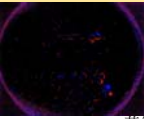



沙石如何“流動” ---- 何謂“流動”?

Rate-dep. transitions from Quicksand, Landslide, to Hourglass



JC Tsai (蔡日強)

with his current research team
 蔡承恩(NCU), 范姜泓杰(NTHU), 黃冠皓
 and numerous contributing members
 黃湘淳, 黃少榆, 黃柏嘉, 周明觀, 費祥霖etc. in the past
Inst. of Physics, Academia Sinica, Taipei, TAIWAN



Special thanks to
 X. Cheng, J. Dijksman, 梁鈞泰, 陳志強, 黃仲仁, 杜其永, 施宏燕
 -- AS-IoP Special Colloquium, 2020/12-

Fluid Mechanics (continuum) vs. Granular flow (particles)

- LAWS of MOTION** in Partial Diff. Eqn.:
 - Euler's equation (est. 1757)
 - for *inviscid* flows, still in use for **air foils**
 - Navier-Stokes equation (est. 1845)
 - as complete theory, with **viscosity** included
- CONSTITUENT RELATIONS**
 - Newtonian vs. **Non-Newtonian “rheology”** (est. 1920+)
 - stress determined **locally** by strain rate

BOTH in debates till now

- limited to special cases
- concerns on **“non-locality”**


relationship to **Solid Mechanics**

* animation by Thierry Dugnolle, CC0, <https://commons.wikimedia.org/w/index.php?curid=18803279>

Scenarios of granular flows

Smooth, “rapid”


Quicksand



1962 film 《Larence in Arabia》
 演員 John Dimech 陷在流沙

Intermittent

Debris flows / Earthquakes





南投縣那坑地區 土石流
 堆積扇 [陳宏宇, 1998, 台灣大學地質科學系]

a block of houses


Smooth, “slow”

Hourglass





南投縣集集鎮 武昌宮
 1999/09/21 01:47




台北市 東星大樓
 (八德路四段與虎林街口)
 200+ km away (Taipei)

Source: 維基百科

Description of “flows” --- common languages (1)

as opposed to pure translation / rotation

Shear Strain



$\gamma \equiv \frac{\Delta x(z)}{\Delta z}$

Strain rate (Shear rate)

$$\dot{\gamma} \equiv \frac{\Delta \dot{x}(z)}{\Delta z} = \frac{\partial V_x(\vec{r})}{\partial z}$$


~ 1/ time

- Stress** \equiv Force transmitted / Cross-sectional area, in average
 - e.g. **Shear stress** (tangential): σ_S
 - Normal stress** (perpendicular): σ_N ~ Pressure
- both as a *tensor* in 3D, strictly speaking.....

Understanding granular flows

Smooth, “rapid”


Quicksand



1962 film 《Larence in Arabia》
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Intermittent


Debris flows / Earthquakes



南投縣那坑地區 土石流
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Smooth, “slow”

Hourglass



Q1: How do we define “FAST” versus “SLOW”?

Q2: How do we define “TIGHT” versus “LOOSE”?

Dimensionless parameters --- common languages (2)

Fluid Mechanics [>200 yr-old]

Long list of examples:

- Reynolds #** $= L V / (\eta / \rho)$, $V \sim \dot{\gamma}$
 - inertial vs. viscous forces
 - turbulent vs. laminar flows
- Peclet #** $= L V / \text{Diffusion constant}$
 - migration vs. thermal motion
- Mach #** $= U / \text{Speed of sound}$
 - importance of compressibility
- Deborah/Weissenburg #**
 - deg. of relaxation (for non-Newtonian fluid)
- Capillary #**
 - viscous force versus surface tension

Granular Flow [in past 20 yrs]

“How fast do particles flow?”

