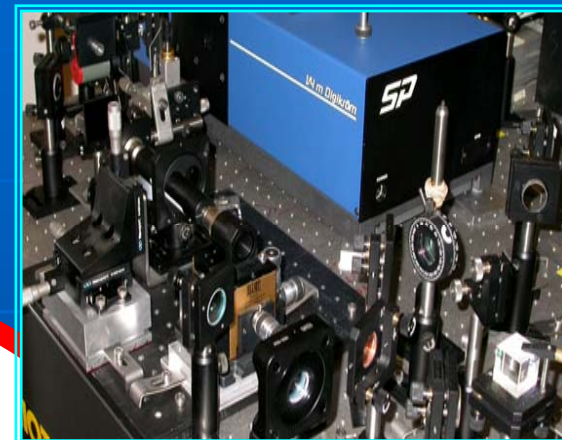
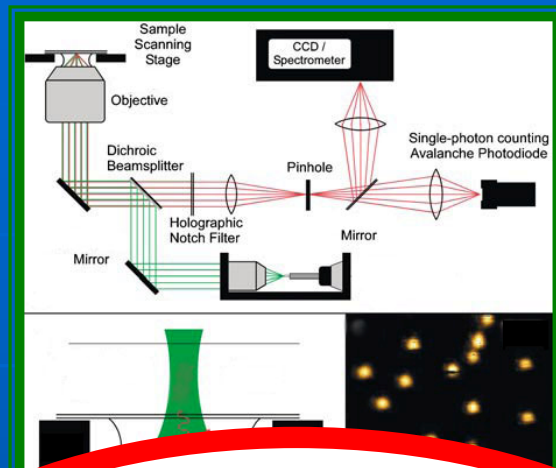
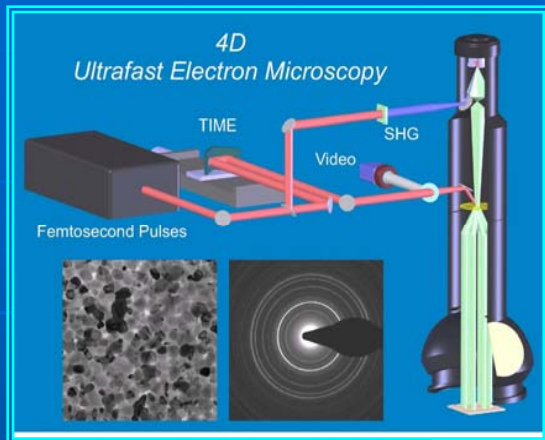


**Studies of light-induced stochastic and coherent
processes in nano-structured materials
using single-molecule techniques and
femtosecond time-resolved electron diffraction**

Jau Tang 湯朝暉

*Research Center for Applied Sciences,
Academia Sinica, Nankang, Taiwan
and*

*Department of Photonics
NCTU, Hsinchu, Taiwan*



Quantum dots, Nanorods,
Biomolecules,
Metallic & Semiconductor films & NCs
Organic conductors, photovoltaics
Light-induced phase transitions



Structure
Dynamics
Functionality



Outline

Time-resolved electron diffraction and microscopy

Impulse-induced structural dynamics & its effects on diffraction patterns

Coherent acoustic wave excitation

Nano-scale heat transfer of laser-heated metals

Light-induced insulator-metal phase transitions

Confocal microscopy and single-molecule techniques

Fluorescence intermittency of illuminated nano-particles

Photo-induced charge transfer

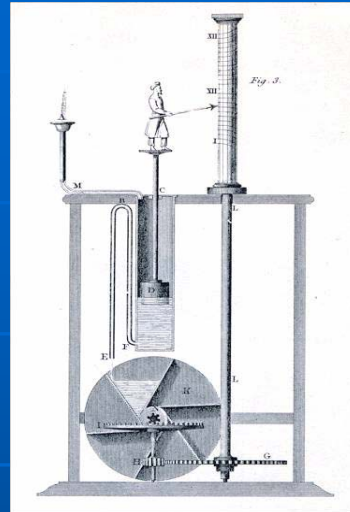
Enzymatic reactions and intermittency of single biomolecules

Förster's fluorescence resonance energy transfer

water clock (1500 BC)



Egyptian water clock (300 BC)



NIST F-1, (cesium fountain clock)



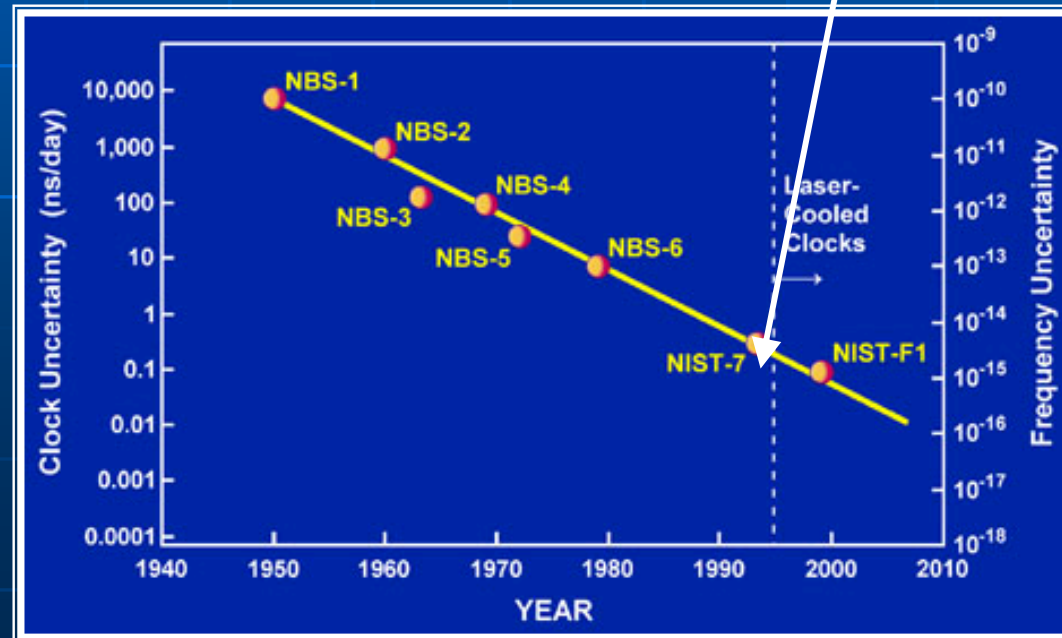
NIST F-1
Atomic Clock

accuracy: 10^{-15}

Sun dial clock (300 BC)



Pendulum clock (1500 AD)



Time arrow



Maser – 1954

CW Laser – 1959

Nanosecond laser – 1960s

Picosecond laser – 1970s

Femtosecond laser – 1980s

A macro, micro and nano view of a human hand, 10 times magnification per frame.



Janssen, Leeuwenhoek
(16th century)

Macro, Micro, Nano

How small is a nanometer? Stepping down in size by powers of 10 takes you from the back of a hand to, at one nanometer, a view of atoms in the building blocks of DNA. The edge of each image denotes a length 10 times longer than its next smallest neighbor. The black square frames the size of the next scene inward.

HAND



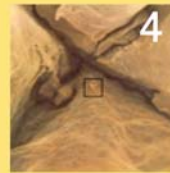
10 centimeters



1 centimeter



1 millimeter



100 microns



Knott, Ruska
(1930)

(magnifying glasses
invented 1st century)



WHITE BLOOD CELL



10 microns



1 micron

DNA



100 nanometers



10 nanometers



1 nanometer

From the classic book *Powers of Ten*,
by Philip and Phylis Morrison and the
office of Charles and Ray Eames.

maximum theoretical resolution
~ 1/2 wavelength

radios ~ miles to mm
infrared ~ mm to 800nm
light ~ 400-800nm
electrons ~ 20kV 0.008 nm

under the best of conditions:

human eye ~ 0.2 mm

light scope ~ 0.2 μ m

SEM ~ 3 nm

TEM ~ 0.2 nm

◆ **Electron Microscopes were not developed until the Twentieth Century**

◆ **The first Transmission Electron Microscope (TEM) was built in 1932**



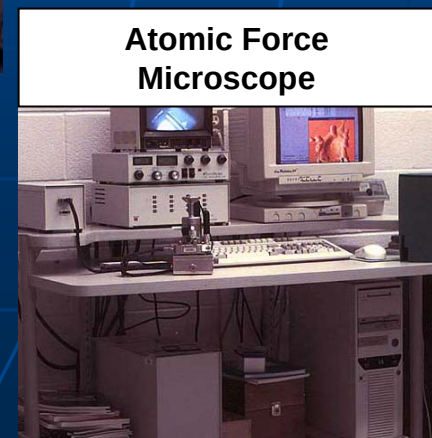
◆ **The first Scanning Electron Microscope (SEM) was built in 1942.**



