

Computational Material Chemistry

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Aim

- Understand the basic theory behind quantum chemistry calculation
- Learn running quantum chemistry program
- Understand what the output is saying
- Get a feeling of what method to use for what problem

Grading

- Oral Presentation (10-15 minutes) by students in class: 50% (of which class evaluation is half);
- Report on the calculation that they performed and compare with experiment if available 25%
- Quiz/Homework :25%

Schedule

1. Born Oppenheimer Approximation, LCAO
 H_2^+ calculation
2. H_2 , homonuclear and heteronuclear
diatomic molecule, Restricted Hartree
Fock.
3. Unrestricted Hartree Fock, Roothan
Equation, Basis Set Gaussian Calculation
Input/structure optimization.
4. Potential Energy Surface, Barrier
Transition State.

Schedule

5. Electron correlation (DFT, MP2, QCISD, CCSD, CASSCF, MRCI, G2 G3)
6. Vibrational Spectroscopy, Electronic Spectroscopy
7. Force field parametrization, intermolecular potential
8. Molecular dynamics simulation of liquid

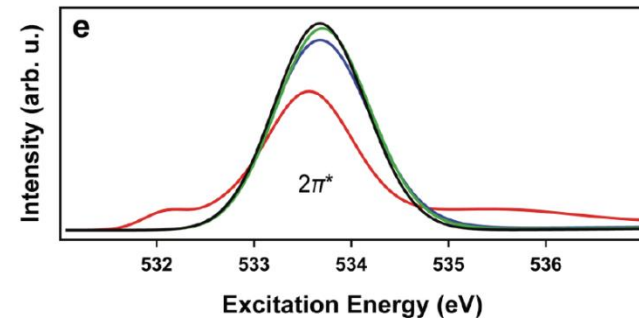
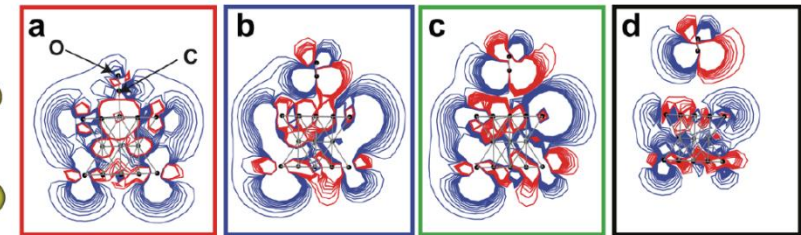
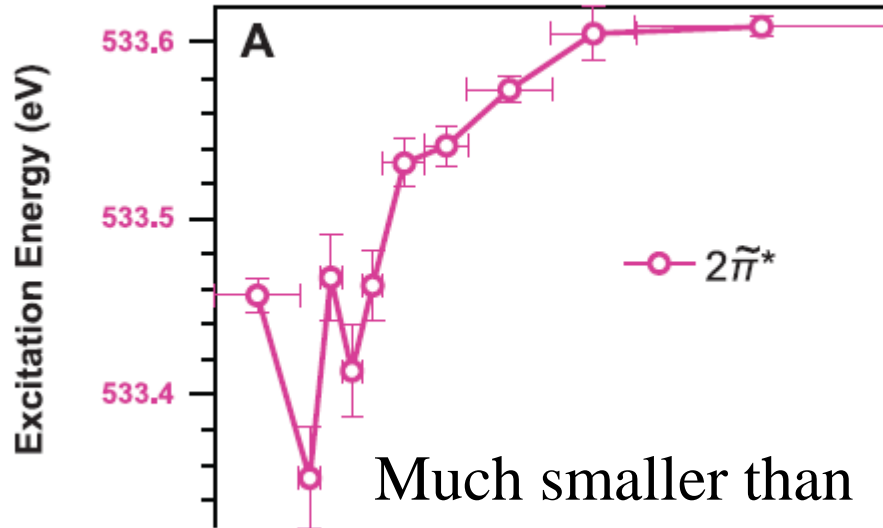
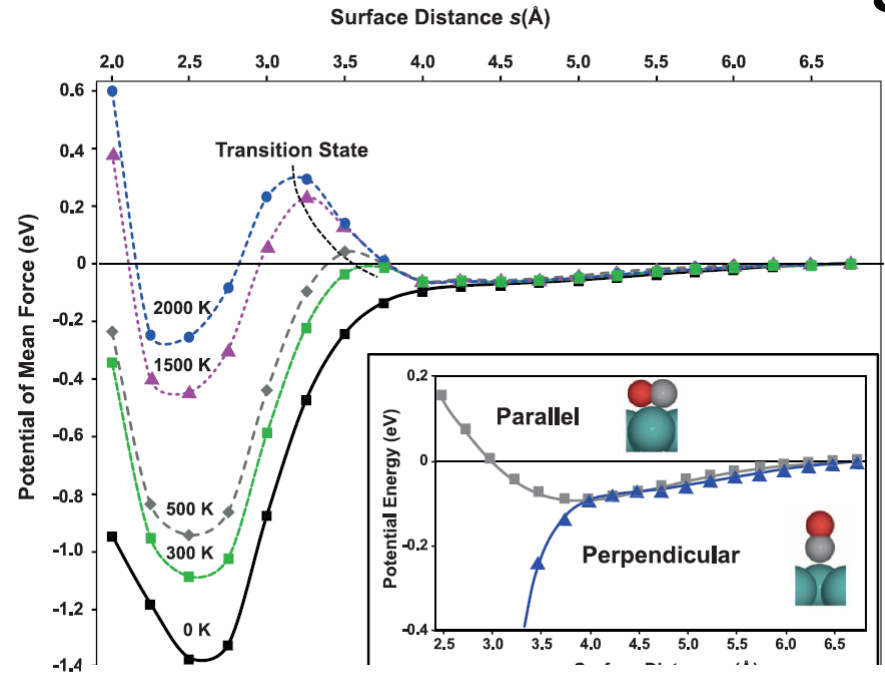
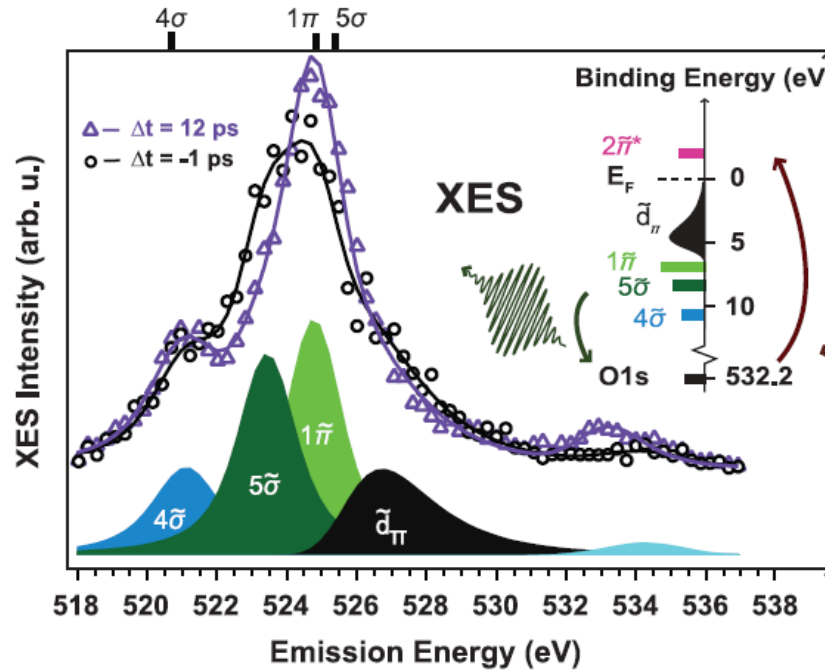
Questions?