- Inertia number** [est. 2004+]
- Viscous number** [est. 2011+]

$I^2 \equiv \rho d^2 \dot{\gamma}^2 / \sigma_N$

Mass volume

Particle size

$J \equiv \eta \dot{\gamma} / \sigma_N$

Fluid viscosity


Confining pressure

→ ARE THEY SUFFICIENT ?

Understanding granular flows

Smooth, "rapid"


Quicksand



1962 film <Larence in Arabia>
演員 John Dimech 陷在流沙

Intermittent


Debris flows / Earthquakes



南投縣部坑地區 土石流
堆積部 [陳宏宇, 1998,
高橋大學地質科學系]

Smooth, "slow"

Hourglass




Q1: How do we define "FAST" versus "SLOW"?
Q2: How do we define "TIGHT" versus "LOOSE"?

Land-mark concepts & simulations: role of friction X_μ

$\phi_{hcp} = \frac{\pi}{\sqrt{18}} \approx 0.74$

↑ Packing Fraction ϕ

FRICTIONLESS (" $X_\mu = 0$ ")

• "Ideal jamming" = 

Random Close Packing (RCP*) or Max. Random Jamming (MRJ)**

$\phi_c \approx 0.64$

FRICIONAL (" $X_\mu \sim O(1)$ ")

• Packing of realistic particles

But, what if X_μ also depends on driving rate $\dot{\gamma}$

$\phi_c \approx 0.55$
Random Loose Packing (RCP)

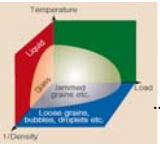
*) Scott & Kilgour (1969)
**) Proposal on MRJ [Torquato et al. 2000-PRL]

Land-mark concepts

--- "jamming" of random, spherical particles

$\phi_{hcp} = \frac{\pi}{\sqrt{18}} \approx 0.74$

↑ Packing Fraction ϕ



Temperature
Density
Liquid
Jammed
Solid

.....[Liu & Nagel, 1998-Nature]

• Jamming diagram (amorphous)

Random Close Packing (RCP*) or Maximally Random Jamming (MRJ)**

$\phi_c \approx 0.64$

*) Scott & Kilgour (1969)
**) Proposal on MRJ [Torquato et al. 2000-PRL]

Short (incomplete) answers:
 Q1: "FAST" vs. "SLOW" → Dimensionless #
 Q2: "TIGHT" vs. "LOOSE" → Dep. on friction

Q3: What about particles immersed in fluid...

(a) Are frictions "eliminated" with things "simplified"?


(b) Are their behaviors covered by physics of suspension [>60 yr-old in *Fluid Mechanics*] ?

Land-mark concepts & simulations: role of friction X_μ

$\phi_{hcp} = \frac{\pi}{\sqrt{18}} \approx 0.74$

↑ Packing Fraction ϕ

FRICTIONLESS (" $X_\mu = 0$ ")

• "Ideal jamming" =  [Liu & Nagel, 1998]

Random Close Packing (RCP*) or Max. Random Jamming (MRJ)**

$\phi_c \approx 0.64$

FRICIONAL (" $X_\mu \sim O(1)$ ")

• Packing of realistic particles

$\phi_c \approx 0.55$
Random Loose Packing (RLP)

MD on isotropic compression of dry particles with varying values of X_μ
.....[Silbert, 2010-SoftMatter]

*) Scott & Kilgour (1969)
**) Proposal on MRJ [Torquato et al. 2000-PRL]

Empirical FACTS versus existing Concepts

$\phi_{hcp} = \frac{\pi}{\sqrt{18}} \approx 0.74$ (HS, ordered)

↑ Solid Volume Fraction ϕ

FRICTIONLESS (" $X_\mu = 0$ ")

Random Close Packing (RCP*)

$\phi_c \approx 0.64$

FRICIONAL (" $X_\mu \sim O(1)$ ")

1) Upper limit observed from stress-controlled shear flow of particles immersed in viscous fluids[Pouliquen, 2011-PRL] $\phi_c \approx 0.58$

$\phi_c \approx 0.55$
Random Loose Packing (RCP)

2) Lower limit known for non-cohesive particles to exhibit a "yield stress".
 $\phi^{(exp.)} \geq 0.48$

Prior experiments by Boyer, Guazzelli, and Pouliquen
[2011-Phys.Rev.Lett.]

