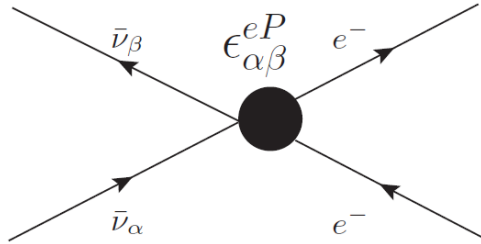
[PRD 82, 033004 (2010)]

– The measurements of neutrino scattering provide a sensitive tools to probe NSI and UP Physics



NSI of Neutrino

$$\mathcal{L}_{\text{eff}} = -\epsilon_{\alpha\beta}^{eP} 2\sqrt{2}G_F (\bar{\nu}_\alpha \gamma_\rho L \nu_\beta) (\bar{e} \gamma^\rho P e)$$



$$\frac{d\sigma(E_\nu, T)}{dT} = \frac{2G_F^2 M_e}{\pi} [(\tilde{g}_L^2 + \sum_{\alpha \neq e} |\epsilon_{\alpha e}^{eL}|^2) + (\tilde{g}_R^2 + \sum_{\alpha \neq e} |\epsilon_{\alpha e}^{eR}|^2) \left(1 - \frac{T}{E_\nu}\right)^2 - (\tilde{g}_L \tilde{g}_R + \sum_{\alpha \neq e} |\epsilon_{\alpha e}^{eL}| |\epsilon_{\alpha e}^{eR}|) m_e \frac{T}{E_\nu^2}]$$

– ν mass models all mechanisms carry some modifications to the structure of the standard EW NC & CC

– **V-A Form**, similar to the four Fermi

- exchange of Higgs
- Supersymmetric scalar bosons
- New heavy gauge boson Z'

– characterized by some new couplings, called NSI

(NU) NSI: ϵ_{ee}^{eLR}

(FC) NSI: $\epsilon_{e\mu}^{eLR} \quad \epsilon_{e\tau}^{eLR}$

$$\tilde{g}_L = g_L + \epsilon_{ee}^{eL}$$

$$\tilde{g}_R = g_R + \epsilon_{ee}^{eR}$$

❖ There is a strict bound on $|\epsilon_{e\mu}^{eLR}| < 7.7 \times 10^{-4}$ derived from

$\mu \rightarrow 3e$ decay

❖ The main parameters will be $\epsilon_{e\tau}^{eLR}$ for **FC NSI** and

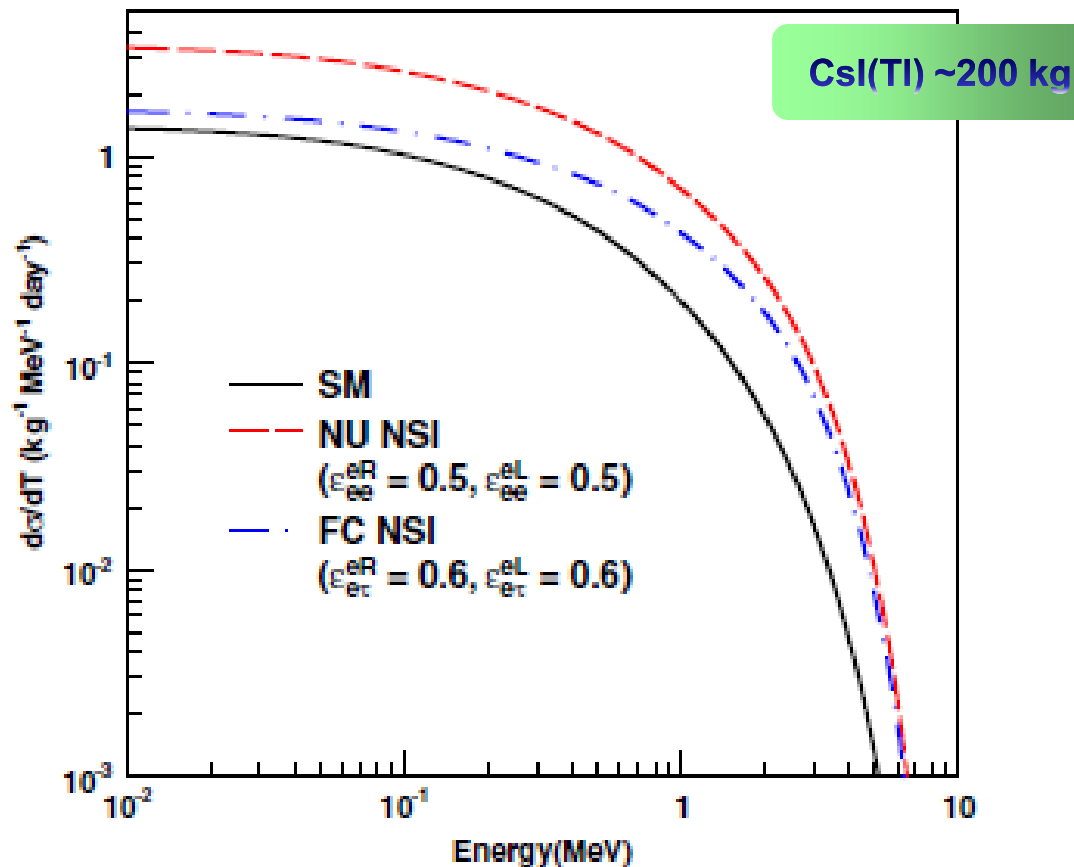
ϵ_{ee}^{eLR} for **NU-NSI**.