- What grade are you in, and what lab are you in.
- Why did you take this class.

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What can we calculate with
quantum chemistry packages?

Understanding Metal Surface Binding



Oxygen Reduction Reaction (ORR) in Hydrogen Fuel Cells

Anode: $\text{H}_2 = 2(\text{H}^+ + \text{e}^-)$

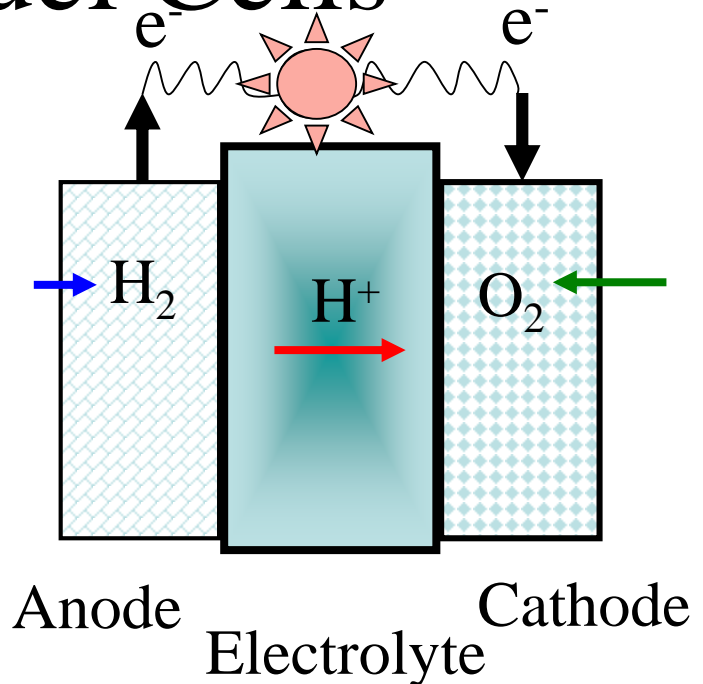
Cathode: $2(\text{H}^+ + \text{e}^-) + \frac{1}{2} \text{O}_2 = \text{H}_2\text{O}$

★ Slow kinetics of the ORR
occurring at the cathode.

Need a better metal than Pt

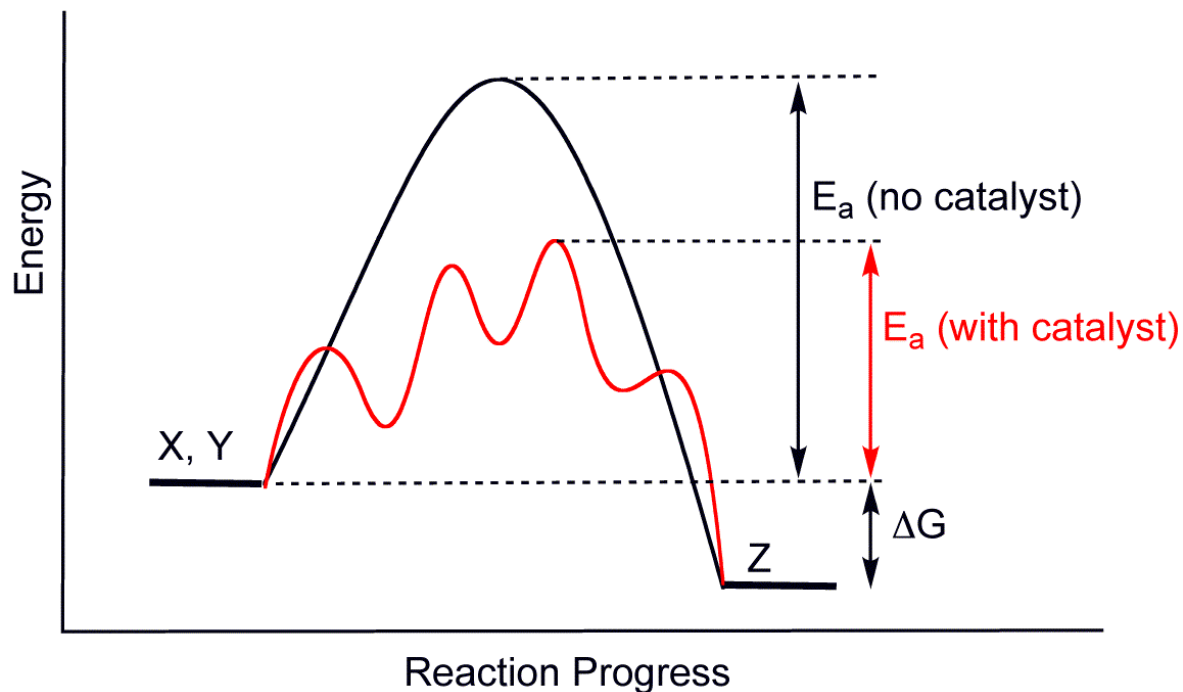
Is alloying the answer?

★ Requirements: (a) low cost, (b) high efficiency, and (c) durability.



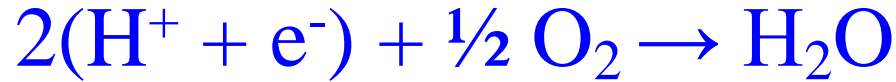
Why metal surfaces?

- Metal surfaces function as catalysts



- Sabatier principle: optimum catalyst must have intermediate affinity for reactants

Oxygen Reduction Reaction (ORR)