--- Cone geometry, fixed # {particles}
--- Stress-controlled
--- Immersed in high-viscosity fluid [$\eta \sim 3000\eta_{water}$] \rightarrow insensitive to properties of particles!?

More Exp. & MDs: reviewed in 2017-JFM known as “*J*-rheology”

$J \equiv \eta \dot{\gamma} / \sigma_N$

fast or small F

Plot (c) shows σ_s / σ_N vs. Solid frac. ϕ . Plot (b) shows σ_s / σ_N vs. Ω for different ϕ .

Rate-dependent behaviors of frictional particles
--- with X_μ up to 0.5

PDMS (polydimethylsiloxane) Particles, $d=0.9\mu m$

Molding Response to uniaxial strain

Instantaneous torque at steady shearing of different Ω

arXiv:2002.07329

Our experiments:
--- a home-built macro-“rheometer”

Side view

Roughened cones, at the scale of particle size $d \sim 1\text{cm}$

--- Controlled at a FIXED TOTAL VOLUME containing N particles
 \rightarrow nominal volume fraction $\phi \equiv \frac{N v_1}{V_{total}}$ ← Volume of single particle

--- Structurally-reinforced for maintaining STEADY ROTATION at a wide range of speeds

effective shear rate $\dot{\gamma} \equiv \frac{\Omega}{2 \tan \zeta}$

** Limited only by the strength of motor and our container

Rough assessment (1/2)
--- definition of σ_s by averaging, and a “flow curve”

Spatial average

Time-averaged values

Instantaneous torque

$\phi = 0.56$

Fluid Viscosity: ∇ 610mP, Δ 80mP, \bullet 10mP

σ_s [kPa]

$\frac{\Omega}{2 \tan \zeta}$ [s⁻¹]

Two sources of information:
force (mean & fluctuations) vs. slice **imaging**

with index-matched interstitial fluid

A) Multiple force sensors (set of 6)

B) Internal fluorescent-contrast images, in-sync with force

Camera

net force(s)

[e.g. dyed PDMS particles]

Rough assessment (2/2): effects of X_μ
---side-by-side comparison, in the same geometry