NSI of Neutrino

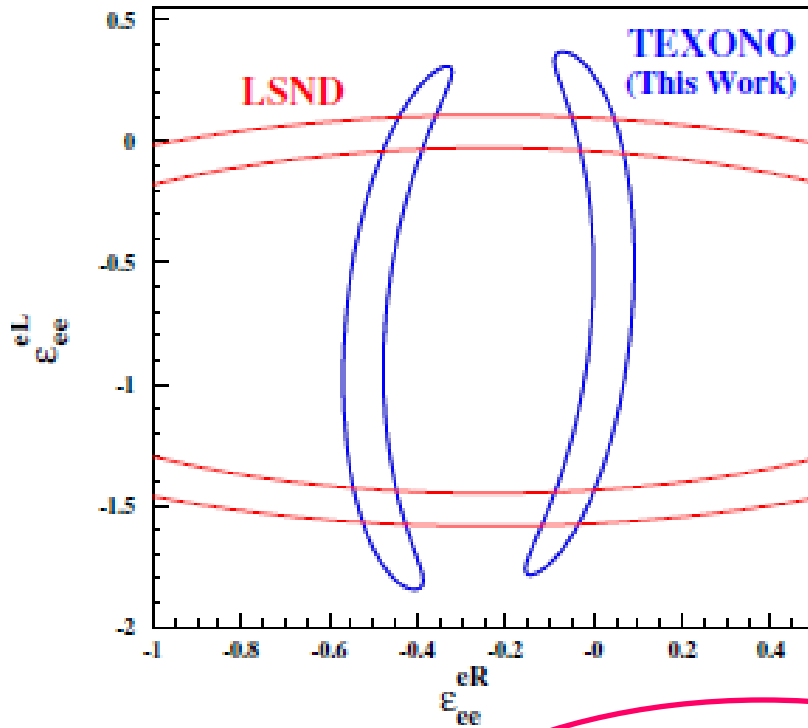
– $\bar{\nu}_e - e^-$ scattering provide a sensitive tool to probe NSI

– The NSI parameters are constrained by the accuracy of the SM cross-section measurements