◆ **The Scanning Tunnelling Microscope was developed in 1982.**

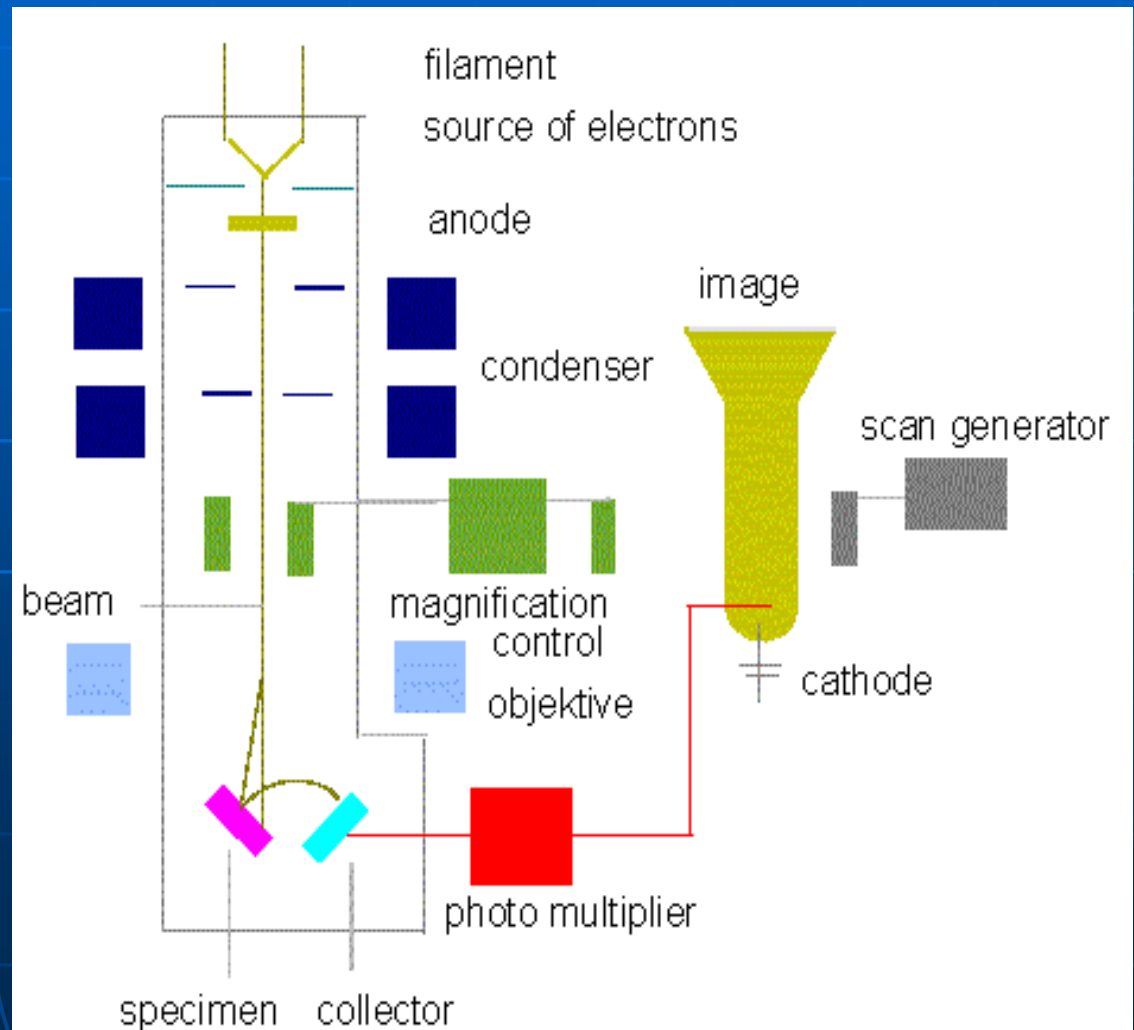
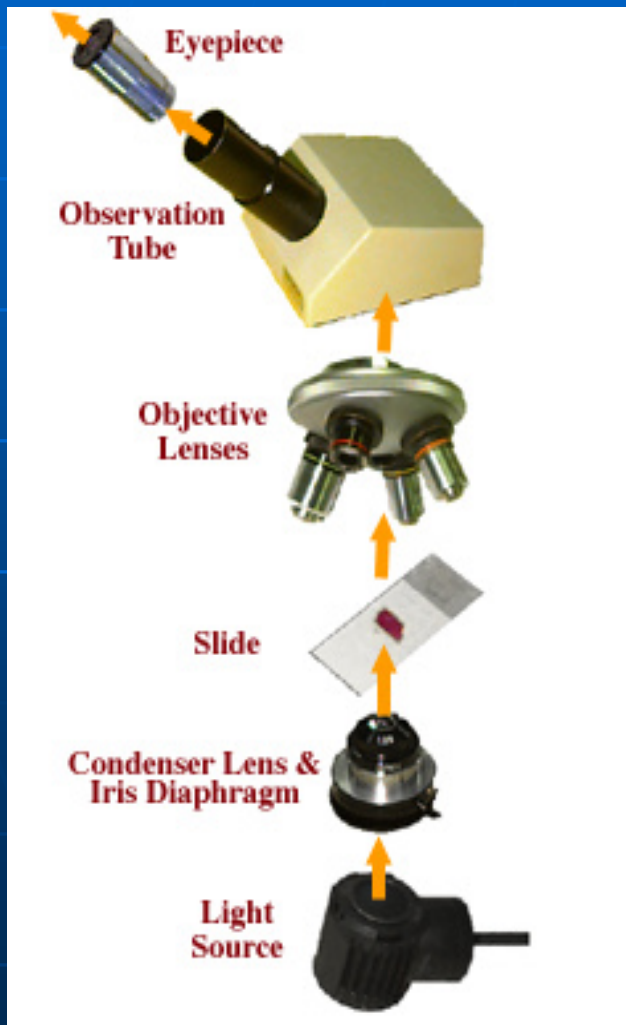


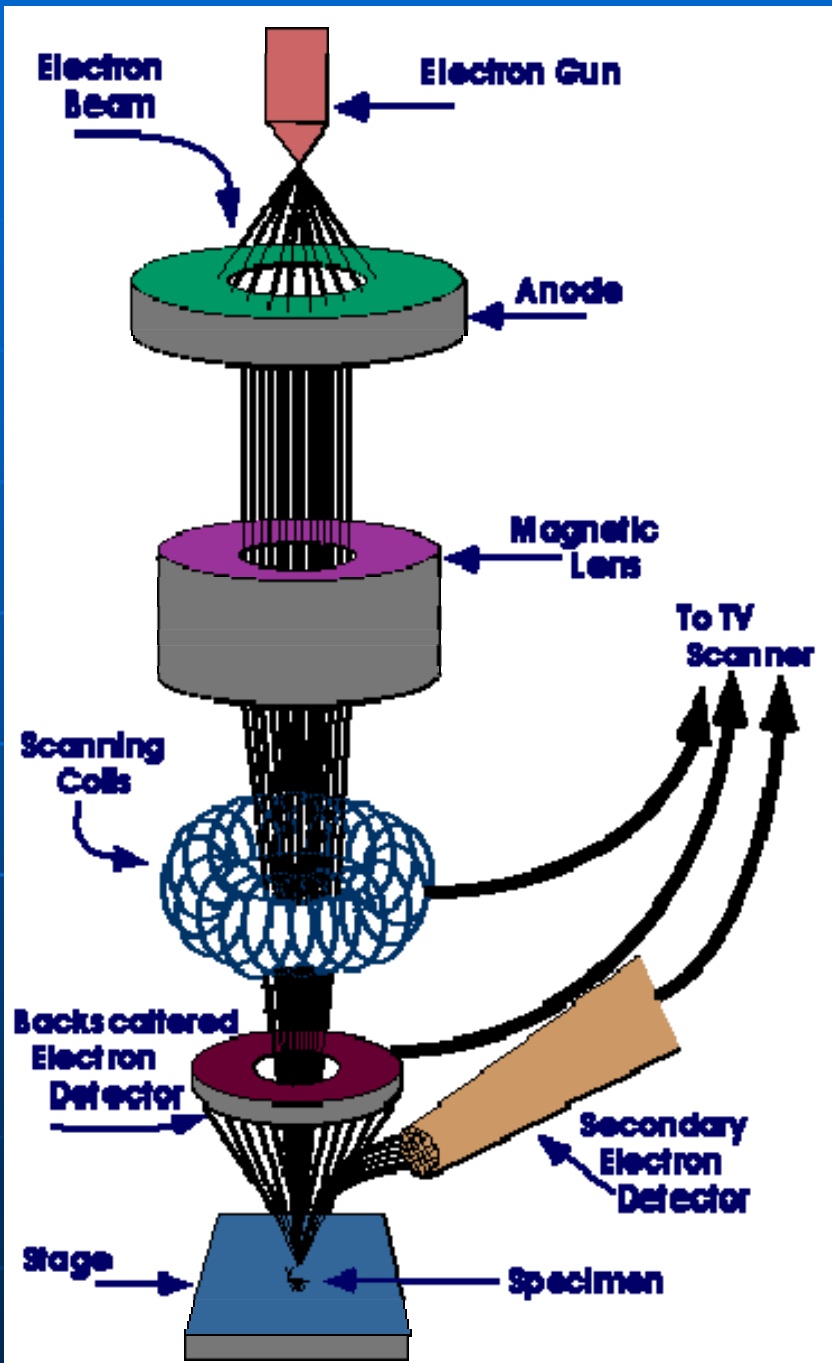
◆ **The Atomic Force Microscope was developed in 1985.**



- Electrons can be focussed using magnetic lenses
- Shorter wavelength than visible light