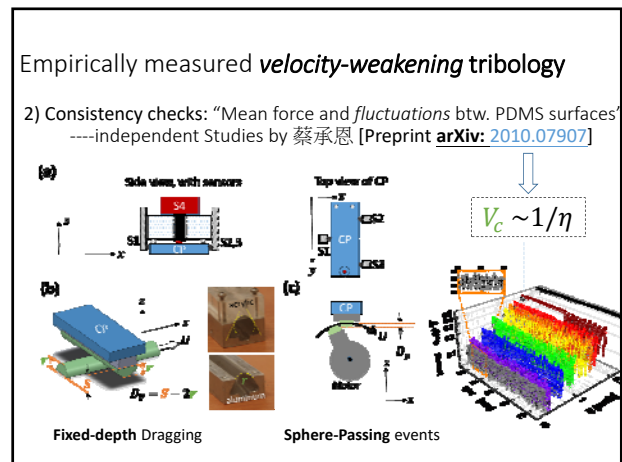
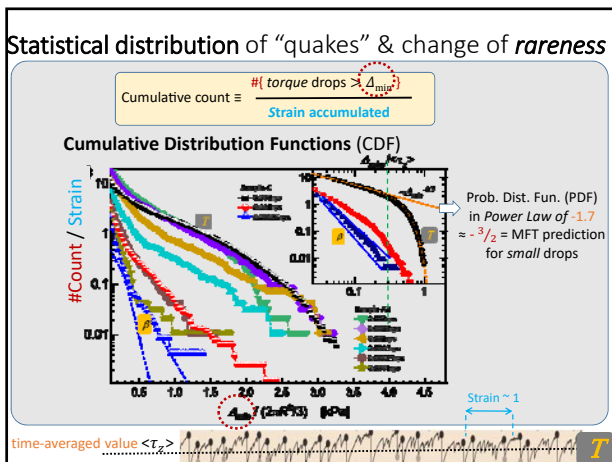
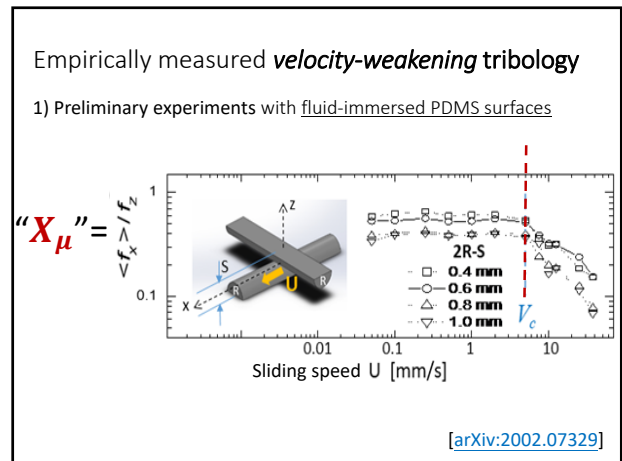
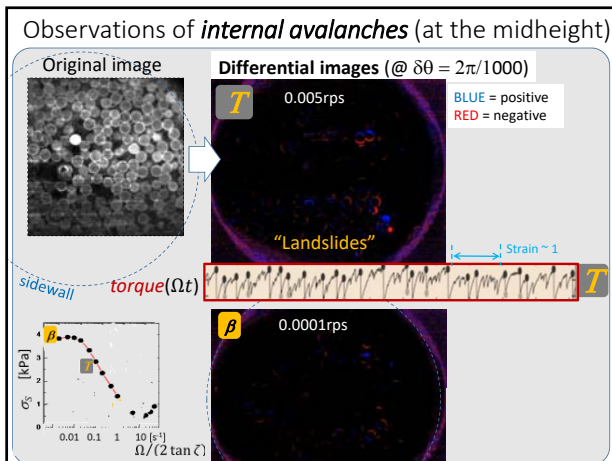
Our prior works on hydrogel balls with $X_\mu < 0.01$ (and a wide range of ϕ , with neutral buoyancy) \rightarrow monotonic increase with Ω

Current work on PDMS spheres with X_μ up to 0.5 (at $\phi = 0.56$ and a similar viscosity) \rightarrow an interval with negative slope

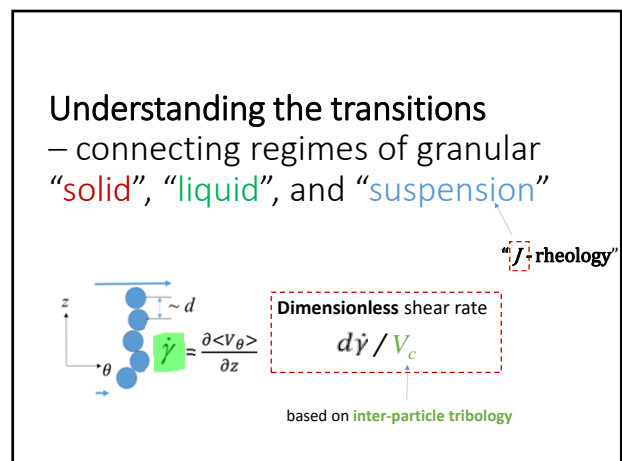
[arXiv:2002.07329]

