

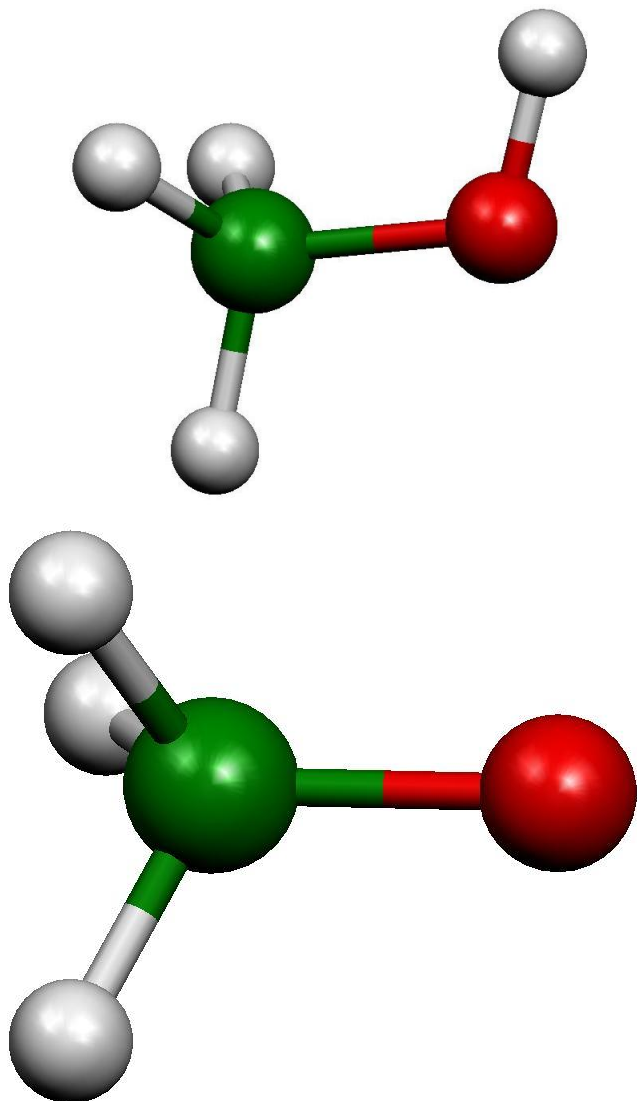
# Picking Basis Set

# Basis Sets

- Minimal Basis: only those in atomic orbital so one 1S orbital for hydrogen, one 1S, 2S, 2P for carbon, oxygen
  - STO-3G
- Split Valence: valence orbital has two, for hydrogen two 1S orbital, one 1S and two 2S, 2P for carbon oxygen
  - 6-31G, 3-21G
- Diffuse: large version of valence orbital
  - +, ++
- Polarization: higher angular momentum add 2P for hydrogen, add 3D for carbon, oxygen
  - \*, \*\*, (d), (d,p)

# Geometry

## Hartree Fock



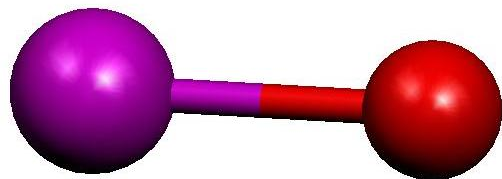
### i. Methoxide Anion Optimizations

We ran geometry optimizations of methanol (gauche form) and methoxide anion using both the 6-31G(d) and 6-31+G(d) basis sets in order to determine the effects of diffuse functions on the predicted structures. Here are the results:

Methanol	6-31G(d)	6-31+G(d)	Experiment
CO bond	1.3966	1.4019	$1.427 \pm 0.007$
CH bond	1.0873	1.0865	$1.096 \pm 0.01$
OH bond	0.9463	0.9464	$0.956 \pm 0.015$
COH angle	109.406	110.346	$108.9 \pm 2.0$
HCH angle	108.4127	108.6555	$109.3 \pm 0.75$
OCH angle	112.008	111.691	

Methoxide anion	6-31G(d)	6-31+G(d)	6-311++G(3df,2pd)
CO bond	1.3107	1.3304	1.3223
CH bond	1.1332	1.121	1.1209
HCH angle	101.5713	103.4298	103.2904
OCH angle	116.537	114.9919	115.1097

Diffuse functions have very little effect on the optimized structure of methanol but do significantly affect the bond angles in negatively charged methoxide anion. We can conclude that they are required to produce an accurate structure for the anion by comparing the two calculated geometries to that predicted by Hartree-Fock theory at a very large basis set (which should eliminate basis set effects).



# PO bond length

B3LYP

Basis Set	Bond Length (Å)
6-31G(d)	1.4986
6-311G(d)	1.4914
6-311G(2d)	1.4818
6-311G(2df)	1.4796
6-311G(3df)	1.4758

Experiment

1.476

# Pick Basis Set Convergence

## Dipole moment of H<sub>2</sub>O

Method	# of basis	Debye
B3LYP/STO-3G	7	1.5936
B3LYP/6-31G	13	2.3986
B3LYP/6-31+G	17	2.5458
B3LYP/6-31+G(d,p)	29	2.1951
B3LYP/6-311G	19	2.4296
B3LYP/6-311++G	25	2.5240
B3LYP/6-311++G(d,p)	37	2.1592
B3LYP/6-311++G(3d,3p)	61	1.8963
B3LYP/6-311++G(3df,3pd)	83	1.8897
B3LYP/aug-cc-pVTZ	105	1.8473
B3LYP/aug-cc-pVQZ	215	1.8458
Exp		1.8550

Computational time  $\sim (\# \text{ of basis})^2$

# Beyond Hartree Fock: Electron Correlation Methods

# Hartree Fock Review

- Hartree Fock Results=By considering one Slater determinant we obtained the best results

Molecular Orbital  $\phi_a(\mathbf{r}_1) = \sum_{u=1}^{Nbasis} C_{ua} \theta_u \quad a = 1, 2, \dots, Nbasis$