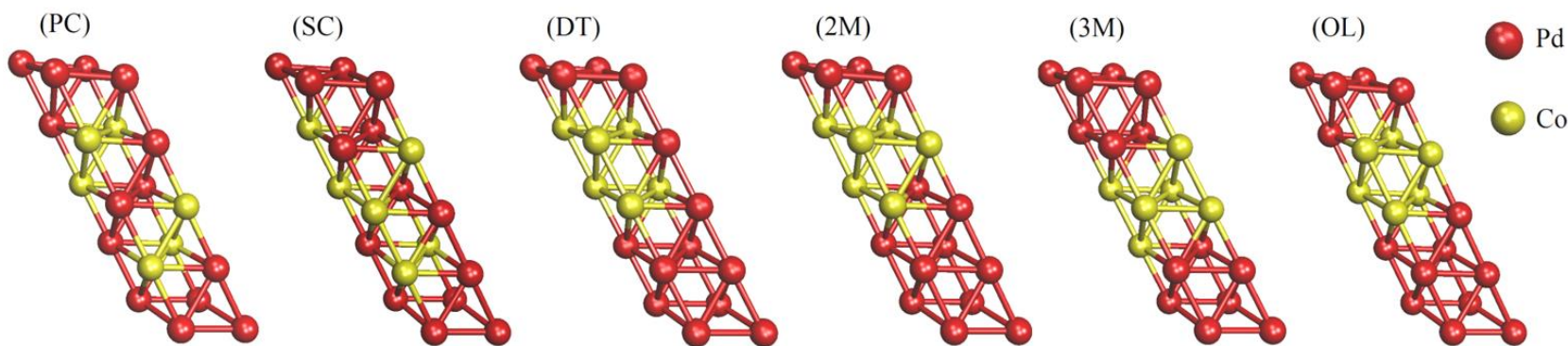
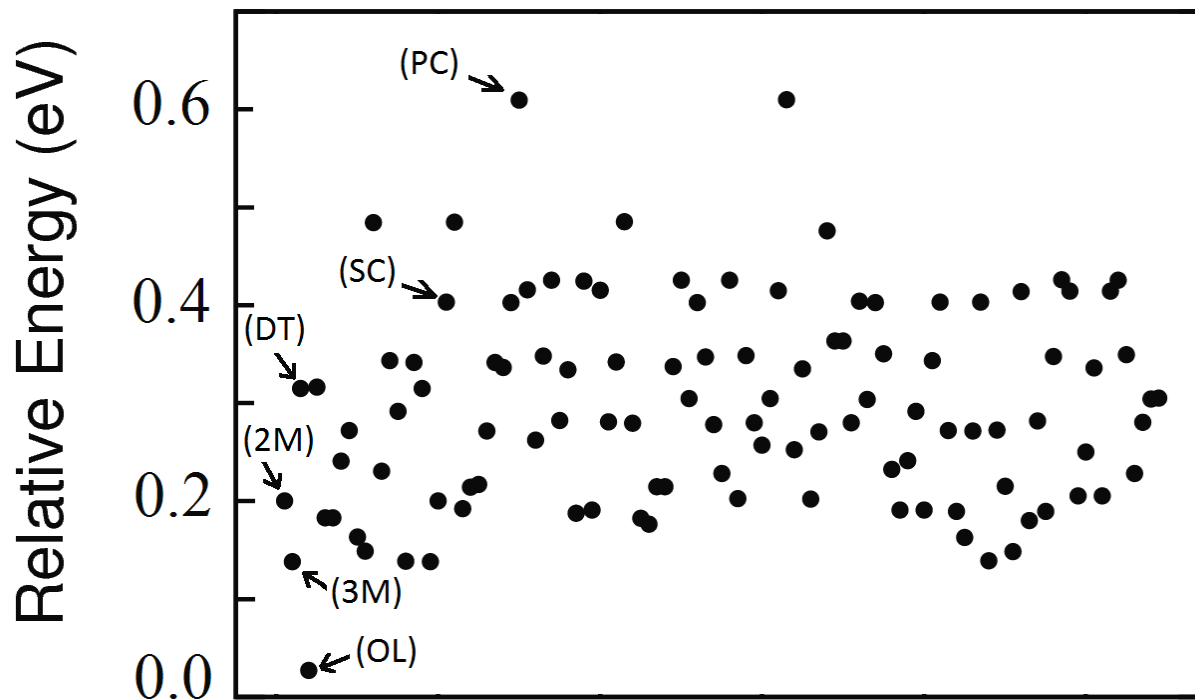
•Mechanism (Intermediates, Rate Determining Steps),
Surf. Sci. 602 (2008) L89-L94, On pure metals (Pt):

Mechanism 1		Mechanism 2		Mechanism 3	
$\text{O}_2 + * \rightarrow \text{O}_2^*$	[1]	$\text{O}_2 + * \rightarrow \text{O}_2^*$	[1]	$\text{O}_2 + * \rightarrow \text{O}_2^*$	[1]
$\text{O}_2^* + * \rightarrow 2\text{O}^*$	[2]	$\text{O}_2^* + (\text{H}^+ + \text{e}^-) \rightarrow \text{OOH}^*$	[3]	$\text{O}_2^* + (\text{H}^+ + \text{e}^-) \rightarrow \text{OOH}^*$	[3]
$\text{O}^* + (\text{H}^+ + \text{e}^-) \rightarrow \text{OH}^*$	[7]	$\text{OOH}^* + * \rightarrow \text{O}^* + \text{OH}^*$	[5]	$\text{OOH}^* + (\text{H}^+ + \text{e}^-) \rightarrow \text{HOOH}^*$	[4]
$\text{OH}^* + (\text{H}^+ + \text{e}^-) \rightarrow \text{H}_2\text{O} + *$	[8]	$\text{O}^* + (\text{H}^+ + \text{e}^-) \rightarrow \text{OH}^*$	[7]	$\text{HOOH}^* + * \rightarrow 2\text{OH}^*$	[6]
		$\text{OH}^* + (\text{H}^+ + \text{e}^-) \rightarrow \text{H}_2\text{O} + *$	[8]	$\text{OH}^* + (\text{H}^+ + \text{e}^-) \rightarrow \text{H}_2\text{O} + *$	[8]