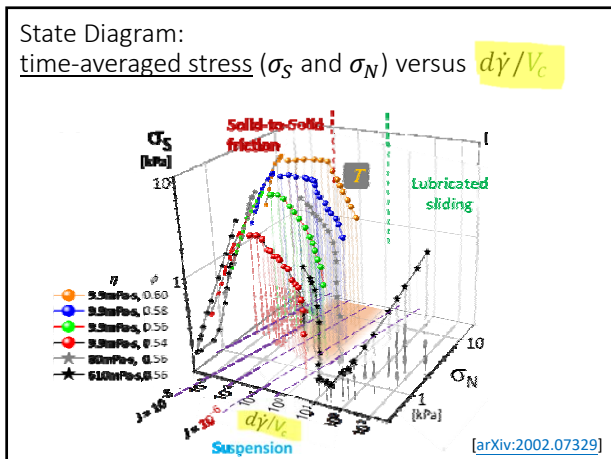
σ_s [Pa]

$\frac{\Omega}{2 \tan \zeta}$ [s⁻¹]



Source of rate dependence:
the velocity-weakening tribology





(a) Numerical experiments with Simple Shear

* Use of well-established LAMMPS (<http://lammps.sandia.gov/>), summing over all contact forces based on a history-dep. Hertzian Mode

Total force on Particle i : $\vec{T}_i^N = \sum_j (T_{ij}^N + T_{ij}^T)$ (Silbert et al. (2001+))

$$\vec{T}_i^N = i\delta_{ij}^2 \left(\frac{R_i R_j}{R_i + R_j} \right)^2 k_c \cdot (\vec{\delta}_{ij} + \dots \vec{\delta}_{ij}^2)$$

$$T_{ij}^T = i\delta_{ij}^2 \left(\frac{R_i R_j}{R_i + R_j} \right)^2 k_c \cdot \Delta \vec{S}_{ij}^T, \text{ if less than } \mu_c |\vec{T}_{ij}^N|;$$

$$X_{\mu} \cdot |\vec{T}_{ij}^N| \frac{1}{|\vec{T}_{ij}^N|}, \text{ otherwise}$$

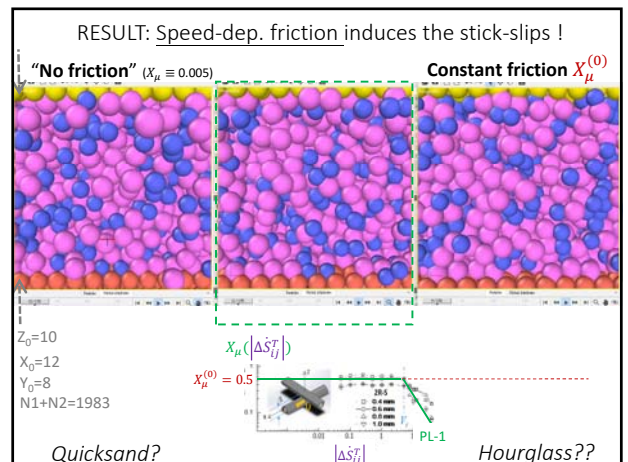
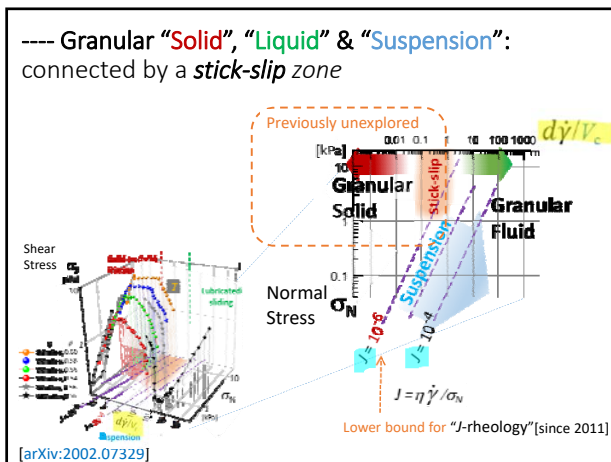
Total torque on Particle i : $\sum_j R_i \frac{r_{ij}}{r_{ij}} \times \vec{T}_{ij}^T$