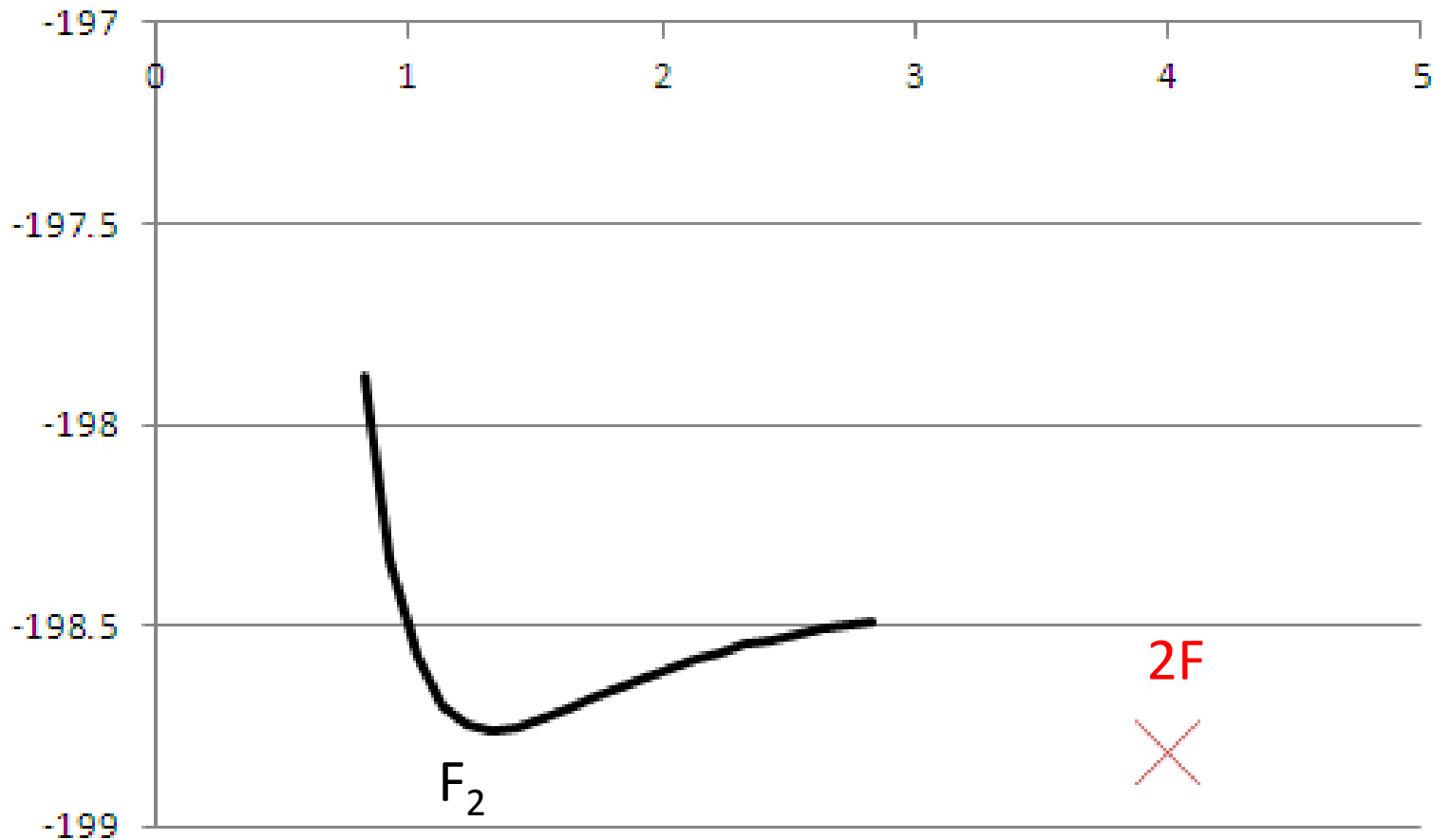
Spin Orbital  $\psi_i(\mathbf{x}) = \phi_a(\mathbf{r}_1) \sigma_m(s) \quad i = 1, 2, \dots, Nbasis, Nbasis + 1 \dots 2Nbasis$   
 $a = 1, 2 \dots Nbasis; \quad m = \alpha, \beta$

$$\Psi(\mathbf{x}_1, \mathbf{x}_2, \dots, \mathbf{x}_n) = \frac{1}{\sqrt{n!}} \begin{vmatrix} \psi_i(\mathbf{x}_1) & \psi_j(\mathbf{x}_1) & \dots & \psi_n(\mathbf{x}_1) \\ \psi_i(\mathbf{x}_2) & \psi_j(\mathbf{x}_2) & \dots & \psi_n(\mathbf{x}_2) \\ \dots & \dots & \dots & \dots \\ \psi_i(\mathbf{x}_n) & \psi_j(\mathbf{x}_n) & \dots & \psi_n(\mathbf{x}_n) \end{vmatrix}$$

$$= \left\| \begin{matrix} \psi_i & \psi_j & \dots & \psi_n \end{matrix} \right\|$$

One Slater Determinant

# Failure of Hartree Fock



The cross is the energy for two F atoms.

For Hartree Fock the minimum energy for F<sub>2</sub> molecule at 1.2 angstrom is higher then the energy for two F atoms = no bond???



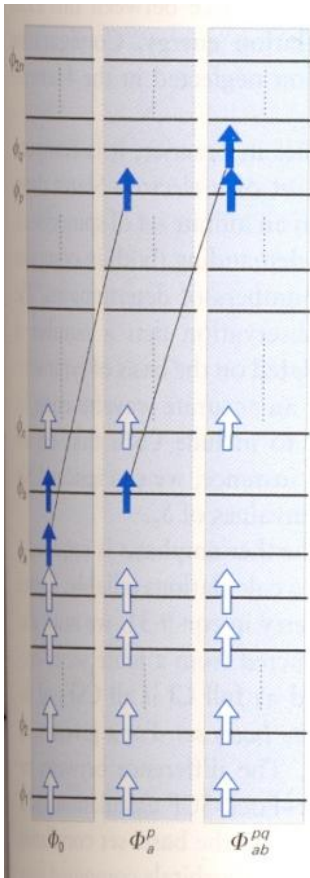
# So to go beyond Hartree Fock

- Use more than one Slater determinant
  - CISD: Configuration Interaction Singles and Doubles
  - CCSD: Coupled Cluster Singles and Doubles
  - MCSCF: Multiconfigurational Self Consistent Field
  - MR-CISD: Multireference CISD
- Add in Corrections to Hartree Fock Energy
  - MP2,3,4: Moller Plesset perturbation theory
- Explicitly Correlated methods F12
- Totally Change the Hamiltonian
  - Density Functional Theory Method B3LYP
- Mixed methods G3, CBS

# Configuration Interaction SD

n-electron system HF solution  $D_0 = \left\| \psi_1 \quad \psi_2 \quad \dots \psi_i \psi_j \dots \psi_n \right\|$

You have  $2N_{\text{basis}} - n$  unoccupied orbitals so you can use them



Single excitation from HF solution

$$D_i^a = \left\| \psi_1 \quad \psi_2 \quad \dots \psi_a \psi_j \dots \psi_n \right\|$$