SEM

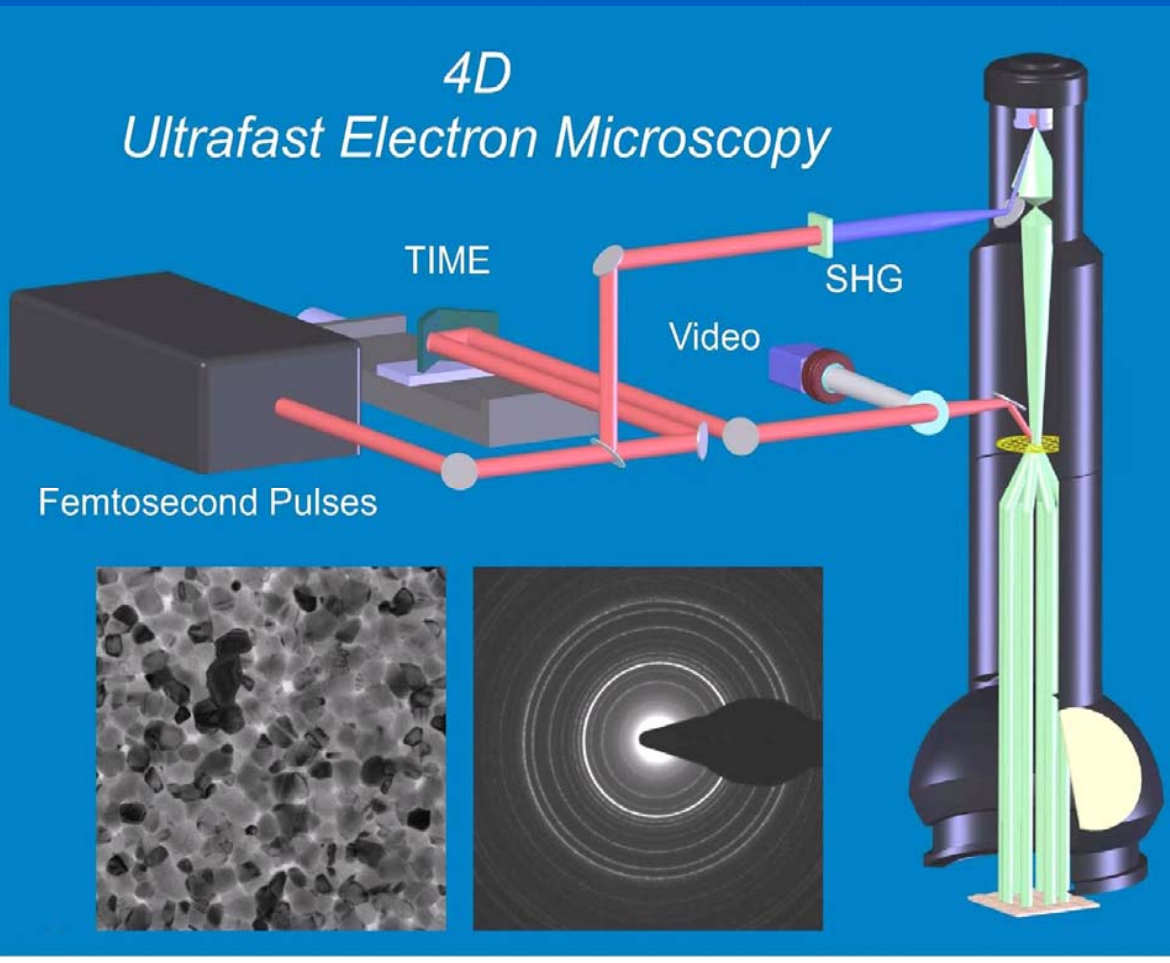




SEM

- **Large depth of field**
More of the sample is in focus at one time
- **Higher resolution**
Smaller features can be imaged
- **Analysis**
The electron beam interacts with the sample enabling information on composition to be collected using additional detectors

Studies of structural dynamics by Ultrafast Electron Diffraction and Microscopy



pulsed X-ray diffraction

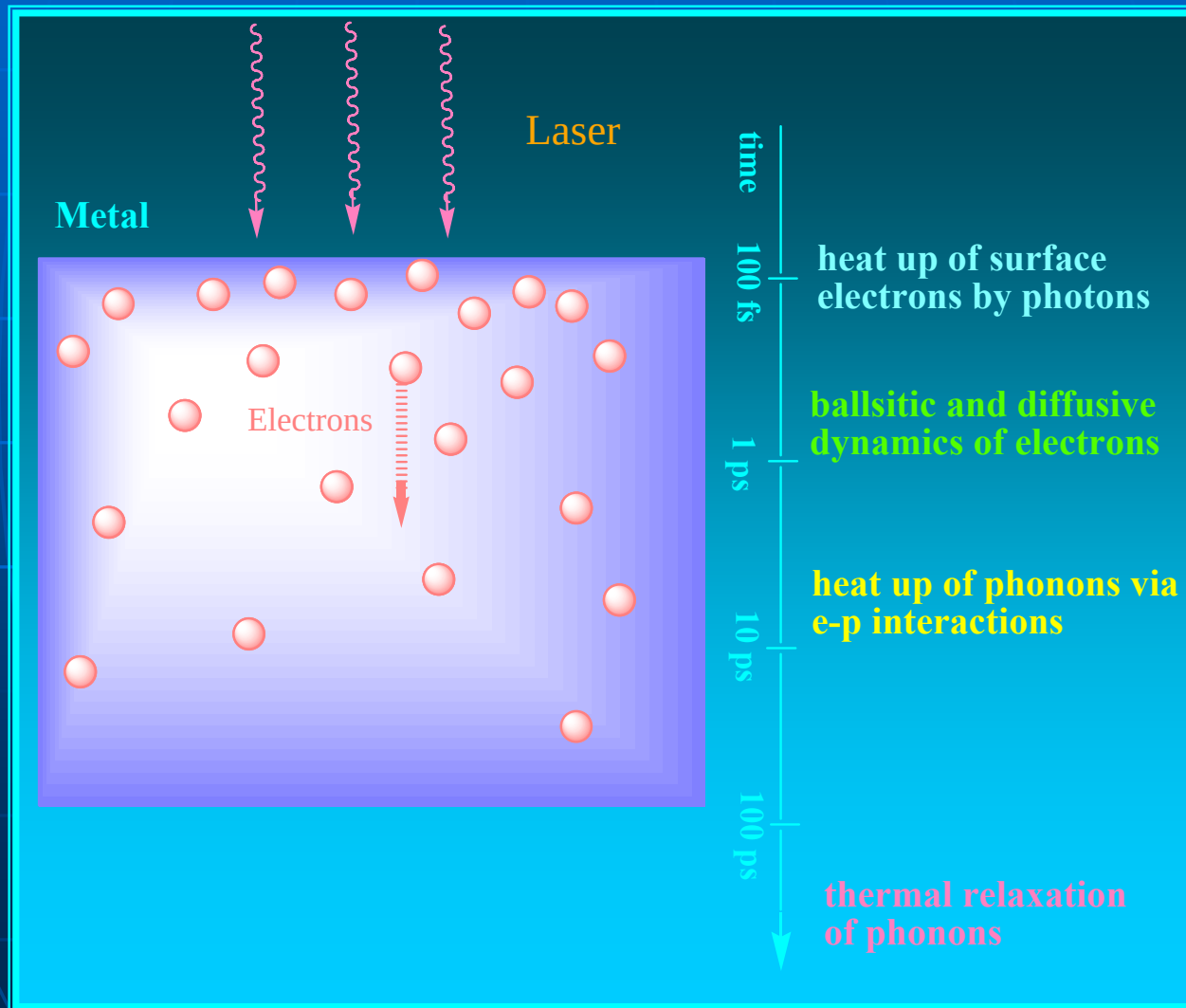
- High synchrotron
- Small scatter cross section
- Deep penetration depth
- Thick and large sample
- Bulk materials

pulsed electron diffraction

- Desk top operation
- Larger (10^6) scatter cross section
- Shallow penetration depth
- Thin and small samples
- nano-structured materials and surface science

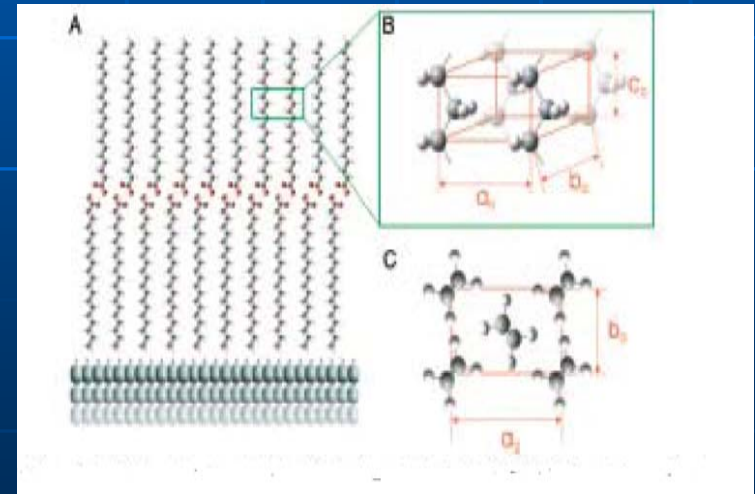
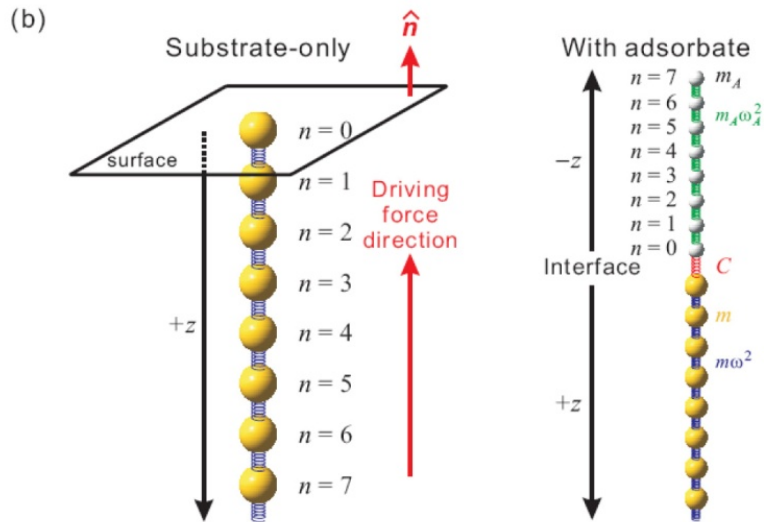
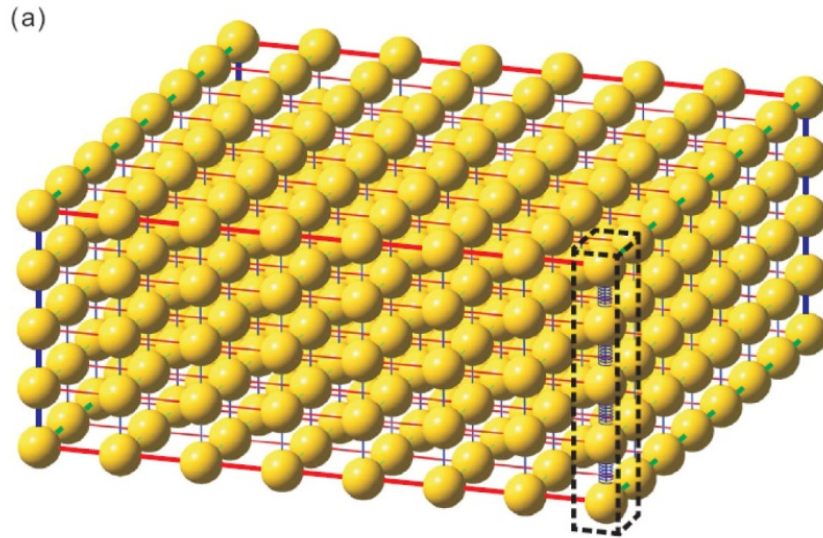
Ultrafast Electron Crystallography of Laser Heated Metals