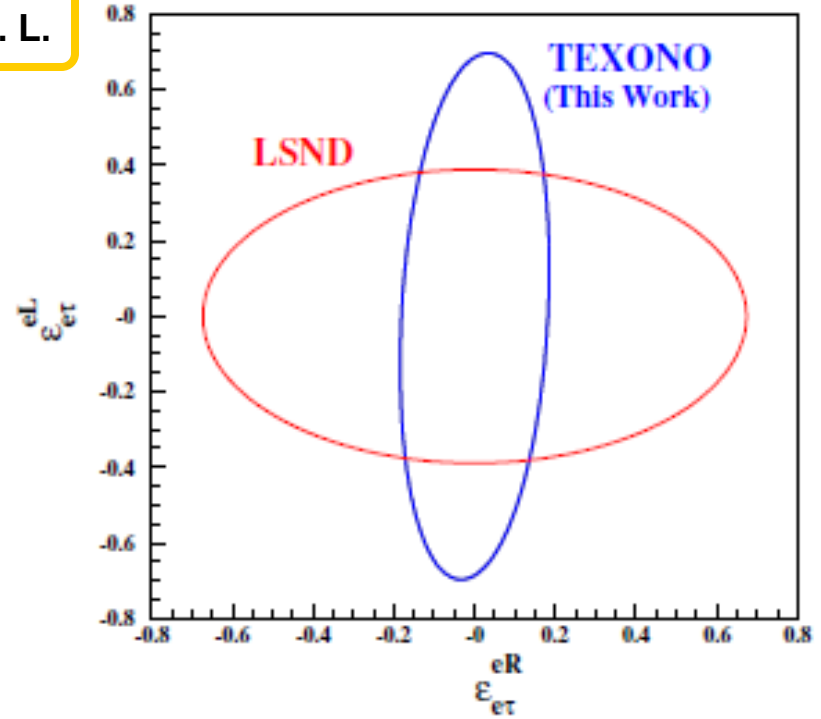
The measurable recoil spectra with typical reactor neutrino “flux”
at typical values of NSI parameters for both NU and FC NSI



Comparison of Bounds of NSI Parameters



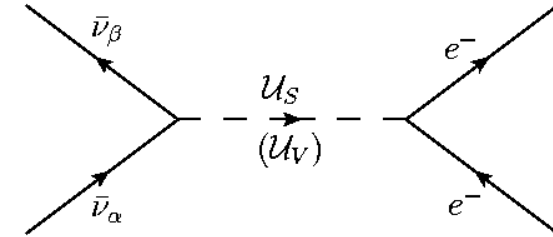
at 90% C. L.



NSI parameters	TEXONO (this work)		Bounds at 90% C.L.	Projected sensitivities	LSND [13]	Combined [14] Bounds at 90% C.L.	Ref. [15]	
	Measurement best fit	$\chi^2/\text{d.o.f.}$						
NU	$\varepsilon_{ee}^{eL} = 0.03 \pm 0.26 \pm 0.17$	8.9/9	$-1.53 < \varepsilon_{ee}^{eL} < 0.38$	± 0.015	$-0.07 < \varepsilon_{ee}^{eL} < 0.11$	$-0.03 < \varepsilon_{ee}^{eL} < 0.08$	$ \varepsilon_{ee}^{eL} < 0.06$	
	$\varepsilon_{ee}^{eR} = 0.02 \pm 0.04 \pm 0.02$	8.7/9	$-0.07 < \varepsilon_{ee}^{eR} < 0.08$	± 0.002	$-1.0 < \varepsilon_{ee}^{eR} < 0.5$	$0.004 < \varepsilon_{ee}^{eR} < 0.151$	$ \varepsilon_{ee}^{eR} < 0.14$	
FC	$\varepsilon_{e\mu}^{eL}$	$\varepsilon_{e\mu}^{eL^2} (\varepsilon_{e\tau}^{eL^2}) = 0.05$	8.9/9	$ \varepsilon_{e\mu}^{eL} < 0.84$	± 0.052	\dots	$ \varepsilon_{e\mu}^{eL} < 0.13$	$ \varepsilon_{e\mu}^{eL} < 0.1$
	$\varepsilon_{e\tau}^{eL}$	$\pm 0.27 \pm 0.24$		$ \varepsilon_{e\tau}^{eL} < 0.84$	± 0.052	$ \varepsilon_{e\tau}^{eL} < 0.4$	$ \varepsilon_{e\tau}^{eL} < 0.33$	$ \varepsilon_{e\tau}^{eL} < 0.4$
	$\varepsilon_{e\mu}^{eR}$	$\varepsilon_{e\mu}^{eR^2} (\varepsilon_{e\tau}^{eR^2}) = 0.008$	8.7/9	$ \varepsilon_{e\mu}^{eR} < 0.19$	± 0.007	\dots	$ \varepsilon_{e\mu}^{eR} < 0.13$	$ \varepsilon_{e\mu}^{eR} < 0.1$
	$\varepsilon_{e\tau}^{eR}$	$\pm 0.015 \pm 0.012$		$ \varepsilon_{e\tau}^{eR} < 0.19$	± 0.007	$ \varepsilon_{e\tau}^{eR} < 0.7$	$0.05 < \varepsilon_{e\tau}^{eR} < 0.28$	$ \varepsilon_{e\tau}^{eR} < 0.27$