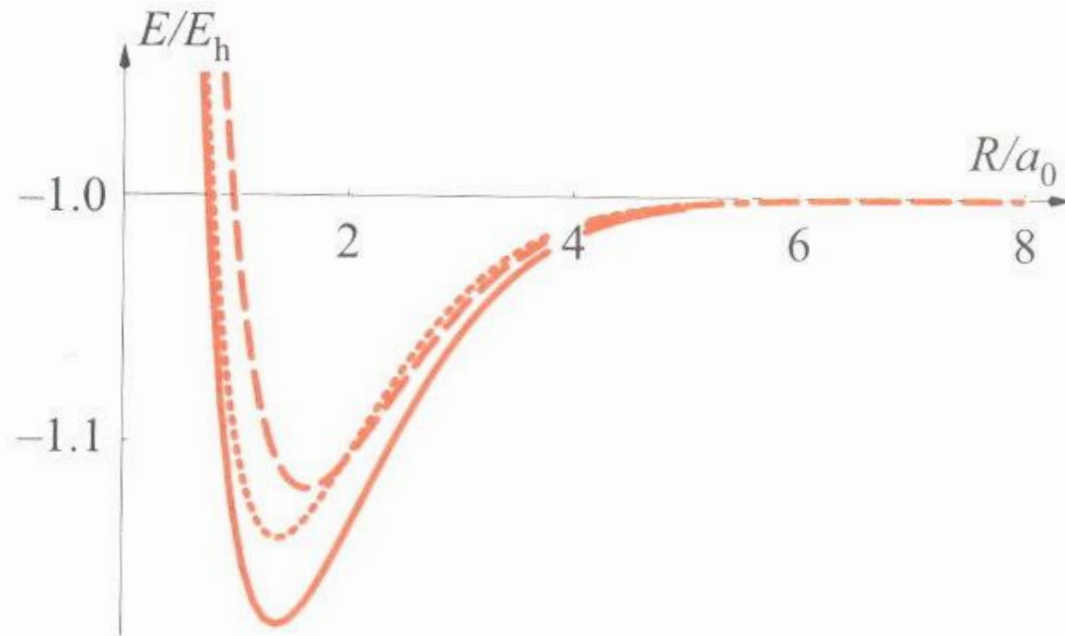
Double excitation from HF solution

$$D_{ij}^{ab} = \left\| \psi_1 \quad \psi_2 \quad \dots \psi_a \psi_b \dots \psi_n \right\|$$

$$CI = C_0 D_0 + \sum_{i,a} C_i^a D_i^a + \sum_{i < j, a < b} C_{ij}^{ab} D_{ij}^{ab} + \sum_{i < j < k, a < b < c} C_{ijk}^{abc} D_{ijk}^{abc} + \dots$$

# H<sub>2</sub> Potential Curve Revisited Review

This is equivalent to using CI double excitation for the wave function



**FIGURE 10.25**

The configuration-interaction energy  $E_{CI}$  of the ground-state energy of H<sub>2</sub> for  $\zeta = 1$  (dashed curve) and for an optimized value of  $\zeta$  (dotted curve) plotted against  $R$ . The “exact” results of Kolos and Wolniewicz (solid curve) are shown for comparison.

# H2 Configuration Interaction Review

Two 1S orbitals can make **TWO** molecular orbitals

Why not use the two and make combinations

$$|\Psi_1\rangle = C_1 \begin{vmatrix} \alpha(1)|+\rangle_1 & \beta(1)|+\rangle_1 \\ \alpha(2)|+\rangle_2 & \beta(2)|+\rangle_2 \end{vmatrix}$$

$$|\Psi_1\rangle \approx |++\rangle(\alpha\beta - \beta\alpha)$$

$$|\Psi_2\rangle = C_2 \begin{vmatrix} \alpha(1)|-\rangle_1 & \beta(1)|-\rangle_1 \\ \alpha(2)|-\rangle_2 & \beta(2)|-\rangle_2 \end{vmatrix}$$

$$|\Psi_2\rangle \approx |--\rangle(\alpha\beta - \beta\alpha)$$

Configuration 1: two  
electron in bonding orbital

$$|\Psi_1\rangle \approx |++\rangle(\alpha\beta - \beta\alpha)$$

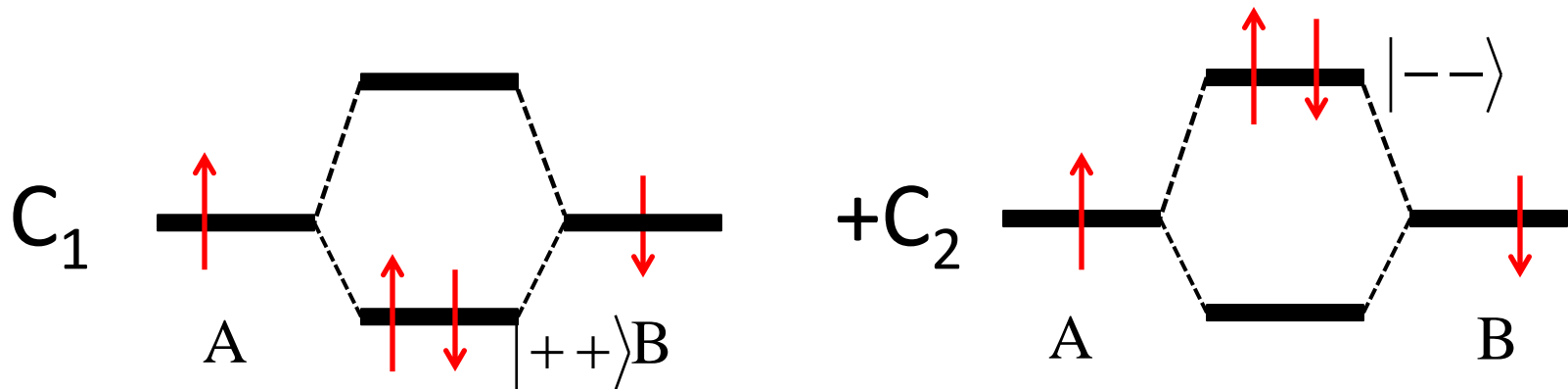
Configuration 2: two  
electron in antibonding orbital

$$|\Psi_2\rangle \approx |--\rangle(\alpha\beta - \beta\alpha)$$

$$|\Psi_{CI}\rangle = C_1|\Psi_1\rangle + C_2|\Psi_2\rangle = C_1|++\rangle + C_2|--\rangle$$

# What we did in H2 is minimal basis CID

Using two 1S orbitals you can get two molecular orbitals



Hartree Fock Solution

Double excitation from  
Hartree Fock Solution

Use variational theory to calculate the values of  $C_1$  and  $C_2$

In most cases due to limitation of computational power we cut at the double excitation CISD, MRCISD, CCSD