J. Tang, D.-S. Yang and A. H. Zewail, J. Phys. Chem. C, (2007)



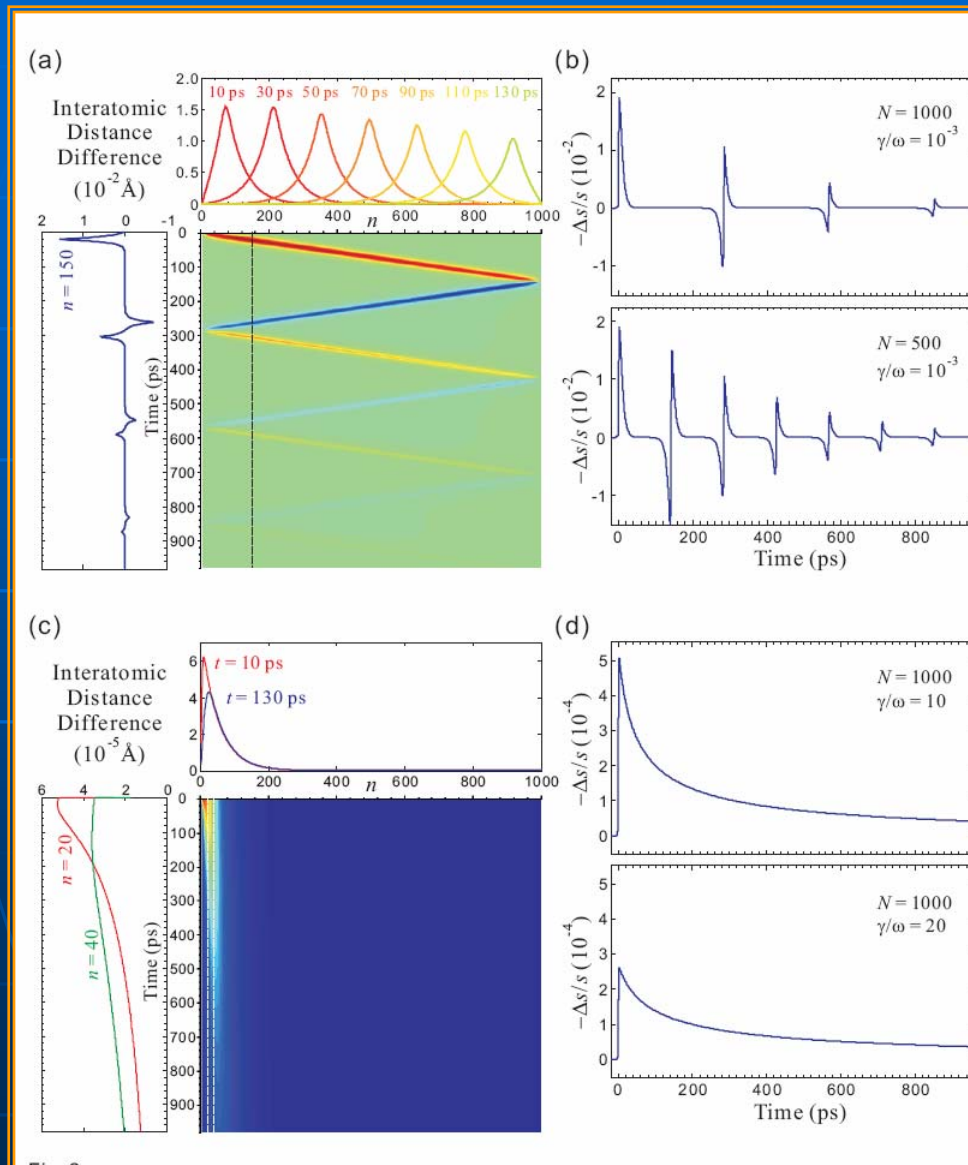
Fermi-Pasta-Ulam model Anharmonic chain

$$H = \sum_{k=0}^{N-1} \frac{p_k^2}{2m_k} + \sum_{k=0}^{N-2} \frac{m_k \omega_k^2}{2} (q_k - q_{k+1})^2 + \sum_{k=0}^{N-2} \alpha_k \frac{m_k \omega_k^2}{3} (q_k - q_{k+1})^3$$