*: binding to surface

Depending on the condition different mechanisms dominate

Co 30% Example

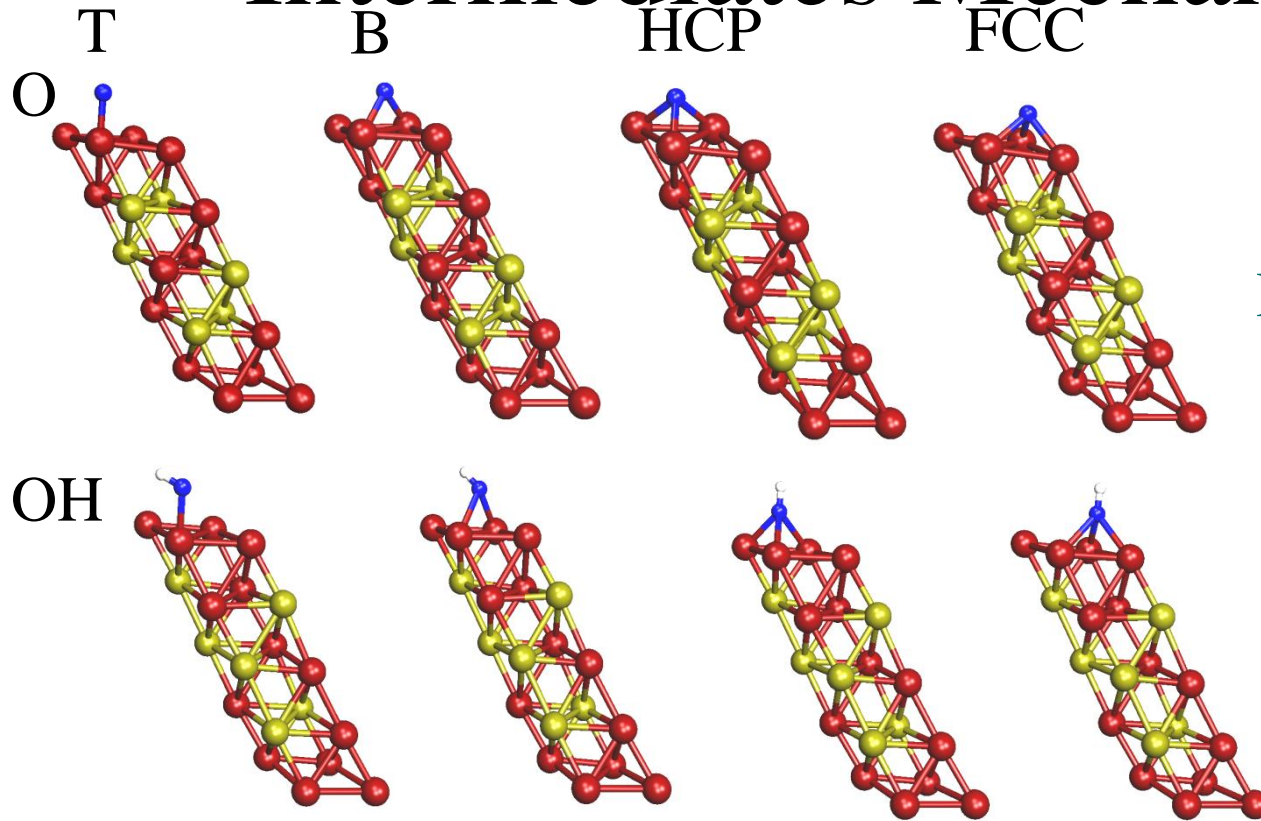


PC: Partially clustering
SC: Scattered

DT: Double-triangle
2M: 2nd-layer monolayer

3M: 3rd-layer monolayer
OL: Octahedral-like 12

Intermediates Mechanism 1



Top: T

Bridge: B

HCP Hollow: HCP

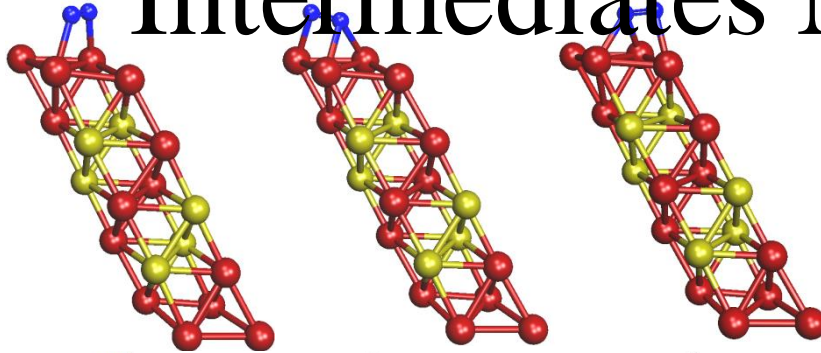
FCC Hollow: FCC

There are two
intermediates:
O and OH.

Intermediates Mechanism 2

OHO vibrational frequency:
2254, 984, 967, 567, 416, 234,
146, 40, 5 (cm^{-1}) all Positive.

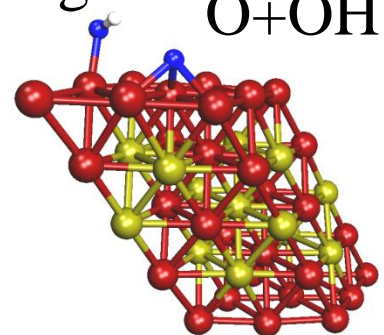
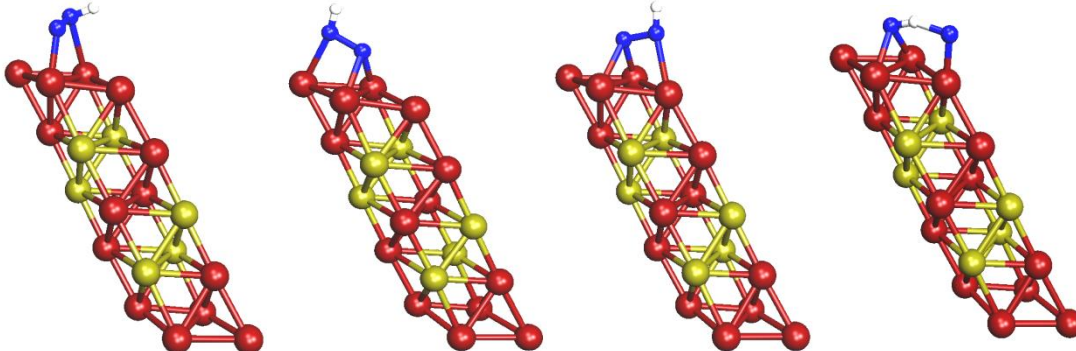
O_2



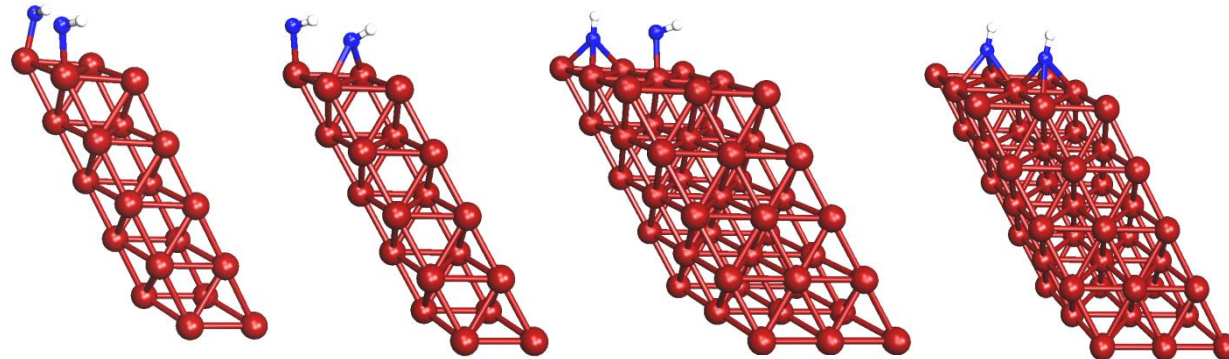
OHO: New Finding

$\text{O}+\text{OH}$

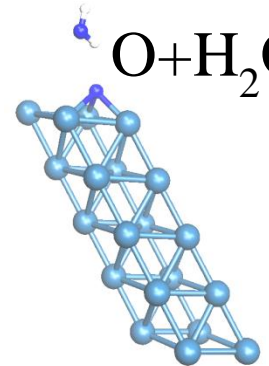
OOH



2OH

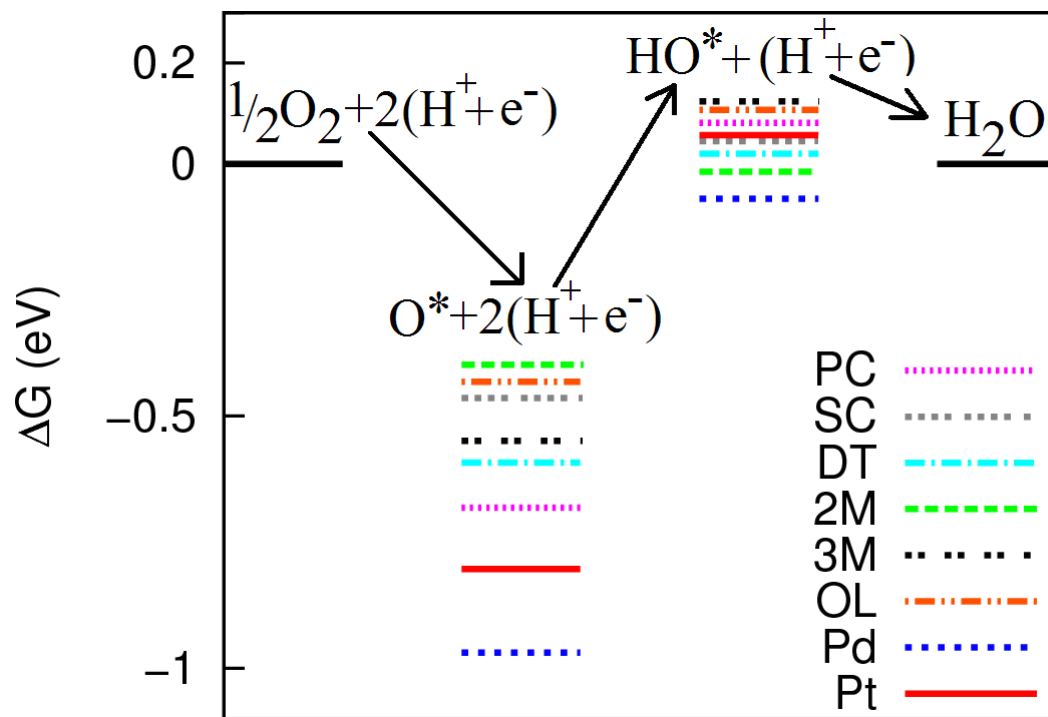


$\text{O}+\text{H}_2\text{O}$



Free Energy Diagram path 1

At equilibrium potential $U = 1.23$ (V), $T = 300$ (K), pressure $P = 1$ (bar), $\text{pH} = 0$. Free energy ΔG calculations take into account reaction energy, zero point energy, and entropy change from DFT results (Computational Methods).