# Brillouin's Theorem

Interaction between the Hartree Fock solution and the one electron excitation determinant is zero

$$\begin{aligned}\langle D_0 | H | D_i^a \rangle &= \langle i | h | a \rangle + \sum_{j=1}^{n_{\text{electron}}} \langle ij | aj \rangle - \langle ij | ja \rangle = h_{ia} + \sum_{j=1}^{n_{\text{electron}}} \langle i | J_j | a \rangle + \langle i | K_j | a \rangle \\ &= \langle i | f | a \rangle = \delta_{ij}\end{aligned}$$

Interaction energy is the cross term of the Fock matrix and by definition it is zero in the canonical Hartree Fock

However two one electron excitations can have interaction and the one electron excitation and two electron interaction can have interaction, so if Hartree Fock is a good solution the two electron excitation contribution is the greatest

# One Electron Operator Matrix Element (two electron system)

$$\begin{aligned}
 \langle D_0 | h_1 | D_i^a \rangle &= 2^{-1} \iint d\mathbf{x}_1 d\mathbf{x}_2 \left( \psi_i(\mathbf{x}_1) \psi_j(\mathbf{x}_2) - \psi_j(\mathbf{x}_1) \psi_i(\mathbf{x}_2) \right)^* \\
 &\quad \times h(\mathbf{r}_1) \left( \psi_a(\mathbf{x}_1) \psi_j(\mathbf{x}_2) - \psi_j(\mathbf{x}_1) \psi_a(\mathbf{x}_2) \right) \\
 &= 2^{-1} \iint d\mathbf{x}_1 d\mathbf{x}_2 \left[ \psi_i^*(\mathbf{x}_1) \psi_j^*(\mathbf{x}_2) h(\mathbf{r}_1) \psi_a(\mathbf{x}_1) \psi_j(\mathbf{x}_2) \right. \\
 &\quad + \psi_j^*(\mathbf{x}_1) \psi_i^*(\mathbf{x}_2) h(\mathbf{r}_1) \psi_j(\mathbf{x}_1) \psi_a(\mathbf{x}_2) \\
 &\quad - \psi_i^*(\mathbf{x}_1) \psi_j^*(\mathbf{x}_2) h(\mathbf{r}_1) \psi_j(\mathbf{x}_1) \psi_a(\mathbf{x}_2) \\
 &\quad \left. - \psi_j^*(\mathbf{x}_1) \psi_i^*(\mathbf{x}_2) h(\mathbf{r}_1) \psi_a(\mathbf{x}_1) \psi_j(\mathbf{x}_2) \right]
 \end{aligned}$$

Due to orthonormality of the spin orbitals last three terms are zero with integration with respect to  $\mathbf{x}_2$

# One Electron Operator 2

$$\begin{aligned}\langle D_0 | h_1 | D_i^a \rangle &= 2^{-1} \iint d\mathbf{x}_1 d\mathbf{x}_2 \left[ \psi_i^*(\mathbf{x}_1) \psi_j^*(\mathbf{x}_2) h(\mathbf{r}_1) \psi_a(\mathbf{x}_1) \psi_j(\mathbf{x}_2) \right] \\ &= 2^{-1} \int \psi_i^*(\mathbf{x}_1) h(\mathbf{r}_1) \psi_a(\mathbf{x}_1) d\mathbf{x}_1 \\ &= 2^{-1} \left( \langle \psi_i | h | \psi_a \rangle \right) \\ \langle D_0 | h_2 | D_i^a \rangle &= 2^{-1} \left( \langle \psi_i | h | \psi_a \rangle \right)\end{aligned}$$



# Two Electron Operator Matrix Element

$$\begin{aligned}
 \left\langle D_0 \left| |\mathbf{r}_{12}|^{-1} \right| D_i^a \right\rangle &= 2^{-1} \iint d\mathbf{x}_1 d\mathbf{x}_2 \left( \psi_i(\mathbf{x}_1) \psi_j(\mathbf{x}_2) - \psi_j(\mathbf{x}_1) \psi_i(\mathbf{x}_2) \right)^* \\
 &\quad \times |\mathbf{r}_{12}|^{-1} \left( \psi_a(\mathbf{x}_1) \psi_j(\mathbf{x}_2) - \psi_j(\mathbf{x}_1) \psi_a(\mathbf{x}_2) \right) \\
 &= 2^{-1} \iint d\mathbf{x}_1 d\mathbf{x}_2 \left[ \psi_i^*(\mathbf{x}_1) \psi_j^*(\mathbf{x}_2) |\mathbf{r}_{12}|^{-1} \psi_a(\mathbf{x}_1) \psi_j(\mathbf{x}_2) \right. \\
 &\quad + \psi_j^*(\mathbf{x}_1) \psi_i^*(\mathbf{x}_2) |\mathbf{r}_{12}|^{-1} \psi_j(\mathbf{x}_1) \psi_a(\mathbf{x}_2) \\
 &\quad - \psi_i^*(\mathbf{x}_1) \psi_j^*(\mathbf{x}_2) |\mathbf{r}_{12}|^{-1} \psi_j(\mathbf{x}_1) \psi_a(\mathbf{x}_2) \\
 &\quad \left. - \psi_j^*(\mathbf{x}_1) \psi_i^*(\mathbf{x}_2) |\mathbf{r}_{12}|^{-1} \psi_a(\mathbf{x}_1) \psi_j(\mathbf{x}_2) \right]
 \end{aligned}$$

Since  $|\mathbf{r}_{12}|$  is equal to  $|\mathbf{r}_{21}|$  perform interchange of electron 1 and electron 2 and first and second are the same and third and fourth are the same

# Two Electron Operator 2

$$\begin{aligned}\langle D_0 | |\mathbf{r}_{12}|^{-1} | D_i^a \rangle &= \iint d\mathbf{x}_1 d\mathbf{x}_2 \left[ \psi_i^*(\mathbf{x}_1) \psi_j^*(\mathbf{x}_2) |\mathbf{r}_{12}|^{-1} \psi_a(\mathbf{x}_1) \psi_j(\mathbf{x}_2) \right. \\ &\quad \left. - \psi_i^*(\mathbf{x}_1) \psi_j^*(\mathbf{x}_2) |\mathbf{r}_{12}|^{-1} \psi_j(\mathbf{x}_1) \psi_a(\mathbf{x}_2) \right] \\ &= \langle ij | aj \rangle - \langle ij | ja \rangle\end{aligned}$$

# Coupled Cluster

- Coupled cluster is a smarter way to do CI

We have Hartree Fock solution

$$D_0 = \left\| \begin{matrix} \psi_1 & \psi_2 & \dots & \psi_i & \psi_j & \dots & \psi_n \end{matrix} \right\|$$

$$CC = \exp^T D_0$$

$$T = T_1 + T_2 + T_3 + \dots$$

$T_n$  is n electron excitation operator

$$T_2 D_0 = \sum_{i < j, a < b} C_{ij}^{ab} D_{ij}^{ab}$$

When using Couple cluster or CCSD or CCSD(T) what do you have to be careful for?