UEC of Gold Film

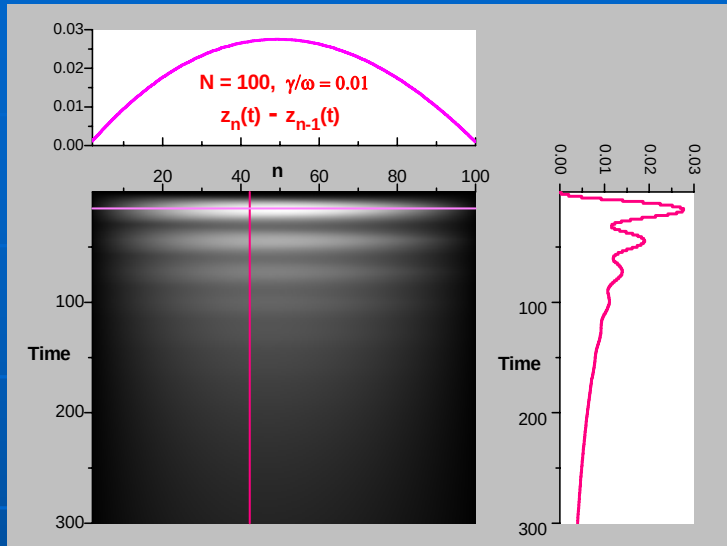
Shallow impulse, $N = 1000$



shallow impulse



Inter-atomic spacing



UEC of a thin gold film

$N = 100$

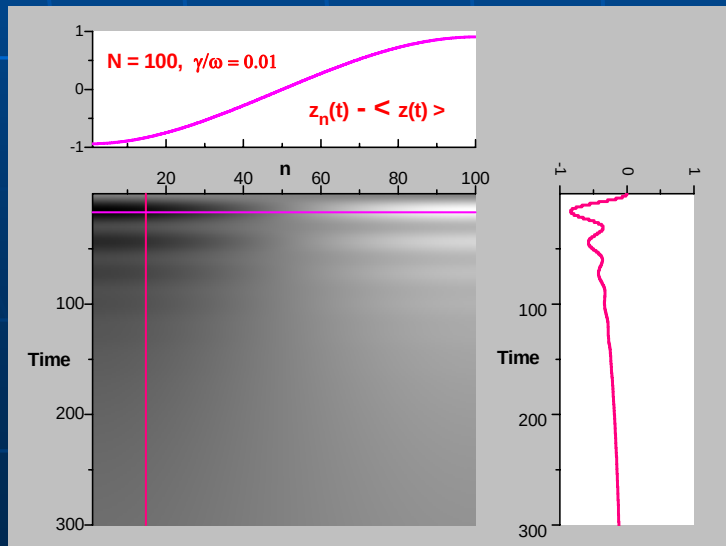
Deep impulse

“breathing” phenomenon

Wavelength / 2 = slab thickness

Period / 2 = thickness / sound speed

Position shift w.r.t. center of mass



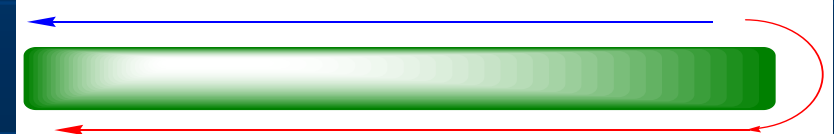
Largest expansion in the middle

constructive interference



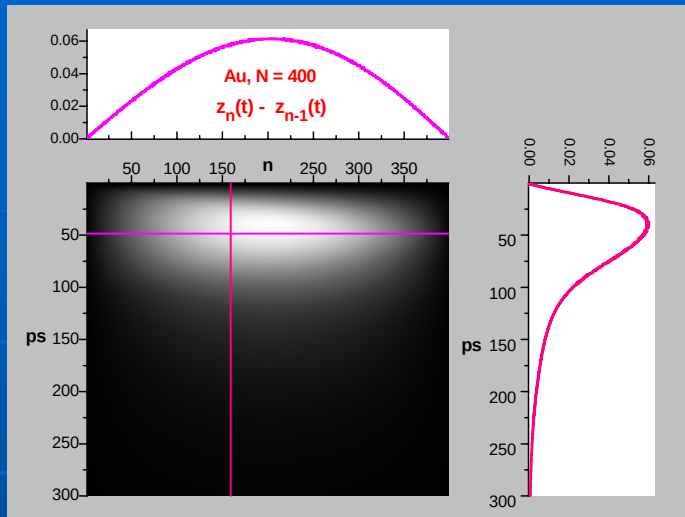
Smallest expansion on both surfaces

destructive interference



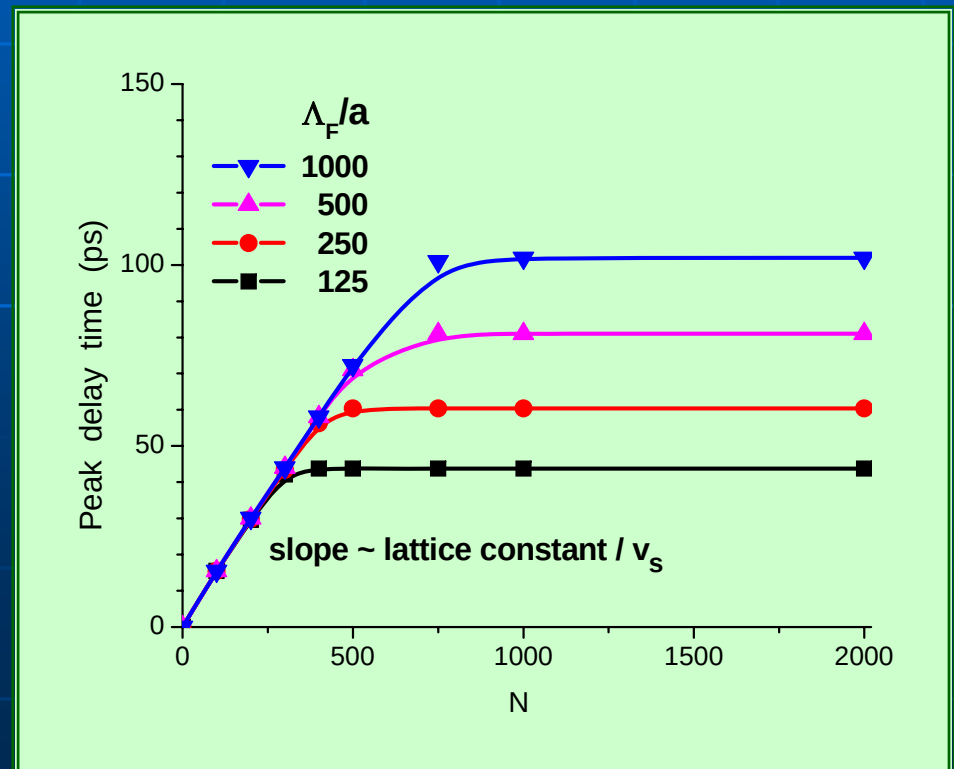
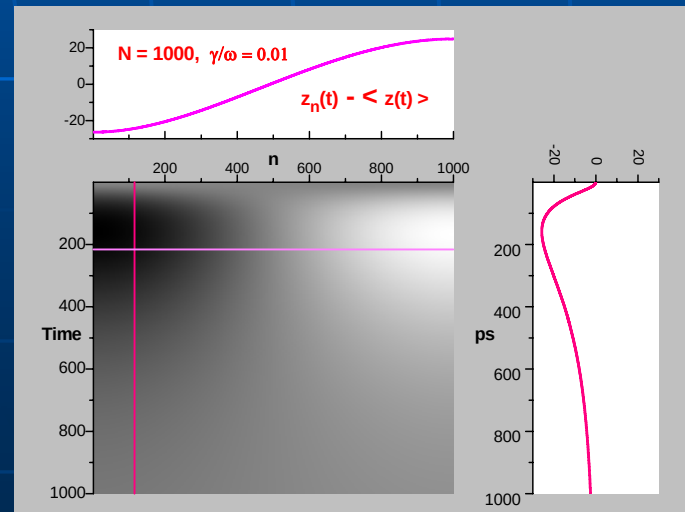
UEC of Gold Film $N = 400$

Inter-atomic spacing

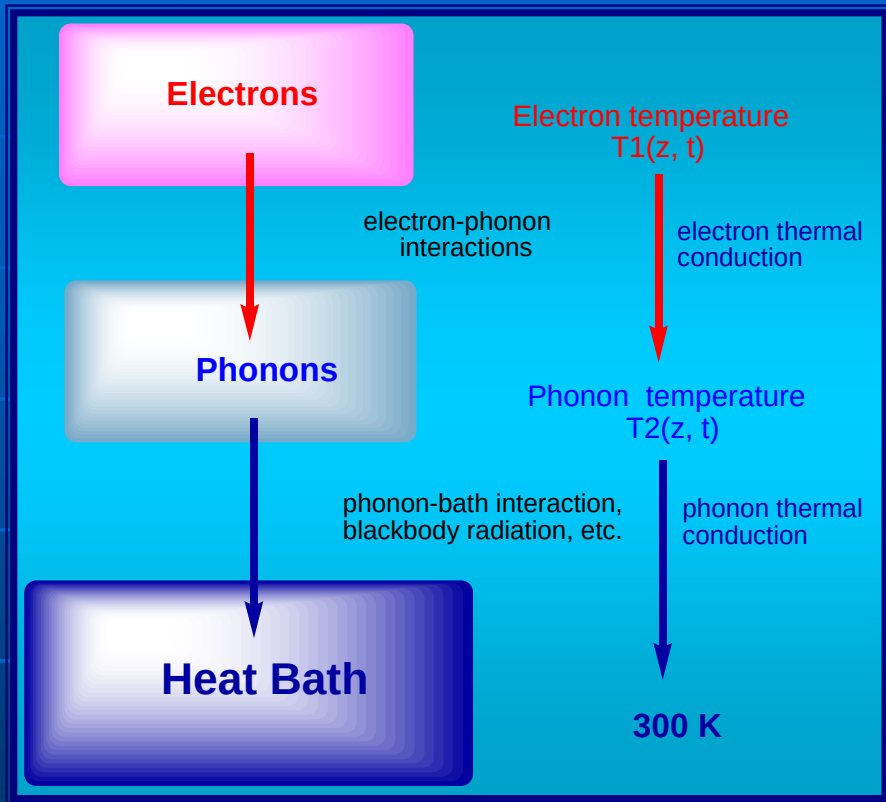


Dependence of the peak delay time on the thickness of a substrate

Position shift w.r.t. center of mass



Laser heating of metals and two-temperature model



$$\frac{\partial}{\partial t}[C_e(T_1)T_1(z,t)] = \frac{\partial}{\partial z}\left(\kappa_e(T_1)\frac{\partial}{\partial z}T_1(z,t)\right) - g(T_1(z,t) - T_2(z,t)) + S(z,t)$$

$$C_2 \frac{\partial}{\partial t}T_2(z,t) = \kappa_L \frac{\partial^2}{\partial z^2}T_2(z,t) - g(T_2(z,t) - T_1(z,t))$$

Temperature-induced impulsive forces

$$F_k(t) = -\gamma_{2,G}C_2(T_{2,k}(t) - T_0)\ell^2 - \gamma_{1,G}C_1T_{1,k}(t)(T_{1,k}(t) - T_0)\ell^2$$

Spatial and temporal dependence of electron and phonon temperatures

→ Impulsive force on each atom →

Transient atomic Displacement & Bragg peak shift

Aluminum: $\gamma_{2,G} \sim 2.2$ phonons
 $\gamma_{1,G} \sim 1.6$ electrons

Heat capacity of electrons $\sim T$

Fermi electron gas model

$$C_e(T) \sim (\pi^2 N k_B / 3 T_F) T, \quad T_F = E_F / k_B$$

Fermi temperature: Au - 6.4×10^4 , Al - 13.6×10^4 , Cu - 8.2×10^4

Heat capacity of phonons $\sim \text{constant in } T$

$$C = \frac{\partial}{\partial T} \sum_k \frac{\hbar \omega_k}{\exp(\hbar \omega_k / k_B T) - 1} \approx 3 N_A k_B$$

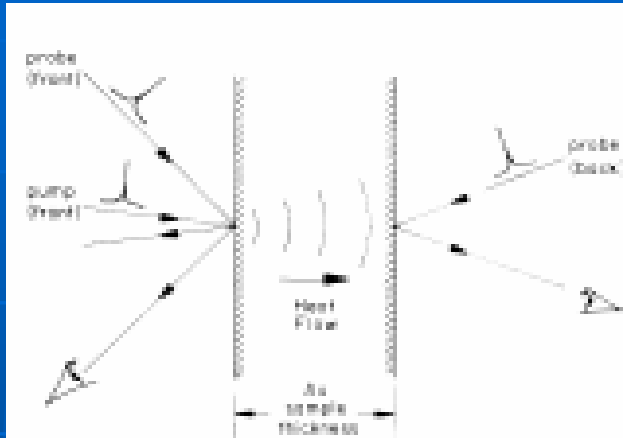
Dulong-Petit value

Electronic thermal conductivity $\sim T$

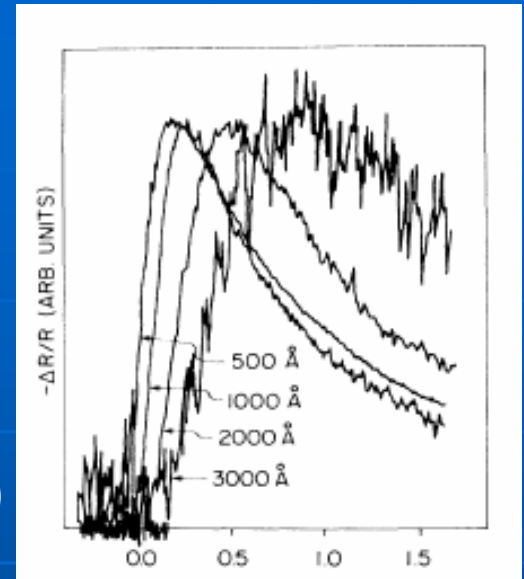
$$\kappa_e = \frac{\tau^2 k_B^2 n_e \tau}{3 m_e} T, \quad \tau \text{ scattering relaxation time}$$

$$\Lambda = v_F \tau \quad \text{mean free path}$$

Transient optical reflectivity of a gold film



(rear surface)