Unparticle Physics

- ❖ The notion of unparticles is introduced by Howard Georgi . A scale invariant sector which decouples at a sufficiently large energy scale exists. [*Phys. Rev. Lett.* **98**, 221601 (2007)]
- ❖ The signatures of Unparticles can also be observed by reactor neutrinos by searching the effects of virtual unparticle exchange between fermionic currents.
- ❖ This interaction can be either exchange of **Scalar Unparticles** or **Vector Unparticles**.



– The UP does not have any well defined invariant mass but rather has a continuous mass spectrum

1. Exchange of Scalar Unparticles

$$\frac{d\sigma_{U_S}}{dT} = \frac{[g_{0e}^{\alpha\beta}(d)]^2}{\Lambda^{(4d-4)}} \frac{2^{(2d-6)}}{\pi E_\nu^2} (m_e T)^{(2d-3)} (T + 2m_e)$$

2. Exchange of Vector Unparticles

$$\frac{d\sigma_{U_V}}{dT} = \frac{1}{\pi} \frac{[g_{1e}^{\alpha\beta}(d)]^2}{\Lambda^{(4d-4)}} 2^{(2d-5)} (m_e)^{(2d-3)} (T)^{(2d-4)} \left[1 + \left(1 - \frac{T}{E_\nu}\right)^2 - \frac{m_e T}{E_\nu^2} \right]$$

For the flavour changing case:

$$\frac{d\sigma_{U_V-SM}}{dT} = \frac{\sqrt{2}G_F}{\pi} \frac{g_{1e}(d)}{\Lambda^{(2d-2)}} (2m_e T)^{(d-2)} m_e \left\{ g_L + g_R \left(1 - \frac{T}{E_\nu}\right)^2 - \frac{(g_L + g_R) m_e T}{2 E_\nu^2} \right\}$$