CC assumes Hartree Fock is a good answer so if Hartree Fock is a bad answer CCSD is bad! CHECK T1 DIAGNOSTIC (tells you how much one electron excitation is contributing to the wave function)

# T1 Diagnostic in Gaussian

140.109.112.238:22 - Tera Term VT

File Edit Setup Control Window Resize Help

```
278659.master      ...BtOn_NoSymOpt garyer      356:46:4 R hp
278699.master      96_water_300_NVT hyming      2052:58: R ibm2
278704.master      job                  lwchou      319:58:1 R ibm
278731.master      ...PWlPW9l_rootl mktsai      243:47:0 R ibm2
278742.master      job                  lwchou      221:59:1 R ibm
278766.master      WLMC                mikechwu    25:39:43 R hp
278767.master      WLMC                mikechwu    25:40:59 R hp
282637.master      MnWl1l1_10         ktliu       274:25:2 R ibm
282669.master      job                  yanjx       282:15:4 R ibm2
282671.master      PTMC                mikechwu    803:47:5 R hp
282673.master      ..._NoSymOpt_PCM   garyer      98:44:52 R hp
282674.master      ..._NoSymOpt_PCM   garyer      96:50:05 R hp
282679.master      ...pt-t_mPWlPW9l  mktsai      46:06:39 R ibm2
282681.master      casino              crhsing     00:00:00 R hp
282682.master      casino              crhsing     00:00:00 R hp
282683.master      casino              crhsing     00:00:00 R ibm
282689.master      smallstuff          kaito       0 R testibm2
```

kaito@master:/lustre/lwork/kaito/kaito/G09/h2/a63+/ccsd> more h2.com

#P CCSD(T1Diag)/aug-cc-pVTZ pop=reg

Title

0 1

H1

H2, H1, RH1H2

RH1H2=0.76092319

kaito@master:/lustre/lwork/kaito/kaito/G09/h2/a63+/ccsd> █

# T1 Diagnostic for H2

If T1 diagnostic greater than 0.03 means Hartree Fock is not a good starting point

Last time we show that Hartree Fock is not good for long distances

140.109.112.238:22 - Tera Term VT

File Edit Setup Control Window Resize Help

```
*****
DDI Dir will call FoFMem 1 times, MxPair= 2
NAB= 1 NAA= 0 NBB= 0.
Norm of the A-vectors is 1.0456581D-05 conv= 1.00D-05.
RLE energy= -0.0398507293
DE(Corr)= -0.39850725E-01 E(CORR)= -1.1724313370 Delta= 8.63D-08
NORM(A)= 0.10094743D+01
Iteration Nr. 7
*****
DDI Dir will call FoFMem 1 times, MxPair= 2
NAB= 1 NAA= 0 NBB= 0.
Norm of the A-vectors is 2.2535773D-06 conv= 1.00D-05.
RLE energy= -0.0398507304
DE(Corr)= -0.39850721E-01 E(CORR)= -1.1724313336 Delta= 3.41D-09
NORM(A)= 0.10094743D+01
CI/CC converged in 7 iterations to DelEn= 3.41D-09 Conv= 1.00D-07 ErrA1= 2.25D-06 Conv= 1.00D-05
T1 Diagnostic = 0.00566332
Largest amplitude= 3.45D-02
Leave Link 913 at Sun Mar 27 15:15:58 2011, MaxMem= 33554432 cpu: 0.2
(Enter /home/software/g09-i7/g09/l60l.exe)
Copying SCF densities to generalized density rwf, IOpC1= 0 IROHF=0.

*****

Population analysis using the SCF density.

*****

Orbital symmetries:
h2.log lines 378-406/1125 30%
```

H<sub>2</sub> at equilibrium 0.76 Angstrom

# T1 Diagnostic for H2

140.109.112.238:22 - Tera Term VT

File Edit Setup Control Window Resize Help

```
Iteration Nr. 9
*****
DDlDir will call FoFMem 1 times, MxPair= 2
NAB= 1 NAA= 0 NBB= 0.
Norm of the A-vectors is 1.0092897D-06 conv= 1.00D-05.
RLE energy= -0.0781418604
DE(Corr)= -0.78141767E-01 E(CORR)= -1.0385485035 Delta= 1.78D-07
NORM(A)= 0.10846184D+01
Iteration Nr. 10
*****
DDlDir will call FoFMem 1 times, MxPair= 2
NAB= 1 NAA= 0 NBB= 0.
Norm of the A-vectors is 1.2925129D-07 conv= 1.00D-05.
RLE energy= -0.0781418335
DE(Corr)= -0.78141845E-01 E(CORR)= -1.0385485816 Delta=-7.80D-08
NORM(A)= 0.10846184D+01
CI/CC converged in 10 iterations to DelEn=-7.80D-08 Conv= 1.00D-07 ErrA1= 1.29D-07 Conv= 1.00D-05
T1 Diagnostic = 0.04519794
Dominant configurations:
*****
Spin Case I J A B Value
ABAB 1 1 2 2 -0.312088D+00
ABAB 1 1 2 4 0.135697D+00
ABAB 1 1 4 2 0.135697D+00
Largest amplitude= 3.12D-01
Leave Link 913 at Sun Mar 27 15:15:59 2011, MaxMem= 33554432 cpu: 0.3
(Enter /home/software/g09-i7/g09/l601.exe)
Copying SCF densities to generalized density rwf, IOpC1= 0 IROHF=0.
```

h2long.log lines 398-426/1153 31%

H<sub>2</sub> at longer 1.76 Angstrom

Not very good to use this CCSD results

# CASSCF/MRCI

- You pick which configuration you want to put into the summation of configuration interaction
- CASSCF optimizes the orbitals in each determinant as well as the coefficient

## Why Use?

1. If you know the problem correctly, you can choose only the important configuration more efficiently

# MP2, 3, 4

- Perturbation Theory: Add in correction one by one

$$E^{MP2} = E^{HF} + E^{correction2}$$

$$E^{MP3} = E^{MP2} + E^{correction3}$$

$$E^{MP4} = E^{MP3} + E^{correction4}$$

Why Use?