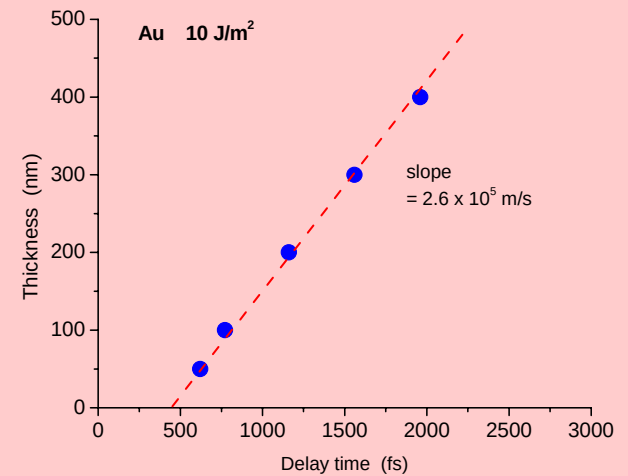
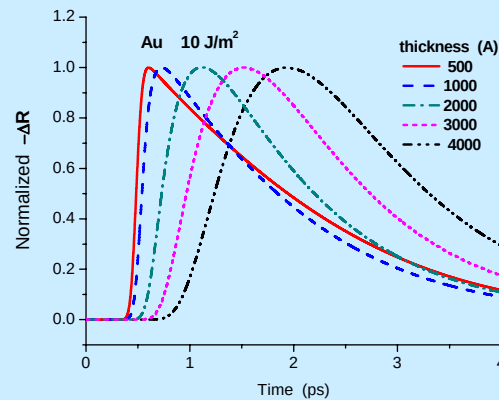
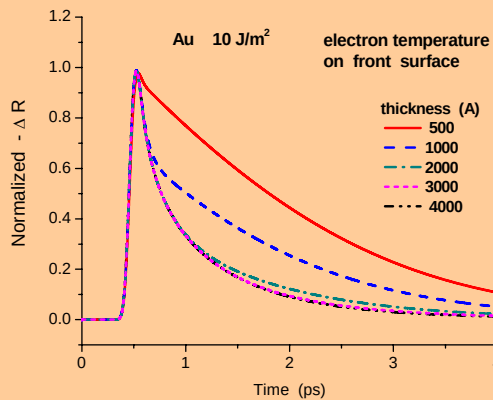


E. P. Ippen et al. Phys. Rev. Lett. 59, 1962, (1987)

Supersonic hot electrons

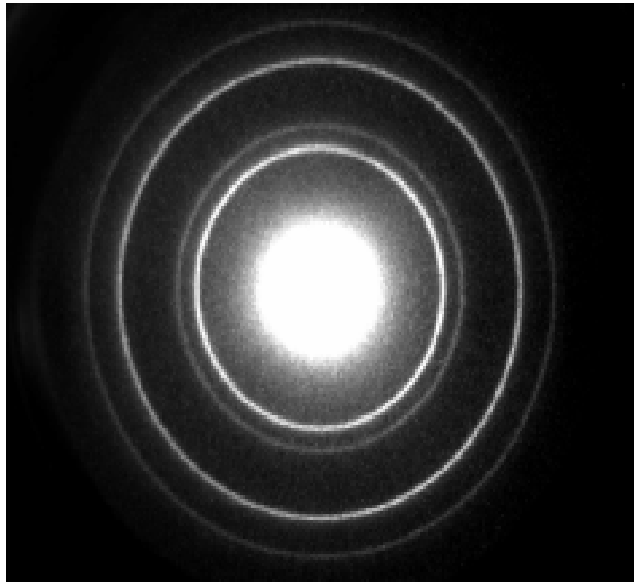
TTM (front surface)

(rear surface)

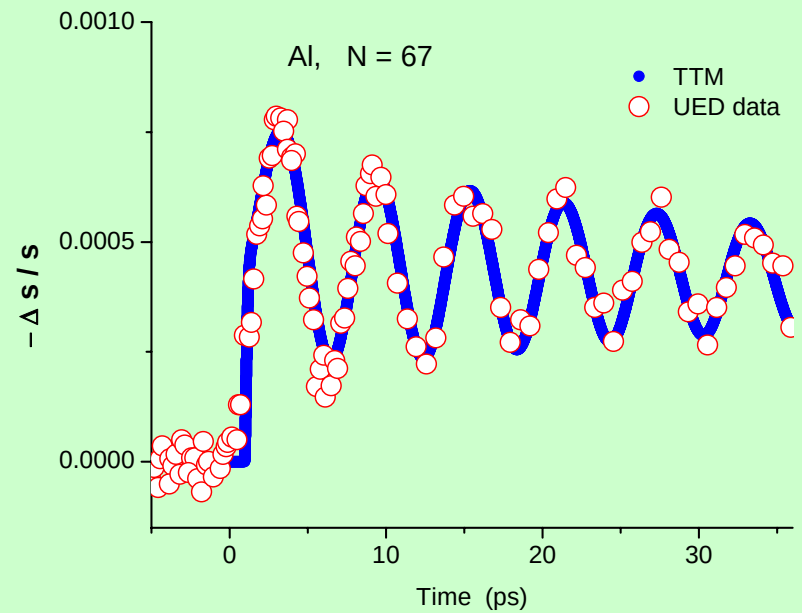
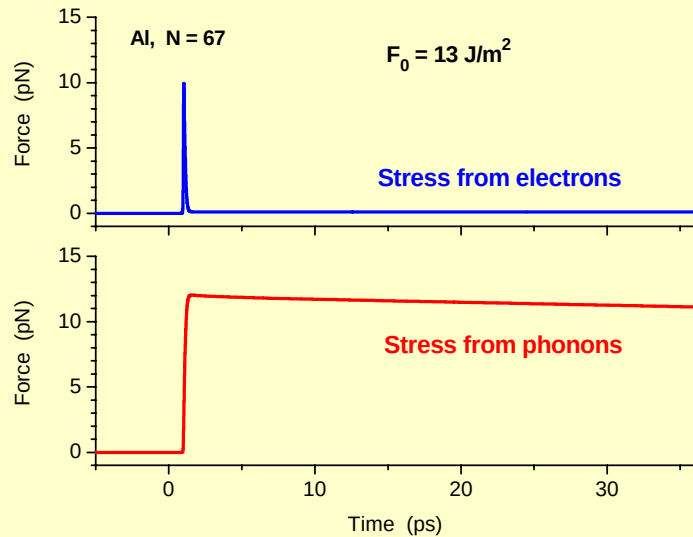


UEC of Aluminum Film

J. Cao et al. Phys Rev. Lett. (2006)
(20 nm thickness, fluence 13 J/m²)



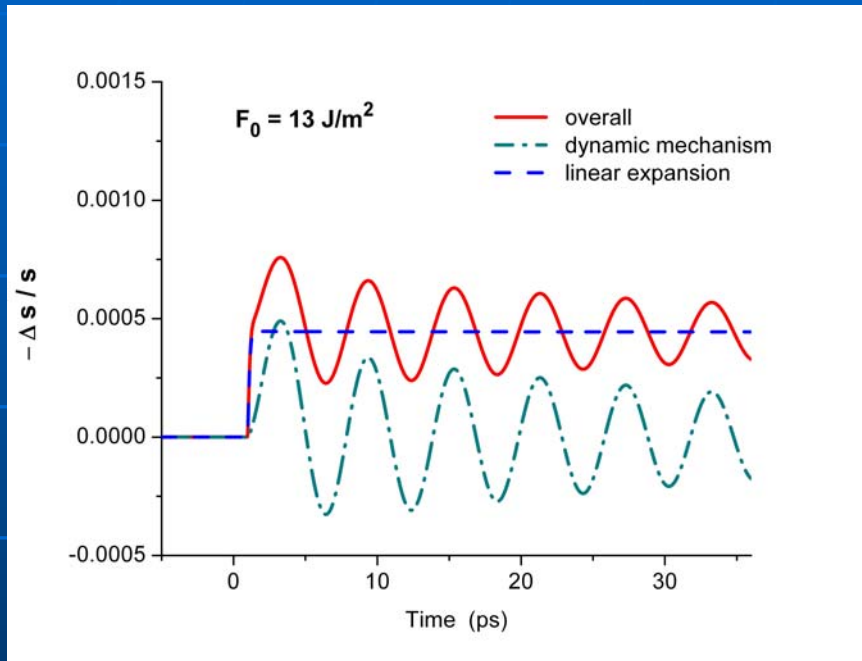
Acoustic wave excitation



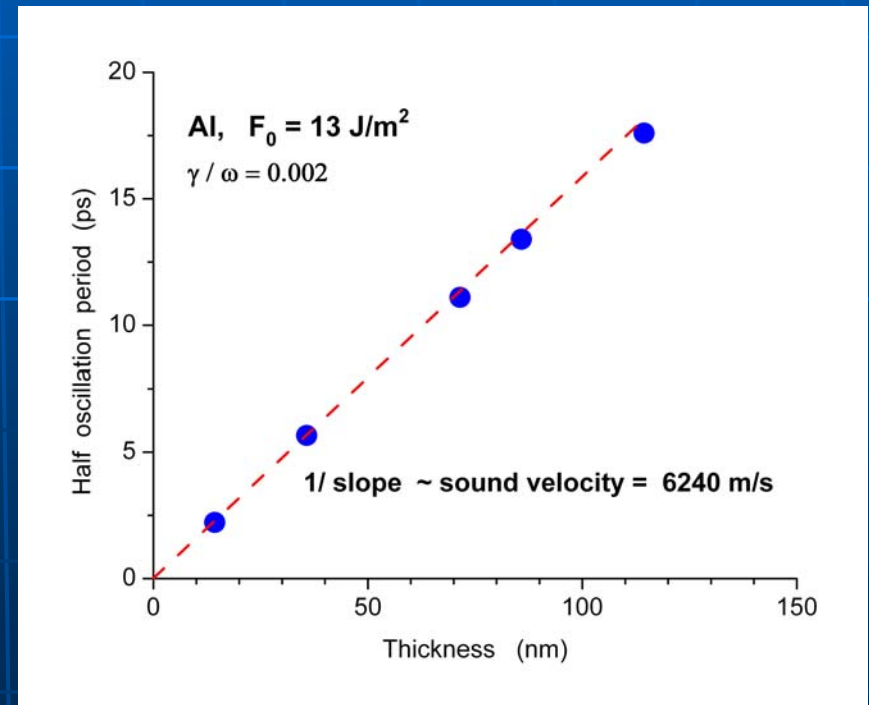
J. Tang, PRL (submitted)