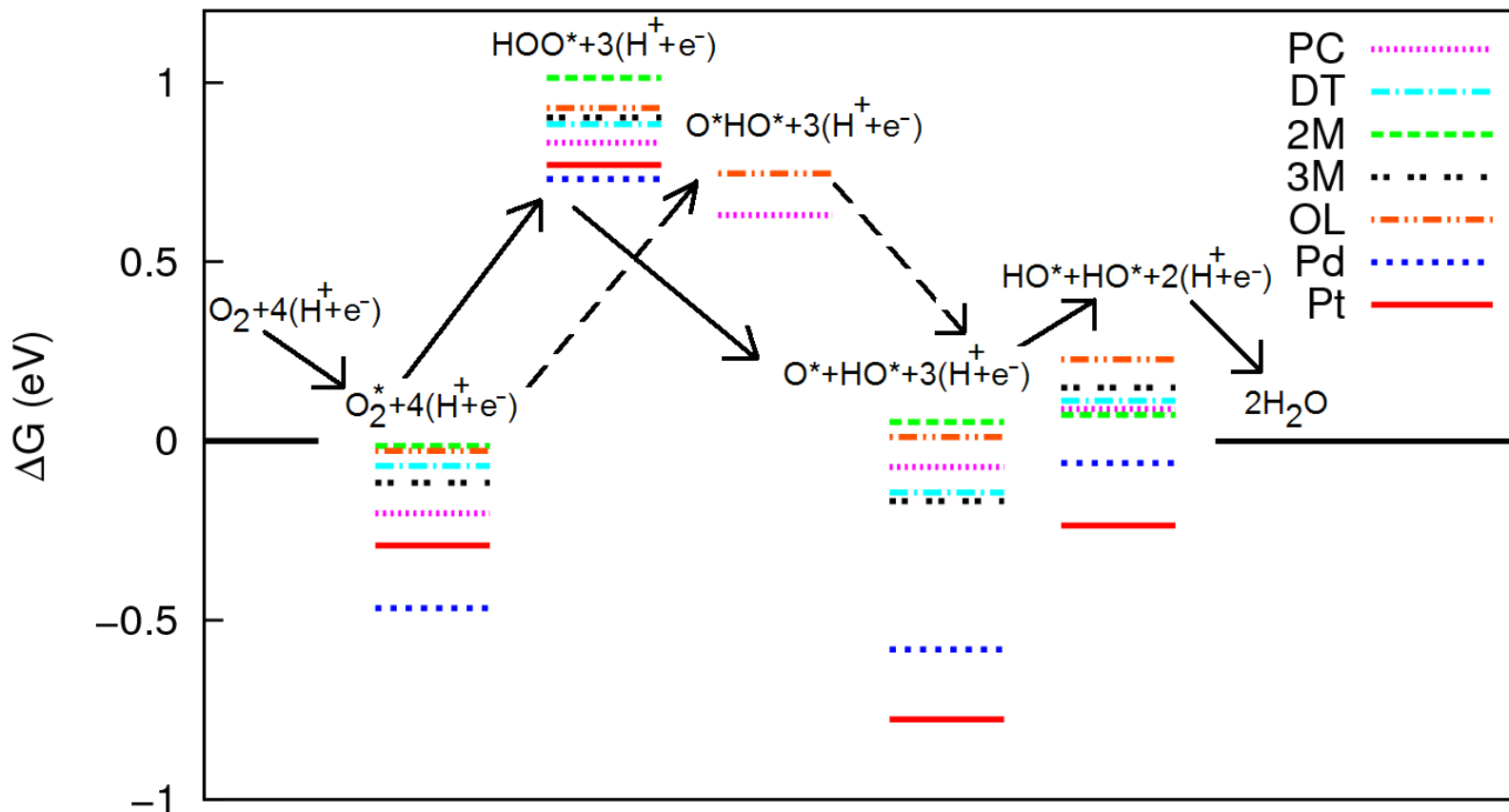


Path 1

Highest thermodynamic barrier locates in the first hydrogenation.

Reference for Pt and Pd:
J. Chem. Theory Comput. 2
 (2006) 1388.

Free Energy Diagram path 2,3



Reference for Pt, Pd: *Surf. Sci.* 602 (2008) L89-L94.

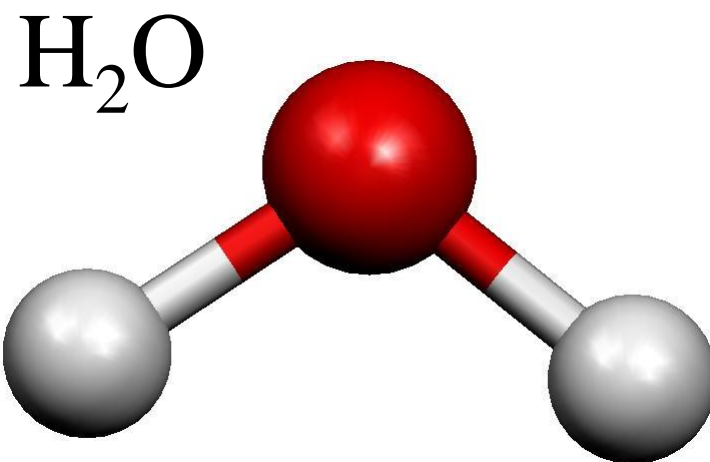
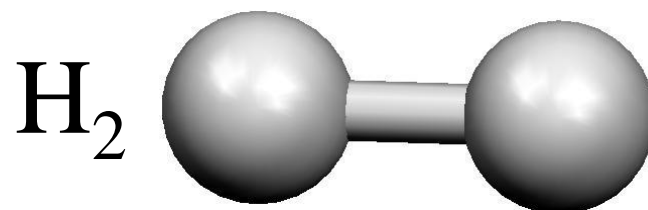
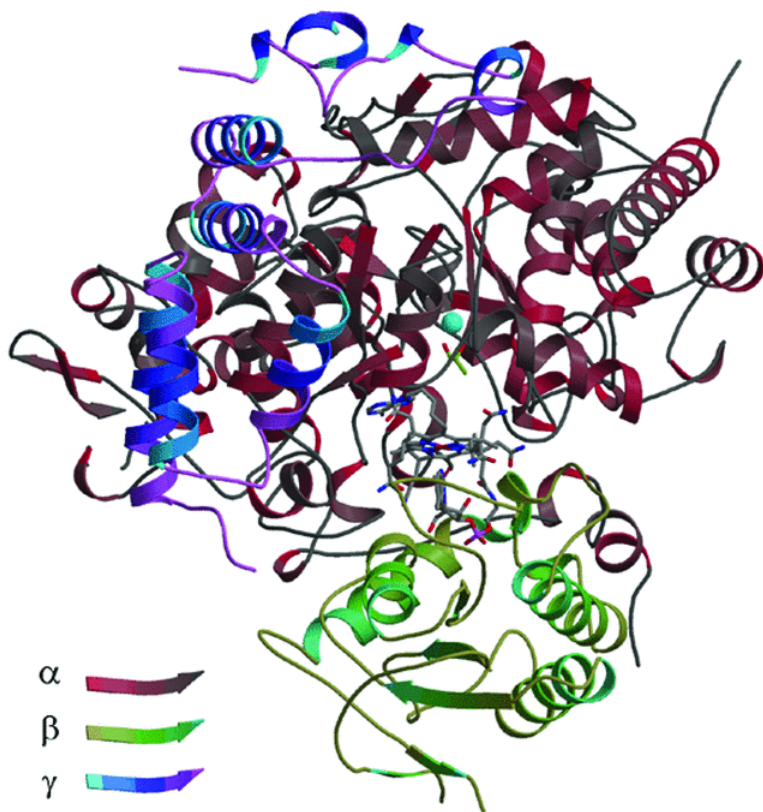
Highest thermodynamic barrier also locates in the first hydrogenation.