1. Easy to do after SCF results
2. Can get good accuracy for van der waals interactions



# Perturbation Convergence

STO-3G

Method	HCN	CN <sup>-</sup>	CN
MP2	-91.82033	-91.07143	-91.11411
MP3	-91.82242	-91.06862	-91.12203
MP4	-91.82846	-91.07603	-91.13538
MP5	-91.83129	-91.07539	-91.14221
MP6	-91.83233	-91.07694	-91.14855
MP7	-91.83264	-91.07678	-91.15276
MP8	-91.83289	-91.07699	-91.15666
Full CI	-91.83317	-91.07706	-91.17006
$\Delta E < 0.001$ at	MP6	MP6	MP19
Full CI – MP4 (kcal-mol <sup>-1</sup> )	-2.96	-0.65	-21.76

# CCSD(T)

- CCSD with contributions coming from Triples excitation is done by perturbation “golden standard” of quantum chemistry

$$E^{CCSD(T)} = E^{CCSD} + E^{Triples}$$

## Explicitly Correlated F12

- MP2-F12, CCSD-F12 are methods where two electron distance  $r_{ij}$  is explicitly in the basis set of the calculation. Usually this uses density fitting and resolution of identity approximation so needs to define three basis.

# Valance and Full Correlation

- Usually most post-Hartree Fock calculation are performed for only the valance electrons, since valance electrons are most important
- However to get high accuracy one has to consider the core correlation energy in that case one is recommended to use basis set that also include core correlation using usual basis that was made for valance correlation is not good

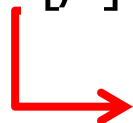
Aug-ccpCVXZ basis

# Density Function Theory

- Instead of getting the wavefunction let's get the correct density

$$H\Psi = E\Psi \Rightarrow E[\rho]$$

$$\rho(r) = \Psi^*(r)\Psi(r)$$



Problem is no one knows this relationship  
Many people have thought of approximate  
solution: B3LYP, PBE, BLYP.....

Why?

1. Can get good answer with SCF
2. Density is much easier to calculate than orbital



Figure 1: Total charge density for the H2O molecule, created by Waveplot, visualised by VMD.

## H2O Molecular Orbitals - Mozilla Firefox

File Edit View History Bookmarks Tools Help

<http://vergil.chemistry.gatech.edu/courses/chem3412/handouts/h2o-mo.html>

H2O molecule orbital

Most Visited Getting Started Latest Headlines Customize Links Free Hotmail

H2O Molecular Orbitals

### H2O Molecular Orbitals



1a<sub>1</sub>

Energy (a.u.): -20.566



2a<sub>1</sub>

Energy (a.u.): -1.325



1b<sub>2</sub>

Energy (a.u.): -0.690



3a<sub>1</sub>

Energy (a.u.): -0.564



1b<sub>1</sub>

Energy (a.u.): -0.496

<http://vergil.chemistry.gatech.edu/courses/chem3412/handouts/h2o-mo.html>

# Single Point of Methanol Time

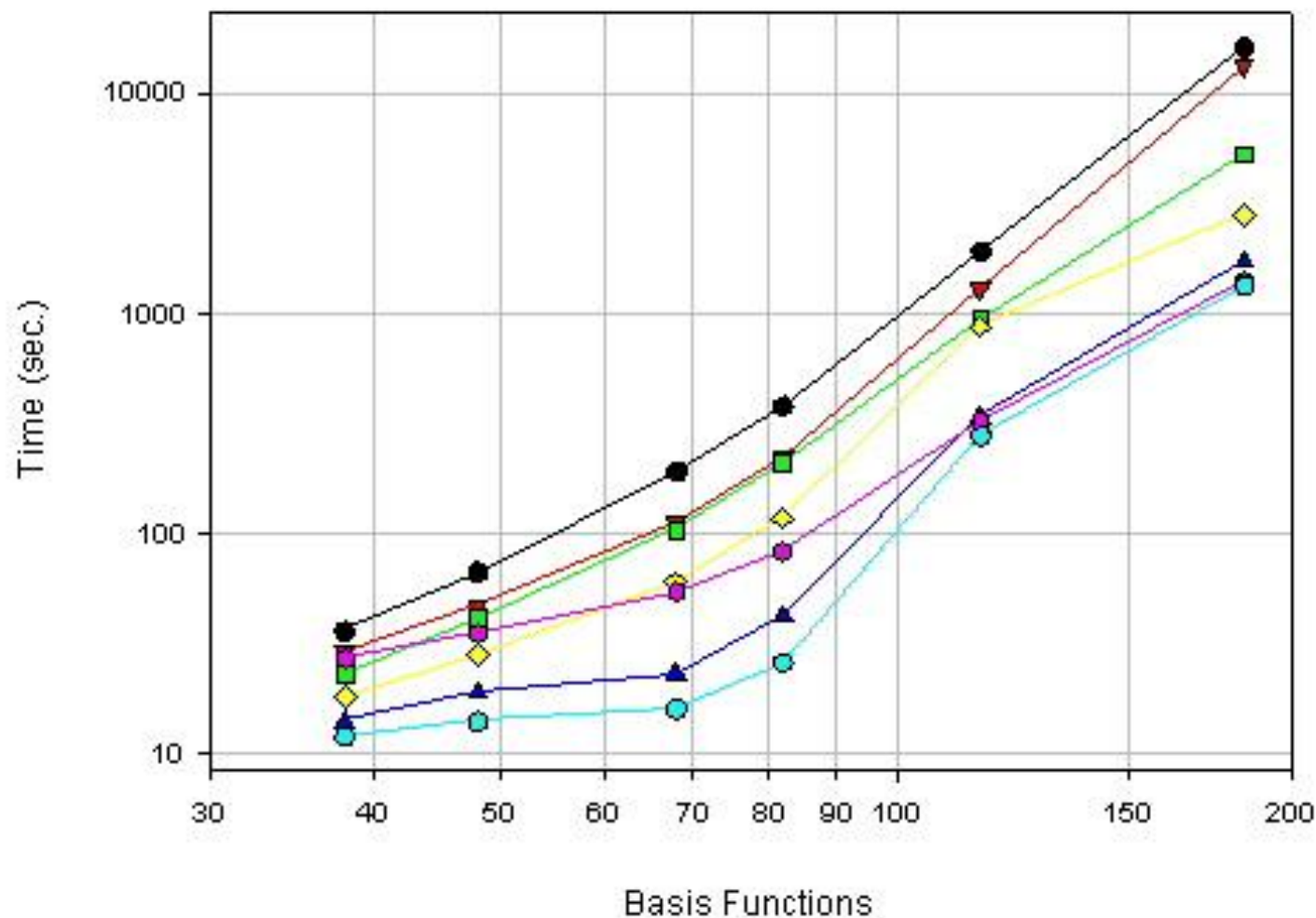
Times are wall-clock seconds for job to run. Gaussian 98 (IBM-RS6000-G98RevA.4) was used