Dynamic mechanism

Static linear thermal expansion

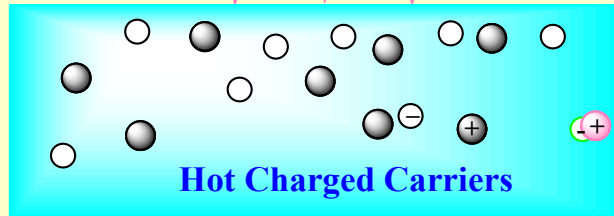


Half oscillation period vs. thickness



Production and heating of
electrons and holes
by photons

Laser



Auger
recombination

Energy transfer to
LO phonons ~ 0.1 ps

Energy transfer to
acoustic phonons ~ 0.5 ps

Optical Phonons

Energy transfer to
acoustic phonons ~ 10 ps

Acoustic Phonons

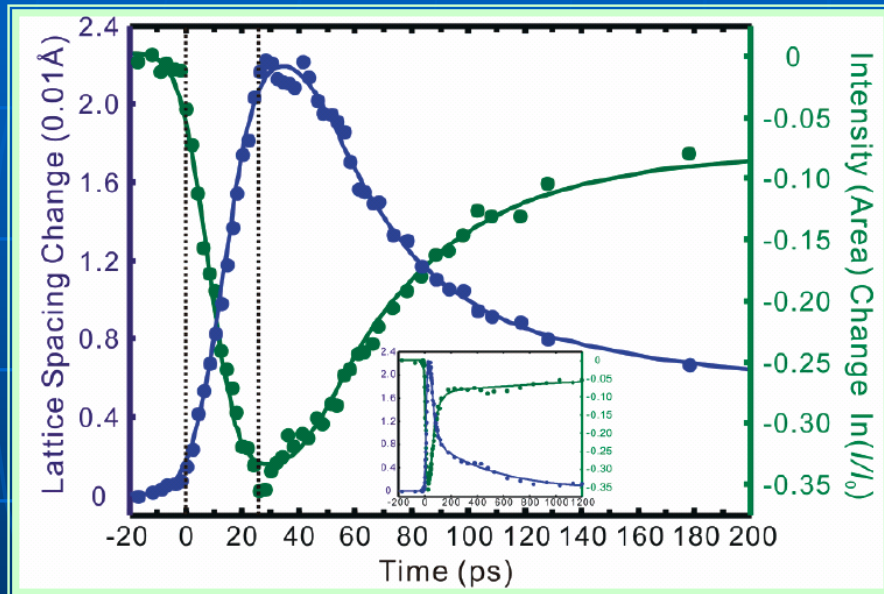
Laser heating of
Si & GaAs
semiconductors and
three-temperature model

Charged carriers

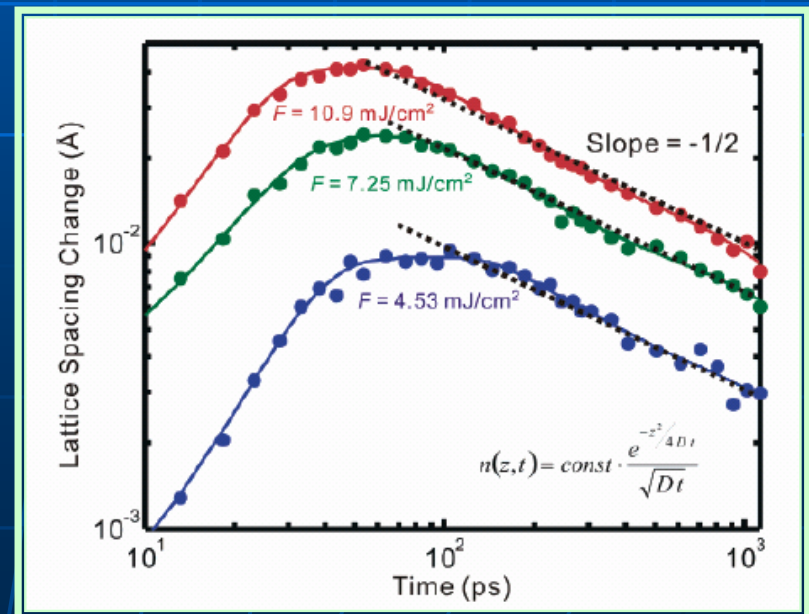
Optical phonons

Acoustic phonons

UEC of GaAs thin film



Yang, Gedik and Zewail, JPC (2007)



Dynamic vs. static mechanisms

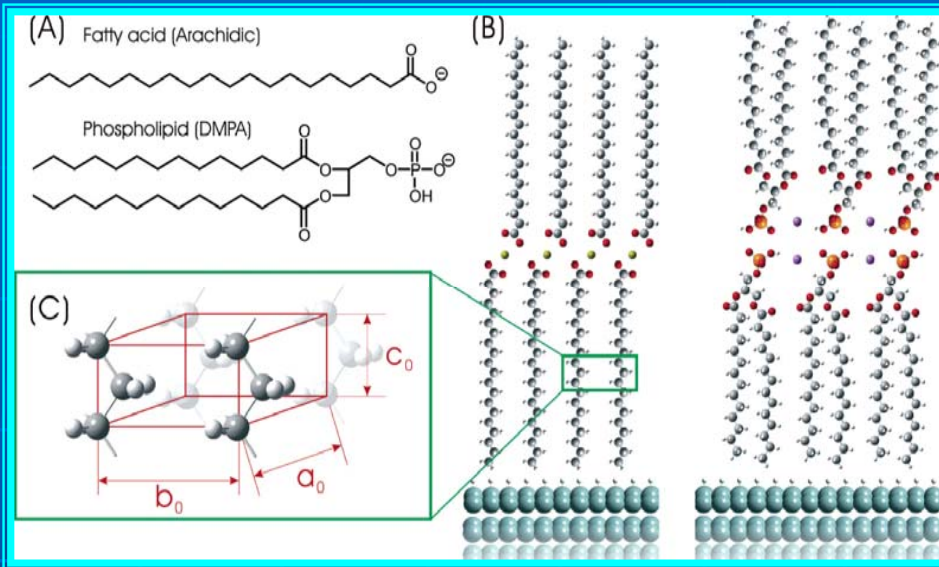
Dynamic

- Periodic expansion and contraction (thin film) or pure expansion (thick film)
- Non-local effect
- coherent (wave propagation)
- The magnitude of expansion is larger
- $\Delta s/s$ proportional to N (not saturated)
- $\Delta L/L \sim L^2$ (middle atoms expand more)

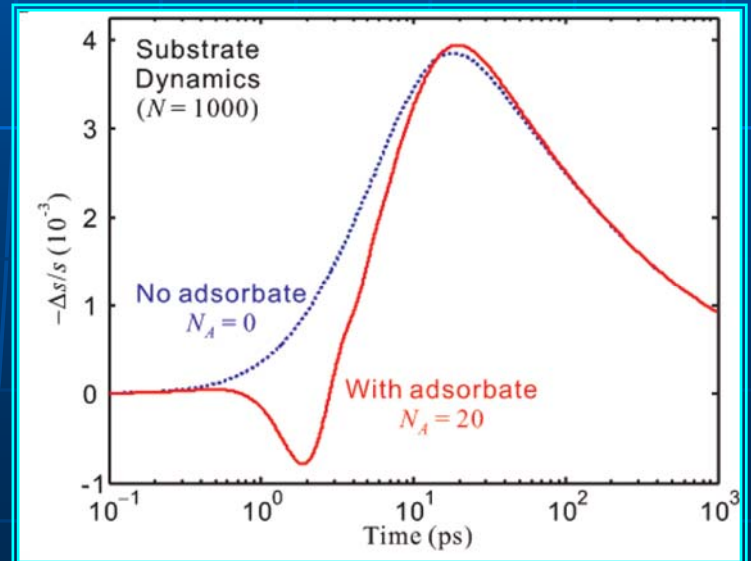
Static

- Only expansion occurs
- Local effect
- Non-coherent
- The magnitude of expansion is smaller
- $\Delta s/s$ independent of N
- $\Delta L/L \sim \Delta T$ independent of L

UEC of adsorbates & interfacial effects



Time dependence of
Bragg-spot shift



Laser-induced phase transition

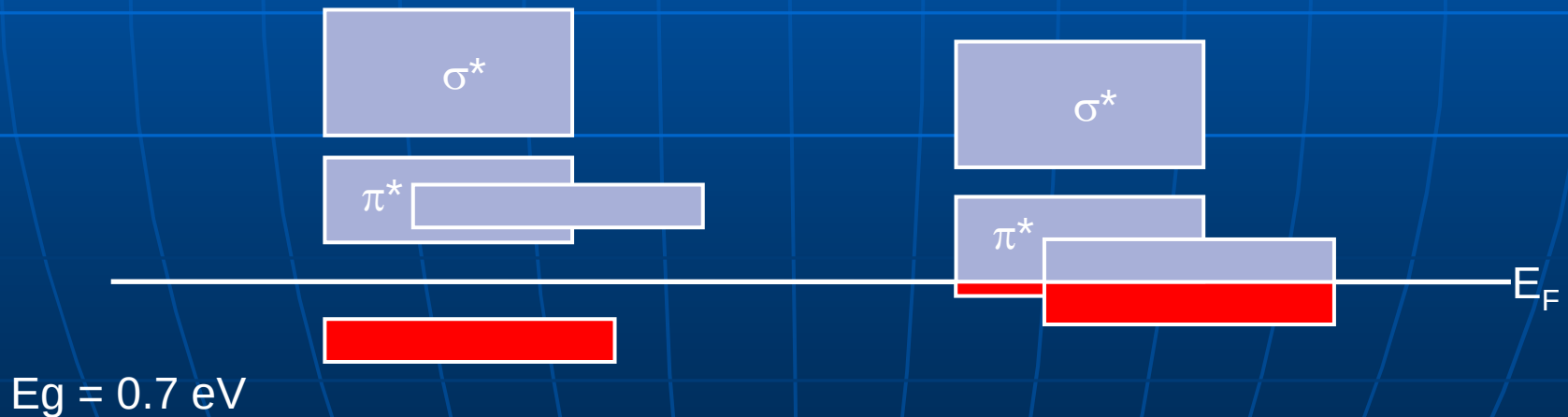
in VO_2

Insulator

$\leftrightarrow$

Metallic

$T_c \sim 340 \text{ K}$



Phase transition in VO_2

Insulator

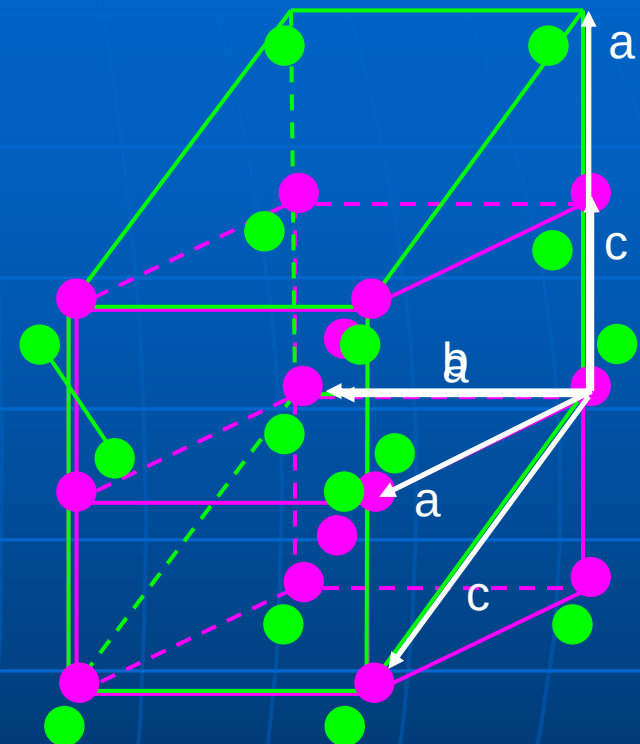


Metallic

Monoclinic

Tetragonal (Rutile)