For the flavour conserving case

$i = \mathbf{0(1)}$: Unparticle scalar/vector field

$\lambda_0 (\lambda_1)$: Scalar(Vector) unparticle couplings

\mathbf{f} : e, u, d

α, β : denotes neutrino flavours

d : Unparticle mass dimension

Λ : Unparticle energy scale

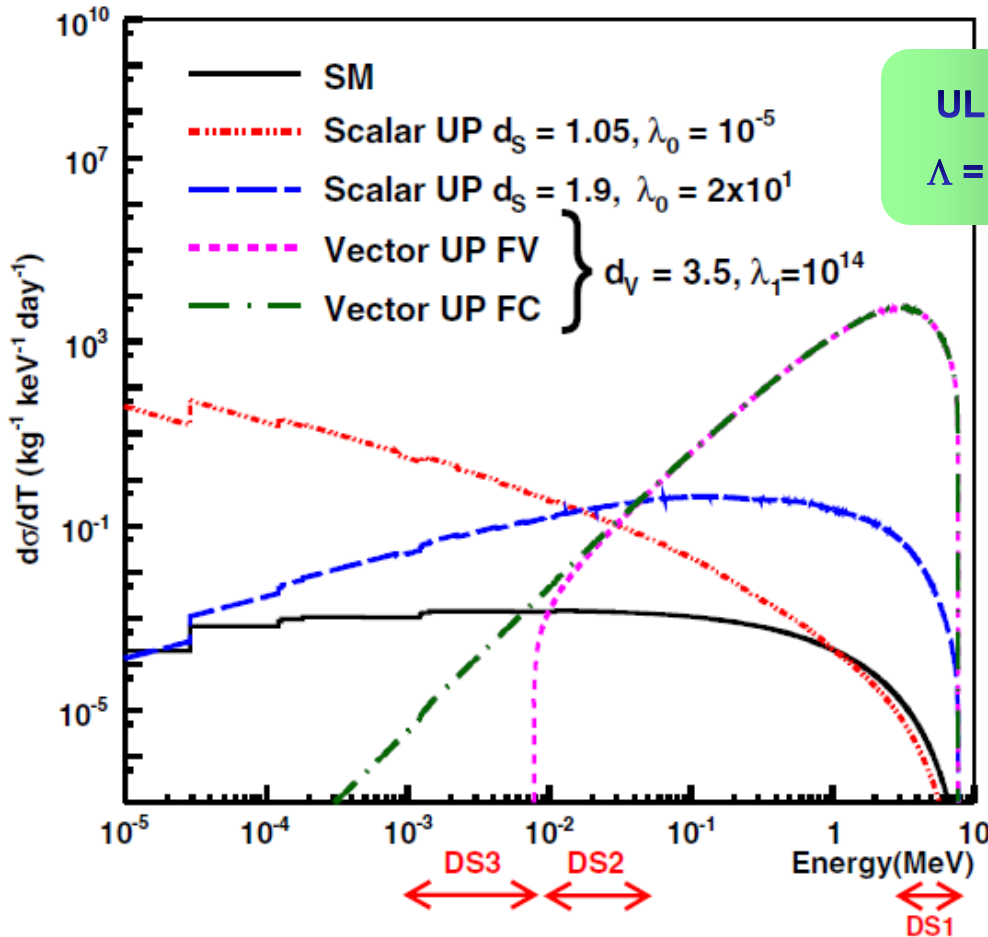
$$g_{if}^{\alpha\beta}(d) = \frac{\lambda_{i\nu}^{\alpha\beta} \lambda_{if}}{2 \sin(d\pi)} A_d \quad A_d = \frac{16\pi^{5/2}}{(2\pi)^{2d}} \frac{\Gamma(d+1/2)}{\Gamma(d-1)\Gamma(2d)}$$

$$\lambda_0 (\lambda_1) = \sqrt{\lambda_{0\nu}^{e\beta} \lambda_{0e}} (\sqrt{\lambda_{1\nu}^{e\beta} \lambda_{1e}})$$

Unparticles



– The differential cross-section of the interaction for $\bar{\nu}_e$ -e scattering via virtual scalar and vector UP exchange



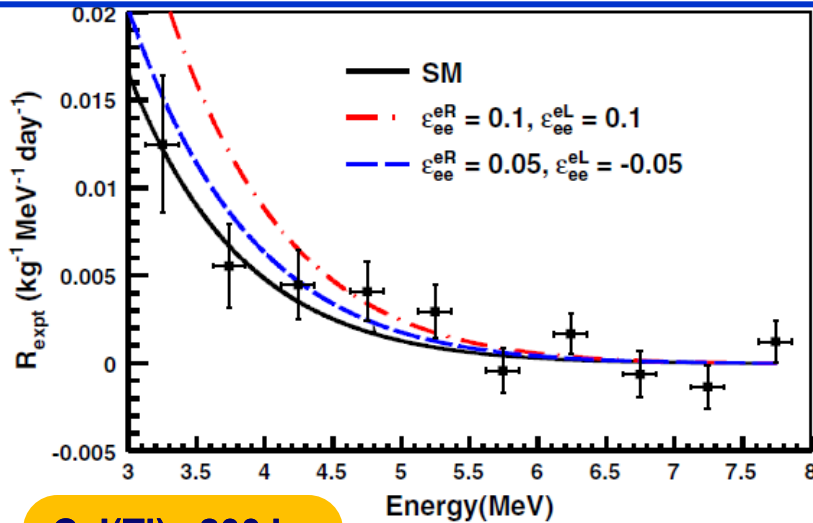
ULB Ge
 $\Lambda = 1 \text{ TeV}$

– measurements with low energy threshold are expected to provide better sensitivities for low values of mass dimension of “d”