Born-Oppenheimer Approximation

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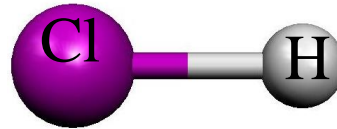
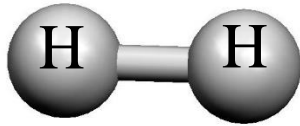
Proteins, Molecules



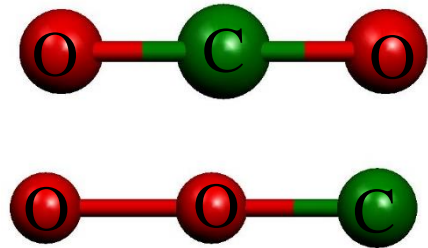
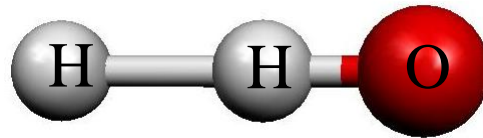
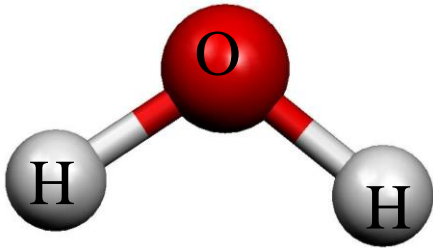
You always write where the nucleus is but you never write the electrons or the electrons are written as a line!!!
YOU ARE ALREADY ASSUMING BORN-OPPENHEIMER APPROXIMATION

Atoms vs Molecules

- Simple case : Diatomic Molecules
 - H_2 , O_2 , F_2 , HCl



- More Complex: Triatomic Molecules
 - H_2O , CO_2



Why graphene like structure is seen for carbon but not nitrogen?

Which structure is the most likely? 19

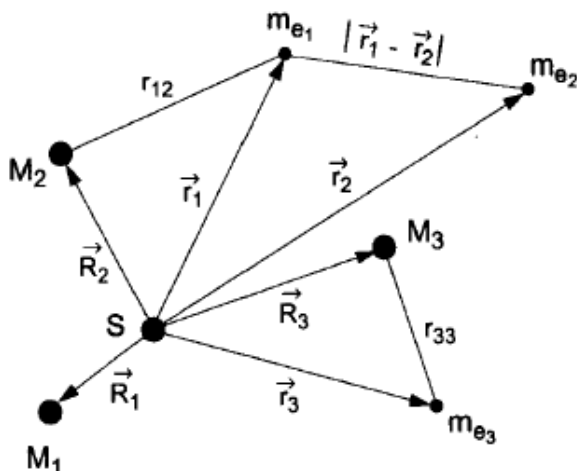
Full Problem

$$\hat{H} = \hat{T} + \hat{V} = -\frac{\hbar^2}{2m_e} \sum_{i=1}^n \nabla_i^2 - \frac{\hbar^2}{2} \sum_{I=1}^N \frac{1}{M_I} \nabla_I^2 + V(\mathbf{r}, \mathbf{R})$$

$$\nabla_i^2 = \frac{\partial^2}{\partial x_i^2} + \frac{\partial^2}{\partial y_i^2} + \frac{\partial^2}{\partial z_i^2}$$

$$\nabla_I^2 = \frac{\partial^2}{\partial X_I^2} + \frac{\partial^2}{\partial Y_I^2} + \frac{\partial^2}{\partial Z_I^2}$$

$$V(\mathbf{r}, \mathbf{R}) = \frac{e^2}{4\pi\epsilon_0} \left[\sum_{I=1}^N \sum_{J \neq I}^N \frac{Z_I Z_J}{|\mathbf{R}_I - \mathbf{R}_J|} - \sum_{I=1}^N \sum_{i=1}^n \frac{Z_I}{|\mathbf{R}_I - \mathbf{r}_i|} + \sum_{i=1}^n \sum_{j \neq i}^n \frac{1}{|\mathbf{r}_i - \mathbf{r}_j|} \right]$$



$$\mathbf{R}_I = \vec{R}_I = X_I \vec{e}_x + Y_I \vec{e}_y + Z_I \vec{e}_z$$

$$\mathbf{r}_i = \vec{r}_i = x_i \vec{e}_x + y_i \vec{e}_y + z_i \vec{e}_z$$

$$\hat{H}\Psi(\mathbf{r}, \mathbf{R}) = E_{el,NU} \Psi(\mathbf{r}, \mathbf{R})$$

Atomic Units

For quantum systems such as electrons and molecules it is easier to use units that fit them=**ATOMIC UNIT**

Use mass of electron (not kg)

Use charge of electron (not coulomb)

Use \hbar for angular momentum (not $\text{kg m}^2 \text{s}^{-1}$)

Use $4\pi\epsilon_0$ for permittivity (not $\text{C}^2 \text{s}^2 \text{kg}^{-1} \text{m}^{-3}$)

TABLE 9.1

Atomic Units and Their SI Equivalents

Property	Atomic unit	SI equivalent
Mass	Mass of an electron, m_e	$9.1094 \times 10^{-31} \text{ kg}$
Charge	Charge on a proton, e	$1.6022 \times 10^{-19} \text{ C}$
Angular momentum	Planck constant divided by 2π , \hbar	$1.0546 \times 10^{-34} \text{ J}\cdot\text{s}$
Length	Bohr radius, $a_0 = \frac{4\pi\epsilon_0\hbar^2}{m_e e^2}$	$5.2918 \times 10^{-11} \text{ m}$
Energy	$\frac{m_e e^4}{16\pi^2\epsilon_0^2\hbar^2} = \frac{e^2}{4\pi\epsilon_0 a_0} = E_h$	$4.3597 \times 10^{-18} \text{ J}$
Permittivity	$\kappa_0 = 4\pi\epsilon_0$	$1.1127 \times 10^{-10} \text{ C}^2\cdot\text{J}^{-1}\cdot\text{m}^{-1}$

Born-Oppenheimer Approximation in words

Mass of electron versus mass of nucleus

1 <<< 1830 (at least)

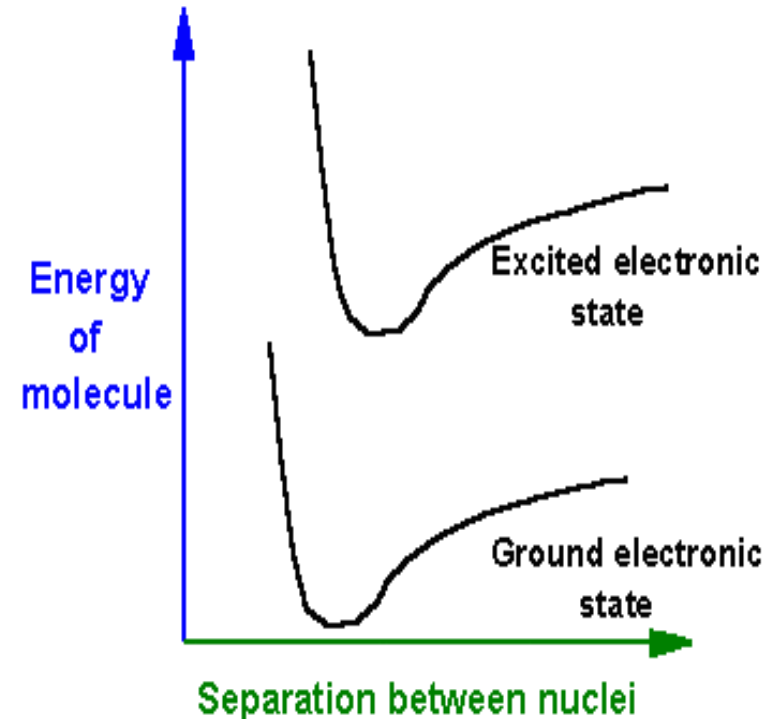
electron moves much faster than nucleus, so
electron can instantly adjust to change in nuclear
coordinate, so nucleus moves in an average field
made up by the electrons