Level	6-31G*	6-311+G**	cc-pVDZ	aug-cc-pVDZ	cc-pVTZ	aug-cc-pVTZ
HF	12	16	14	26	277	1335
BLYP	26	55	36	82	283	1189
B3LYP	27	54	35	83	325	1402
B3PW91	28	55	36	86	316	1404
MP2FC	14	23	19	42	340	1717
MP2FU	14	23	18	43	346	1750
MP4FC	23	104	41	209	946	5278
CID	15	49	24	101	843	6222
CISD	17	57	27	119	876	6458
CCD	18	60	28	118	871	2816
QCISD	23	82	38	153	1130	12087
CCSD	29	113	47	221	1298	13147
QCISD(T)	31	158	57	307	1653	15168
CCSD(T)	36	191	67	380	1912	16073
<i>basis functions</i>	<i>38</i>	<i>68</i>	<i>48</i>	<i>82</i>	<i>116</i>	<i>184</i>

For a graph of part of the table above click [here](#)

# Single Point of Methanol Time

Single Point Times CH<sub>3</sub>OH



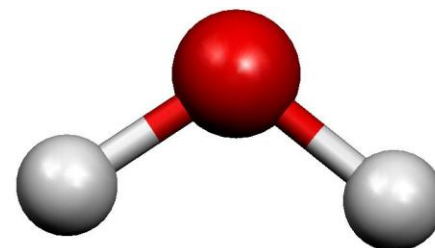
# Present State of Art Gas Phase Water results: using time independent

**Table 1.** Predicted VBOs for various theoretical models. Results are presented as differences from the observed values (Obs) in  $\text{cm}^{-1}$  (34). The standard deviation,  $\sigma$ , is for all experimentally known VBOs. 5Z, aug-cc-pV5Z MRCI calculation; 6Z, aug-cc-pV6Z MRCI calculation; CBS, MRCI calculation extrapolated to the complete basis set limit; PS, partially augmented cc-pV5Z MRCI

calculation plus core correlation owing to Partridge and Schwenke (5); CBS + CV, CBS with core correlation correction; Rel, CBS + CV with relativistic effects included; QED, Rel with one electron Lamb shift included; BODC, QED with Born–Oppenheimer diagonal correction included; Nonad, BODC with vibrational nonadiabatic effects included. Dashes indicate no data available.

State	Obs	5Z	6Z	CBS	PS	CBS + CV	Rel	QED	BODC	Nonad
(010)	1,594.74	-2.99	-2.29	-0.32	-2.79	0.48	-0.81	-0.75	-0.32	-0.27
(020)	3,151.63	-4.22	-2.38	-0.78	-5.38	1.16	-1.57	-1.44	-0.56	-0.44
(030)	4,666.78	-6.30	-3.24	-1.52	-7.91	2.05	-2.37	-2.16	-0.78	-0.60
(040)	6,134.01	-9.81	-5.53	-2.74	-10.38	3.20	-3.30	-3.00	-1.06	-0.83
(050)	7,542.43	-14.70	-9.18	-4.71	-12.90	4.82	-4.45	-4.02	-1.41	-1.14
(101)	7,249.81	12.51	10.76	9.32	-4.78	-5.35	1.70	1.43	0.60	2.00
(201)	10,613.35	18.72	16.46	13.97	-6.96	-7.47	2.98	2.57	1.23	-
(301)	13,830.93	25.72	22.81	18.74	-8.41	-8.95	4.59	4.06	2.05	-
(401)	16,898.84	32.56	28.92	23.06	-9.47	-10.17	6.11	5.49	2.74	-
(501)	19,781.10	40.72	35.96	28.68	-9.31	-10.72	9.04	8.28	4.65	-
(601)	22,529.44	51.14	43.41	34.17	-7.61	-11.88	11.69	10.81	5.94	-
(701)	25,120.27	63.29	51.75	38.66	-5.49	-13.13	13.70	12.75	6.46	-
All	$\sigma$	22.84	19.74	16.56	10.44	7.85	4.23	3.83	1.90	-

Complete basis set gets you to 15  $\text{cm}^{-1}$  accuracy  
 Addition of core valence gets you to 8  $\text{cm}^{-1}$   
 Relativistic lowers 4  $\text{cm}^{-1}$ , QED does not change much  
 and addition of nonBO gets you to 2  $\text{cm}^{-1}$



Include transitions to 30,000  $\text{cm}^{-1}$

O L Polyansky et al. Science 299, 539 (2003)

# Mixed Methods G2

Components of G1 and G2 Total Energies

Step	Job	Result	Notes
1	HF/6-31G(d) Opt Freq	ZPE	Scale by 0.8929.
2	MP2(Full)/6-31G(d) Opt	<i>geometry</i>	Start from HF results; use this geometry for all later jobs.
3	MP4/6-311G(d,p) <sup>†</sup>	E <sup>base</sup>	Base level energy.
4	MP4/6-311+G(d,p)	$\Delta E^+$	= Energy - E <sup>base</sup> .
5	MP4/6-311G(2df,p)	$\Delta E^{2df}$	= Energy - E <sup>base</sup> (set to 0 if > 0).
6	QCISD(T)/6-311G(d,p) <sup>†</sup>	$\Delta E^{QCI}$	= Energy - E <sup>base</sup>
7	Any job	$\Delta E^{HLC}$	= -0.00019n <sub>α</sub> + -0.00595n <sub>β</sub>
8	MP2/6-311+G(3df,2p)	$\Delta^{G2}$	= Energy - E <sup>Step5(MP2)</sup> - E <sup>Step4(MP2)</sup> + E <sup>Step3(MP2)</sup>
9	Any job	$\Delta^{HLC}$	= +0.00114n <sub>β</sub>

<sup>†</sup> These quantities are computed in a single job.