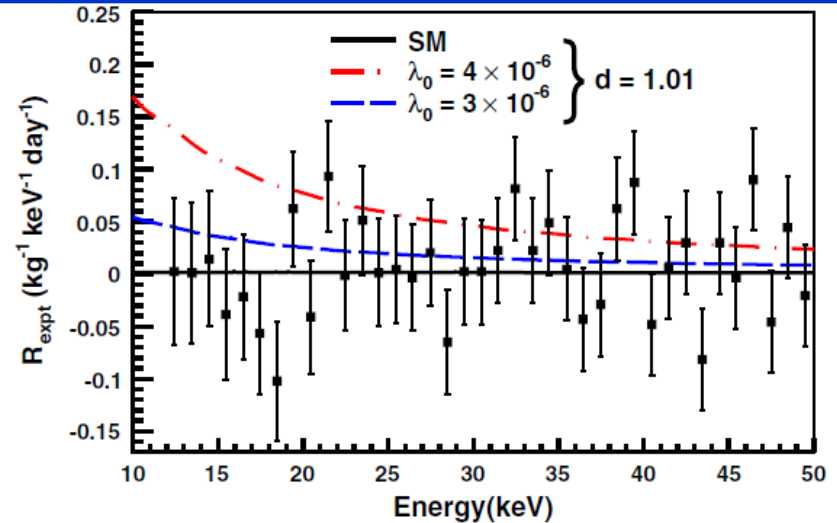
– high energy experiments are preferred to probe UP due to the large values of “d”

Unparticle - Exclusion Plots

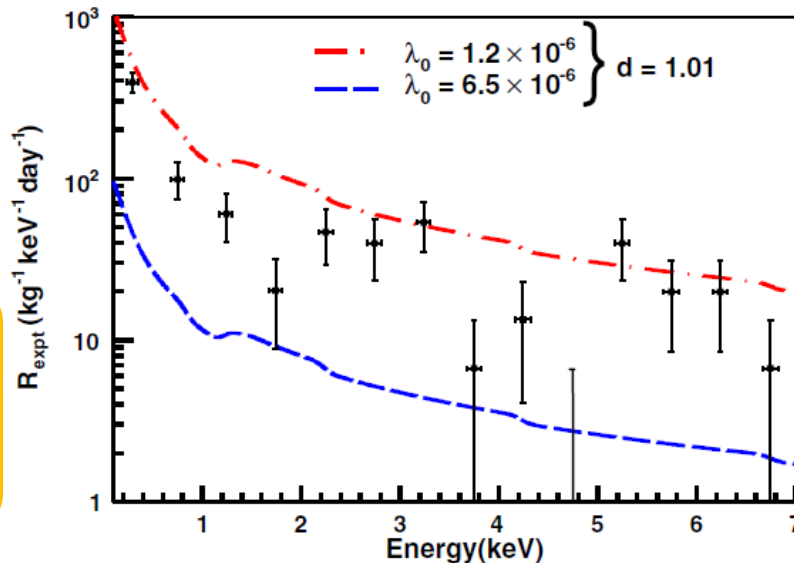
– For the illustrations (only) the measurable spectra of an excluded and allowed parameter space where the data provides the stringent bounds for typical values of parameters