At a given value for the nuclear geometry \mathbf{R} ,
there exist a well defined electronic state
distribution and Energy

$$\phi_n^{el}(\mathbf{r}; \mathbf{R}); E_n(\mathbf{R})$$

This electronic state depends on position of nucleus,
but not on the momentum



BO Approximation

Real problem: N nucleus n electron problem

$$\hat{H} = \hat{T} + \hat{V} = -\frac{\hbar^2}{2m_e} \sum_{i=1}^n \nabla_i^2 - \frac{\hbar^2}{2} \sum_{I=1}^N \frac{1}{M_I} \nabla_I^2 + V(\mathbf{r}, \mathbf{R}) = -\frac{\hbar^2}{2} \sum_{I=1}^N \frac{1}{M_I} \nabla_I^2 + \hat{H}^0$$

Born-Oppenheimer Approximation ignore nonadiabatic coupling

Solve the electronic state at fixed nuclear geometry,

$$\hat{H}^0(\mathbf{r}; \mathbf{R}) \phi_n^{el}(\mathbf{r}; \mathbf{R}) = E_n(\mathbf{R}) \phi_n^{el}(\mathbf{r}; \mathbf{R})$$

Quantum chemistry

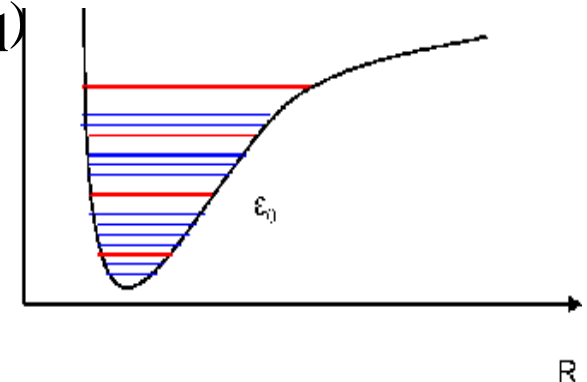
many nuclear geometries: **adiabatic potential energy surface**

$$\left(-\frac{1}{2} \sum_{I=1}^N \frac{1}{M_I} \nabla_I^2 + E_n(\mathbf{R}) \right) \chi_n(\mathbf{R}) = E_{el,NU} \chi_n(\mathbf{R})$$



Classical Molecular Dynamics,
Statistical Modeling

Quantum Dynamics
(Vibrational Schroedinger
Eq)



BO Approximation in equation 1

$$\hat{H} = \hat{T} + \hat{V} = -\frac{\hbar^2}{2m_e} \sum_{i=1}^n \nabla_i^2 - \frac{\hbar^2}{2} \sum_{I=1}^N \frac{1}{M_I} \nabla_I^2 + V(\mathbf{r}, \mathbf{R})$$

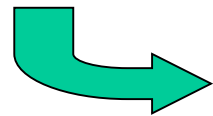
$$= \hat{T}_{el} + \hat{T}_{NU} + \hat{V} = \hat{T}_{el} + \hat{V} + \hat{T}_{NU} = \hat{H}^0 + \hat{T}_{NU}$$

$$\hat{H}^0(\mathbf{r}; \mathbf{R}) = -\frac{1}{2} \sum_{i=1}^n \nabla_i^2 + V(\mathbf{r}, \mathbf{R}) \quad \hat{T}_{NU} = -\frac{1}{2} \sum_{I=1}^N \frac{1}{M_I} \nabla_I^2$$

$$\hat{H}^0(\mathbf{r}; \mathbf{R}) \phi_n^{el}(\mathbf{r}; \mathbf{R}) = E_n(\mathbf{R}) \phi_n^{el}(\mathbf{r}; \mathbf{R})$$

$$\int \phi_n^{el*}(\mathbf{r}; \mathbf{R}) \phi_m^{el}(\mathbf{r}; \mathbf{R}) d\mathbf{r} = \delta_{mn}$$

$$\hat{H}\Psi(\mathbf{r}, \mathbf{R}) = E_{el,NU} \Psi(\mathbf{r}, \mathbf{R})$$



$$\Psi(\mathbf{r}, \mathbf{R}) = \sum_m^{\infty} \chi_m(\mathbf{R}) \phi_m^{el}(\mathbf{r}; \mathbf{R})$$

$\chi_m(\mathbf{R})$ = Expansion Coefficient

BO Approximation 2

$$\left(\hat{H} - E_{el,NU}\right)\Psi(\mathbf{r}, \mathbf{R}) = 0 \quad \begin{array}{l} \text{multiply } \phi_n \\ \text{then integrate with } \mathbf{r} \end{array}$$

$$\int \phi_n^{el*}(\mathbf{r}; \mathbf{R}) \left(\hat{H} - E_{el,NU}\right) \sum_m^{\infty} \chi_m(\mathbf{R}) \phi_m^{el}(\mathbf{r}; \mathbf{R}) d\mathbf{r} = 0$$

$$\int \phi_n^{el*}(\mathbf{r}; \mathbf{R}) \left(\hat{H}^0 + \hat{T}_{NU} - E_{el,NU}\right) \sum_m^{\infty} \chi_m(\mathbf{R}) \phi_m^{el}(\mathbf{r}; \mathbf{R}) d\mathbf{r} = 0$$

$$\int \phi_n^{el*}(\mathbf{r}; \mathbf{R}) \left(\hat{H}^0 - E_{el,NU} + \hat{T}_{NU}\right) \sum_m^{\infty} \chi_m(\mathbf{R}) \phi_m^{el}(\mathbf{r}; \mathbf{R}) d\mathbf{r} = 0$$

$$\begin{aligned} \sum_m^{\infty} \chi_m(\mathbf{R}) \left[\int \phi_n^{el*}(\mathbf{r}; \mathbf{R}) \left(\hat{H}^0 - E_{el,NU}\right) \phi_m^{el}(\mathbf{r}; \mathbf{R}) d\mathbf{r} \right] \\ + \int \phi_n^{el*}(\mathbf{r}; \mathbf{R}) \left(\hat{T}_{NU}\right) \sum_m^{\infty} \chi_m(\mathbf{R}) \phi_m^{el}(\mathbf{r}; \mathbf{R}) d\mathbf{r} = 0 \end{aligned}$$