$\vec{\delta}_{ij} \equiv \frac{\vec{r}_{ij}^R}{|\vec{r}_{ij}^R|} (|\vec{r}_{ij}^R| - (R_i + R_j))$ instantaneous overlap

$\Delta \vec{S}_{ij}^T \equiv$ integrated displacement since the start of contact

$\Delta S_{ij}^T \equiv$ Component of $\Delta \vec{S}_{ij}^T$ perpendicular to $\vec{\delta}_{ij}$

** Our amendment: making its Coulomb coefficient X_{μ} dependent on the sliding speed $|\Delta \dot{S}_{ij}^T|$, with parameters determined by lab measurements.



NEXT:

- Towards a full understanding
 - (a) Numerical simulations [with 范姜泓杰 @NTHU]
 - (b) New experimental setup
- Beyond granular flows

(b) Innovated setup with Circular Shearing
---- in construction

$\vec{s} \equiv s_x + i s_y \equiv e^{i\Omega t}$

Previous work [Denisov et al. 2016]

$R(X, t) \equiv \frac{z}{H} \vec{s}(t)$

$\vec{\gamma} \equiv \frac{\partial R}{\partial X} = \frac{s \Omega}{H} (\hat{z} \times \hat{s}) \otimes \hat{z}$

Collaborations with theorists [Prof. 洪在明 @NTHU]

[arXiv:2010.08739] [arXiv:2010.08289]*

Other collaborations
 --with Dr. 阮文滔 @CMU [Cell, 2019], Patent pending
 --with Prof. 楊馥菱 @NTU-ME

Take-home messages

- (1) Inter-particle tribology: essential ingredient for changes behind **Hourglass** – **Intermittent flows** – **Quicksand**
 [solid-to-solid] [mixed dyn.] [fully-lubricated]
 with a new “dimensionless shear rate” identified, while
- (2) the “loose vs. tight” for packed grains is *coupled* to the driving rate, NOT being determined by just the solid volume fraction alone.

- Force fluctuations & images: complimentary in probing the *dynamics*

→ We need further basis to move on dialogues between communities of Fluid vs. Solid Mechanics.

Recap: What have we learned for flows of packed grains?

Summary, Outlooks & Thanks to AS-IoP

Hydrogel spheres: Suspension with a yield stress determined solely by packing fraction ϕ [Soft Matter, 2020]

FRictional particles (PDMS) in “steady” shearing exp. ----- fluctuation & switches [arXiv:2010.08289]*

Tribology between lubricated PDMS surfaces/particles [arXiv:2010.07907]

Rate-dependent State Diagram $\sigma_S - \sigma_N - [\dot{\gamma}]$

Connection with “J- Rheology” for dense suspensions [pending, with PRL editors; available at arXiv:2002.07329]

Molecular Dynamics Simulations (LAMMPS) on Simple Shearing

Innovated setup (in construction) on Circular Shearing

On-going & Future plans: Collaborations with NTHU & new projects

new Dimensionless Parameter $d\dot{\gamma}/v_c$ based on tribology

in parallel $\dot{\gamma} \equiv \frac{\partial R}{\partial X}$
 $\frac{s \Omega}{H} (2 \times S) \Omega$

Studies on densely-packed granular particles

Hydrogel spheres: Suspension with a yield stress determined solely by packing fraction ϕ [Soft Matter, 2020]

FRictional particles (PDMS) in “steady” shearing exp. ----- fluctuation & switches [arXiv:2010.08289]*

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Rate-dependent State Diagram $\sigma_S - \sigma_N - [\dot{\gamma}]$

Connection with “J- Rheology” for dense suspensions [pending, with PRL editors; available at arXiv:2002.07329]

Previously unexplored

new Dimensionless # for “fast vs. slow” $d\dot{\gamma}/v_c$

Innovated setup (in construction) on Circular Shearing