Csl(Tl) ~200 kg
This Talk
PRD 2010



ULB Ge 1 kg
v MM Data Set
PRD 2001, PRD 2007

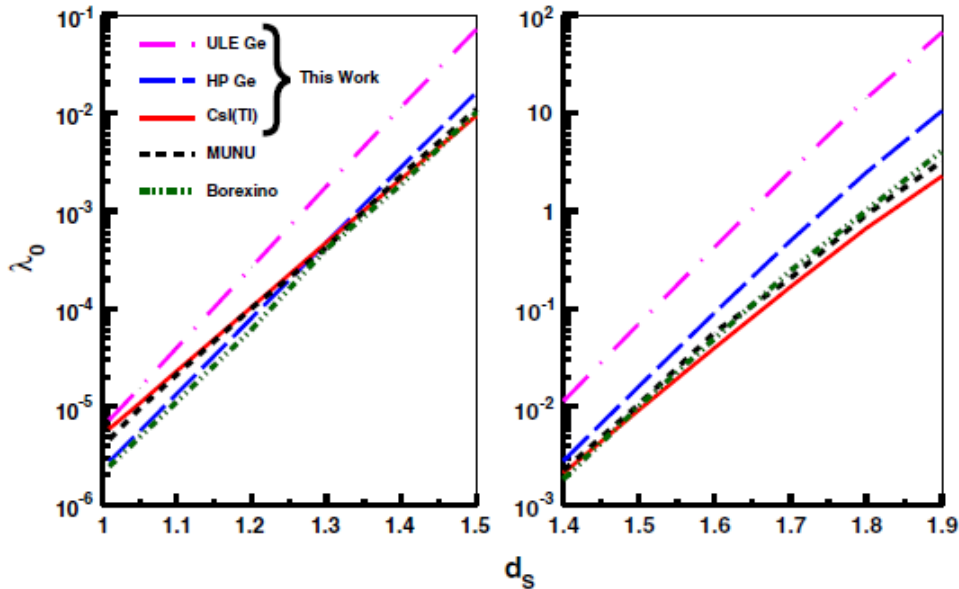


ULE Ge 20 g
WIMP data set
PRD 2009

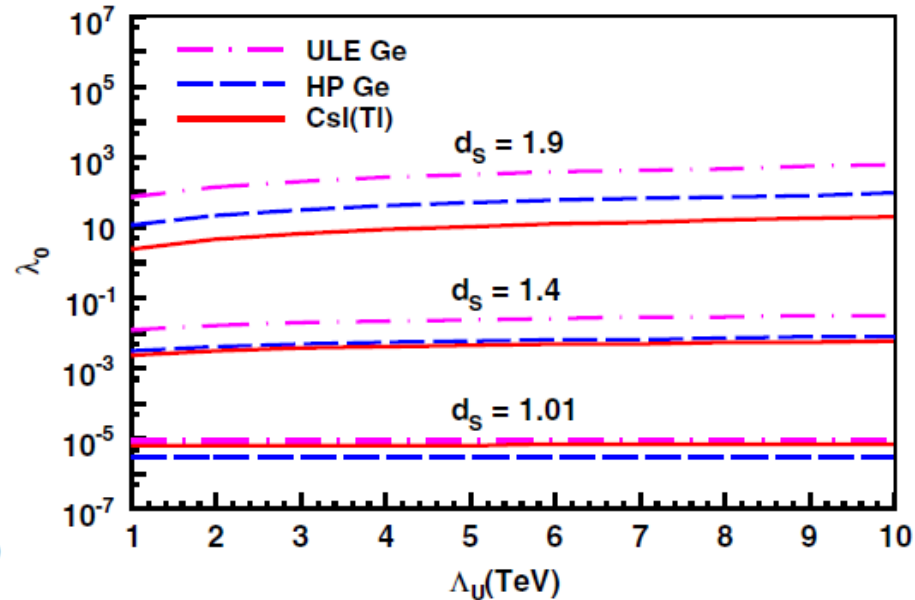
Scalar Unparticle

$$\frac{d\sigma_{US}}{dT} = \frac{[g_{0e}^{\alpha\beta}(d)]^2 2^{(2d-6)}}{\Lambda^{(4d-4)} \pi E_\nu^2} (m_e T)^{(2d-3)} (T + 2m_e)$$