BO Approximation 3

$$\sum_m^{\infty} \chi_m(\mathbf{R}) \left[\int \phi_n^{el*}(\mathbf{r}; \mathbf{R}) (\hat{H}^0 - E_{el,NU}) \phi_m^{el}(\mathbf{r}; \mathbf{R}) d\mathbf{r} \right] \quad \text{Remember} \quad \int \phi_n^{el*}(\mathbf{r}; \mathbf{R}) \phi_m^{el}(\mathbf{r}; \mathbf{R}) d\mathbf{r} = \delta_{mn}$$

$$= \sum_m^{\infty} \chi_m(\mathbf{R}) [(E_m(\mathbf{R}) - E_{el,NU}) \delta_{nm}] = (E_n(\mathbf{R}) - E_{el,NU}) \chi_n(\mathbf{R})$$

$$\int \phi_n^{el*}(\mathbf{r}; \mathbf{R}) (\hat{T}_{NU}) \sum_m^{\infty} \chi_m(\mathbf{R}) \phi_m^{el}(\mathbf{r}; \mathbf{R}) d\mathbf{r}$$

$$= \int \phi_n^{el*}(\mathbf{r}; \mathbf{R}) \left(-\frac{1}{2} \sum_{I=1}^N \frac{1}{M_I} \nabla_I^2 \right) \sum_m^{\infty} \chi_m(\mathbf{R}) \phi_m^{el}(\mathbf{r}; \mathbf{R}) d\mathbf{r}$$

$$|n\rangle = \psi_n(x) \quad \langle m| = \psi_m^*(x)$$

Using Bra-Ket
Notation

$$\langle n | \hat{x} | m \rangle = \int_{-\infty}^{\infty} \psi_n^*(x) x \psi_m(x) dx$$

$$\langle n | \hat{1} | m \rangle = \langle n | m \rangle = \langle n || m \rangle = \int_{-\infty}^{\infty} \psi_n^*(x) \psi_m(x) dx \quad 26$$

BO Approximation 4

$$\begin{aligned}
 & \int \phi_n^{el*}(\mathbf{r}; \mathbf{R}) \left(-\frac{1}{2} \sum_{I=1}^N \frac{1}{M_I} \nabla_I^2 \right) \sum_m \chi_m(\mathbf{R}) \phi_m^{el}(\mathbf{r}; \mathbf{R}) d\mathbf{r} \quad \text{Divide integration to different parts} \\
 &= \int \phi_n^{el*}(\mathbf{r}; \mathbf{R}) \sum_m \left(-\frac{1}{2} \sum_{I=1}^N \frac{1}{M_I} \nabla_I^2 \chi_m(\mathbf{R}) \right) \phi_m^{el}(\mathbf{r}; \mathbf{R}) d\mathbf{r} \\
 &\quad + \int \phi_n^{el*}(\mathbf{r}; \mathbf{R}) \sum_m \left(-\frac{1}{2} \sum_{I=1}^N \frac{1}{M_I} \nabla_I^2 \phi_m^{el}(\mathbf{r}; \mathbf{R}) \right) \chi_m(\mathbf{R}) d\mathbf{r} \\
 &\quad + \int \phi_n^{el*}(\mathbf{r}; \mathbf{R}) \sum_m \left(-\sum_{I=1}^N \frac{1}{M_I} \nabla_I \phi_m^{el}(\mathbf{r}; \mathbf{R}) \right) \nabla_I \chi_m(\mathbf{R}) d\mathbf{r} \\
 &= \sum_m \left(-\frac{1}{2} \sum_{I=1}^N \frac{1}{M_I} \nabla_I^2 \chi_m(\mathbf{R}) \right) \langle n \| m \rangle_r = -\frac{1}{2} \sum_{I=1}^N \frac{1}{M_I} \nabla_I^2 \chi_n(\mathbf{R}) \\
 &\quad + \sum_m \left\langle n \left| -\frac{1}{2} \sum_{I=1}^N \frac{1}{M_I} \nabla_I^2 \right| m \right\rangle \chi_m(\mathbf{R}) + \sum_m \left\langle n \left| -\sum_{I=1}^N \frac{1}{M_I} \nabla_I \right| m \right\rangle \nabla_I \chi_m(\mathbf{R})
 \end{aligned}$$

BO Approximation 5

Bring all
the things
together

$$\int \phi_n^{el*}(\mathbf{r}; \mathbf{R}) \left(E_n(\mathbf{R}) + \hat{T}_{NU} - E_{el,NU} \right) \sum_m \chi_m(\mathbf{R}) \phi_m^{el}(\mathbf{r}; \mathbf{R}) d\mathbf{r} = 0$$

$$= (E_n(\mathbf{R}) - E_{el,NU}) \chi_n(\mathbf{R}) - \frac{1}{2} \sum_{I=1}^N \frac{1}{M_I} \nabla_I^2 \chi_n(\mathbf{R}) + \sum_m C_{nm} \chi_m(\mathbf{R}) = 0$$

$$\sum_m C_{nm} \chi_m(\mathbf{R})$$

$$\equiv \sum_m \left\langle n \left| -\frac{1}{2} \sum_{I=1}^N \frac{1}{M_I} \nabla_I^2 \right| m \right\rangle \chi_m(\mathbf{R}) + \sum_m \left\langle n \left| -\sum_{I=1}^N \frac{1}{M_I} \nabla_I \right| m \right\rangle \nabla_I \chi_m(\mathbf{R})$$



Collect the parts other than total energy

$$\left(-\frac{1}{2} \sum_{I=1}^N \frac{1}{M_I} \nabla_I^2 + E_n(\mathbf{R}) \right) \chi_n(\mathbf{R}) + \sum_m C_{nm} \chi_m(\mathbf{R}) = E_{el,NU} \chi_n(\mathbf{R})$$

$$\hat{H}^0(\mathbf{r}; \mathbf{R}) \phi_n^{el}(\mathbf{r}; \mathbf{R}) = E_n(\mathbf{R}) \phi_n^{el}(\mathbf{r}; \mathbf{R})$$

BO Approximation 6

Born-Oppenheimer Approximation: **ignore all C_{nm}**

$$\left(-\frac{1}{2} \sum_{I=1}^N \frac{1}{M_I} \nabla_I^2 + E_n(\mathbf{R}) \right) \chi_n(\mathbf{R}) = E_{el,NU} \chi_n(\mathbf{R})$$

$$\hat{H}^0(\mathbf{r}; \mathbf{R}) \phi_n^{el}(\mathbf{r}; \mathbf{R}) = E_n(\mathbf{R}) \phi_n^{el}(\mathbf{r}; \mathbf{R})$$

Nuclear wavefunction is given by the expansion coefficient!

$$\hat{H}_{NU}(\mathbf{R}) \chi_{n,v_n}^{NU}(\mathbf{R}) = E_{el,NU} \chi_{n,v_n}^{NU}(\mathbf{R})$$

$$\hat{H}_{NU}(\mathbf{R}) = \left(-\frac{1}{2} \sum_{I=1}^N \frac{1}{M_I} \nabla_I^2 + V(\mathbf{R}) \right)$$

The potential energy that the nucleus feels $V(\mathbf{R})$ is the energy of the electron $E_n(\mathbf{R})$ at that geometry \mathbf{R} !

You have separated the motion of the electron and nucleus. 29

Conclusion of Born-Oppenheimer Approximation

$$\hat{H}^0(\mathbf{r}; \mathbf{R}) \phi_n^{el}(\mathbf{r}; \mathbf{R}) = E_n(\mathbf{R}) \phi_n^{el}(\mathbf{r}; \mathbf{R})$$

$$\hat{H}_{NU}(\mathbf{R}) \chi_{n,v_n}^{NU}(\mathbf{R}) = E_{el,NU} \chi_{n,v_n}^{NU}(\mathbf{R})$$

Now you can say nuclear wave function on the n-th electronic state

$$\Psi(\mathbf{r}, \mathbf{R}) = \phi_n^{el}(\mathbf{r}; \mathbf{R}) \chi_{n,v_n}^{NU}(\mathbf{R})$$

Electron moves so fast that it does not care about **how “fast”** (much slower compared to electron) the nucleus is moving.

Another way to consider is when the electron goes from a certain state to another the nucleus does not move:

Franck-Condon principle.