Use different methods to obtain values and add up contributions that are estimated by small basis set



# Mixed Methods CBS-Q

## Components of CBS Methods

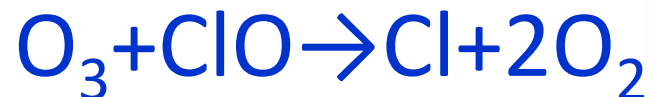
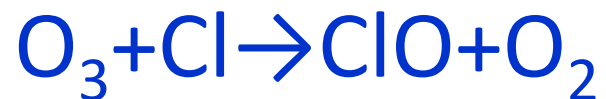
Energy Component	CBS-4	CBS-Q
Optimized geometry	HF/3-21G(d)	MP2/6-31G(d)
ZPE (scale factor)	HF/3-21G(d) (0.91671)	HF/6-31G <sup>†</sup> (0.91844)
SCF energy	HF/6-311+G(3d2f,2df,p)	HF/6-311+G(3d2f,2df,2p)
2nd order correlation	MP2/6-31+G <sup>†</sup>	MP2/6-311+G(3d2f,2df,2p)
CBS extrapolation	≥5 configurations	≥10 configurations
Higher order correlation	MP4(SDQ)/6-31G	MP4(SDQ)/6-31+G(d(f),d,f) QCISD(T)/6-31+G <sup>†</sup>
Additional empirical corrections	1 and 2-electron higher-order corrections (size-consistent), spin contamination	2-electron higher-order correction (size-consistent), spin contamination, core correlation for sodium

# Level of Accuracy and time

Model Chemistry	MAD	Standard Deviation	Absolute Max. Error
CBS-Q	1.0	0.8	3.8
G2	1.2	0.9	5.1
G2(MP2)	1.5	1.2	6.2
G1	1.6	1.4	9.2
CBS-4	2.0	1.5	7.0
B3LYP/6-311+G(3df,2df,2p) // B3LYP/6-31G(d)	2.7	2.6	12.5
B3LYP/6-311+G(2d,p) // B3LYP/6-311+G(2d,p)	3.1	3.0	19.7
B3LYP/6-311+G(2d,p) // B3LYP/6-31G(d)	3.2	3.0	20.1
B3LYP/6-311+G(2d,p) // HF/3-21G(d)	3.2	3.0	21.2
BLYP/6-31+G(d,p) // BLYP/6-31+G(d,p)	3.9	3.2	15.2
BLYP/6-311+G(2d,p) // BLYP/6-311+G(2d,p)	3.9	3.2	15.9
B3LYP/6-31+G(d,p) // B3LYP/6-31+G(d,p)	3.9	4.2	33.8
B3LYP/6-31+G(d,p) // B3LYP/6-31G(d)	4.0	4.2	33.9
B3LYP/6-31G(d) // B3LYP/6-31G(d)	7.9	9.5	54.2
B3LYP/6-31G(d) // HF/3-21G(d)	8.0	9.4	54.2
MP2/6-311+G(2d,p) // B3LYP/6-31G(d)	8.9	7.8	39.2
MP2/6-311+G(2d,p) // MP2/6-311+G(2d,p)	8.9	7.8	39.2
B3LYP/6-31G(d) // AM1	10.5	11.3	54.2
MP2/6-31+G(d,p) // MP2/6-31+G(d,p)	11.4	8.1	44.0
MP2/6-31+G(d,p) // HF/6-31G(d)	11.8	8.2	43.2
PM3 // PM3	17.2	14.0	69.9
SVWN5/6-311+G(2d,p) // SVWN5/6-311+G(2d,p)	18.1	19.8	81.0
AM1 // AM1	18.8	16.9	95.5
SVWN/6-311+G(2d,p) // SVWN/6-311+G(2d,p)	24.9	19.2	89.3
HF/6-311+G(2d,p) // HF/6-31G(d)	46.1	40.0	173.8
HF/6-311+G(2d,p) // B3LYP/6-31G(d)	46.6	40.5	174.6
HF/6-31+G(d,p) // HF/6-31G(d)	46.6	40.7	179.9
HF/6-31+G(d,p) // HF/6-31+G(d,p)	46.7	40.6	179.8
HF/6-31+G(d,p) // AM1	49.4	43.1	206.1
HF/6-31G(d) // HF/6-31G(d)	51.0	41.2	184.2
HF/6-31G(d) // AM1	54.2	43.1	207.2
HF/3-21G(d) // HF/3-21G(d)	58.4	50.1	215.2
HF/STO-3G // HF/STO-3G	93.3	66.3	313.9

Sample Relative CPU Times			
Model	PH <sub>3</sub>	F <sub>2</sub> CO	SiF <sub>4</sub>
CBS-4	1.0	1.0	1.0
G2(MP2)	2.4	10.3	11.5
CBS-Q	2.8	8.4	12.7
G2	3.2	25.9	59.1

# Ozone Hole Problem



Method	O <sub>3</sub>	$\begin{matrix} D_0 \\ \text{O}_2 \end{matrix}$	ClO	$\Delta H$
HF/6-31G(d)	-14.2	26.9	-1.8	-39.3
MP2/6-31G(d)	101.1	115.3	44.6	-58.7
MP4/6-31G(d)	96.1	105.1	43.0	-52.0
B3LYP/6-31G(d)	138.9	122.1	57.9	-41.2
QCISD(T)/6-31G(d)	108.4	103.0	45.3	-39.9
HF/6-31+G(d)	-15.6	26.1	-2.0	-39.7
MP2/6-31+G(d)	100.1	113.6	45.3	-58.8
MP4/6-31+G(d)	95.3	103.5	43.9	-52.1
B3LYP/6-31+G(d)	133.7	118.1	57.3	-41.7
QCISD(T)/6-31+G(d)	89.7	101.1	46.5	-58.0
HF/6-311+G(3df)	-7.1	31.1	4.9	-43.1
MP2/6-311+G(3df)	120.2	124.9	58.0	-62.7
MP4/6-311+G(3df)	117.2	117.3	56.7	-56.8
B3LYP/6-311+G(3df)	138.8	121.5	65.3	-47.9
QCISD(T)/6-311+G(3df)	127.3	113.5	58.6	-44.8
Experiment	142.2	118.0	63.3	-39.1

# When is Hartree Fock Bad

- WHEN electronic state can not be described by one Slater determinant (configuration)  
Hartree Fock is bad
1. Metal systems with small HOMO LUMO gap
  2. Bond dissociation
  3. Transition state
  4. Intermolecular interaction

# Small Rules to Remember

- Equilibrium geometry present day state of art calculation is from CCSD(T)
- Transition metals you should use multireference methods, d orbitals are hard!
- If you want to correctly write the potential energy curve of dissociating problems must use multireference methods
- Van de waals interaction must use MP2, CCSD CI, MRCI methods to accurately get
- You have to do test calculations for your problem to find the best way to do your problem

# Quantities other than Energy

- Dipole Moment of the Molecule: Rotational Spectroscopy
- Dipole Derivative of the Molecule: Infrared Absorption Spectroscopy
- Polarizability Derivative of the Molecule: Infrared Raman Spectroscopy
- Nuclear Shielding by the electron related to NMR spectrum