$\Lambda_U = 1 \text{ TeV}$



$\Lambda_U = 1 - 10 \text{ TeV}$



Results are improved over those from Borexino and MUNU experiments

Unparticle effects decreases as Λ_U increases

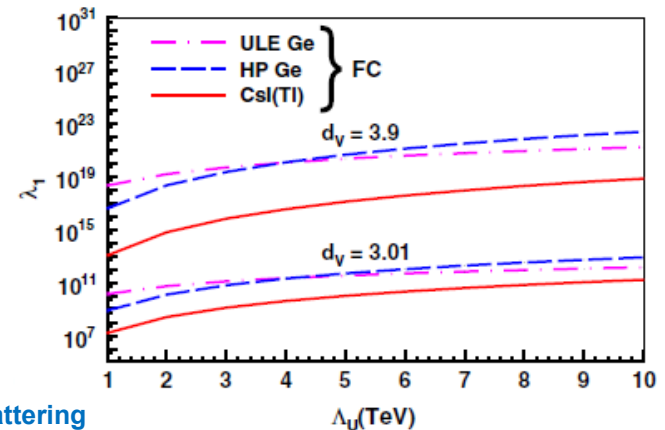
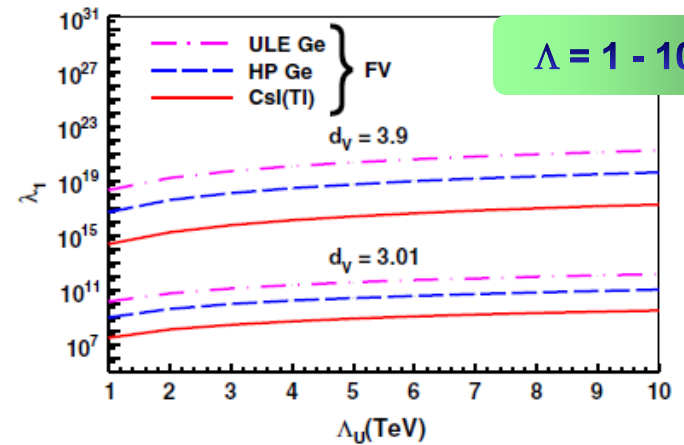
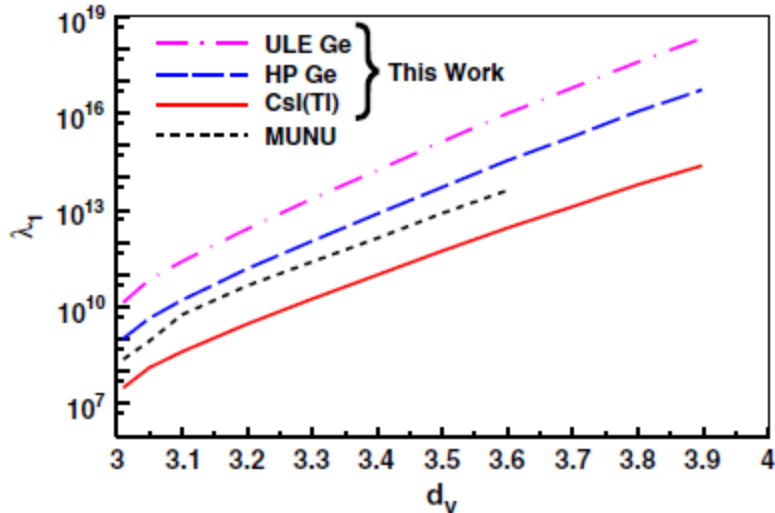
Vector Unparticle

$$\frac{d\sigma_{UV}}{dT} = \frac{1}{\pi} \frac{[g_{1e}^{\alpha\beta}(d)]^2}{\Lambda^{(4d-4)}} 2^{(2d-5)} (m_e)^{(2d-3)} (T)^{(2d-4)} \left[1 + \left(1 - \frac{T}{E_\nu}\right)^2 - \frac{m_e T}{E_\nu^2} \right]$$

$$\frac{d\sigma_{UV-SM}}{dT} = \frac{\sqrt{2}G_F}{\pi} \frac{g_{1e}(d)}{\Lambda^{(2d-2)}} (2m_e T)^{(d-2)} m_e \left\{ g_L + g_R \left(1 - \frac{T}{E_\nu}\right)^2 - \frac{(g_L + g_R) m_e T}{2 E_\nu^2} \right\}$$

– both FC and FV scenario are considered and analysed

$\Lambda = 1 \text{ TeV}$



Summary

■ **Detector:** CsI(Tl) Scintillating Crystal Array (~ 200 kg)

- **Threshold:** 3 MeV

■ **Analysis Results:**

- $\sigma(\bar{\nu}_e - e^-)$ with $\sim 25\%$ accuracy
- **Weak Mixing Angle** with $\sim 15\%$ accuracy
- Verify SM **negative interference**
- $\mu_{\bar{\nu}}$ sensitivity $\sim 10^{-10} \mu_B$
- **neutrino charge radius** sensitivity $\sim 10^{-32} \text{ cm}^2$

■ **Probing new Physics : NSI and UP**

- Current bounds are improved over those from the previous experiments

Thank YOU