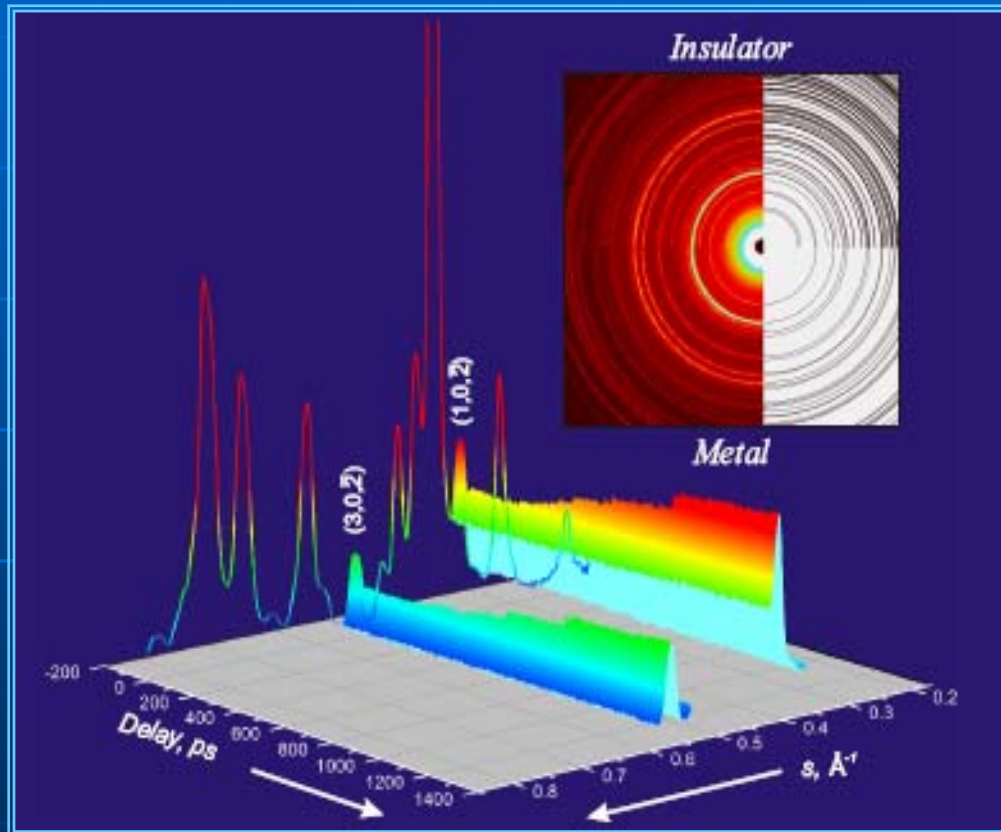
$$\begin{aligned} a_m &= 2c_r \\ b_m &= a_r \\ c_m &= a_r - c_r \end{aligned}$$



V-V bond

Electron diffraction pattern

Measured and simulated electron diffraction patterns



Electron gun: 120 keV
Laser : 776 nm
Fluence: 0.05 mJ/cm^2

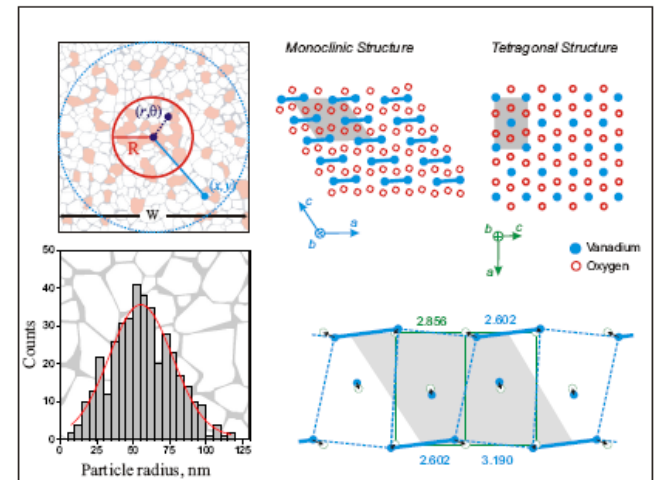
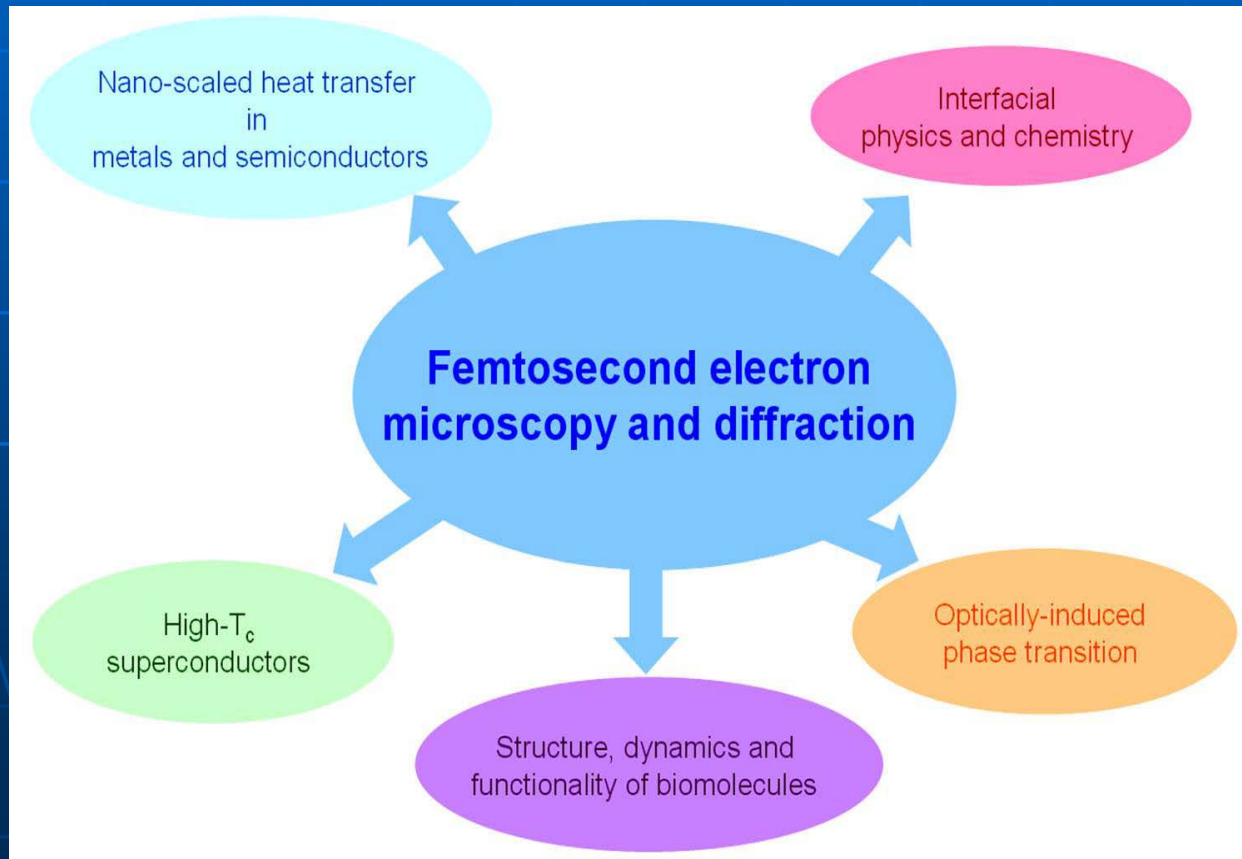
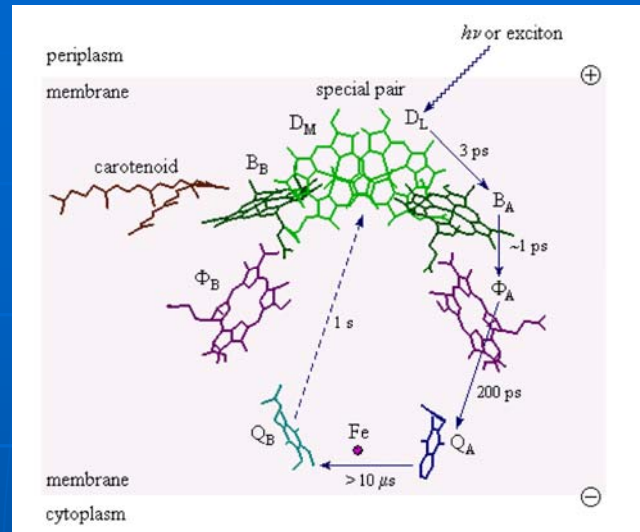


Figure 4. Lobastov, Weissenfelder, Tang & Zewail

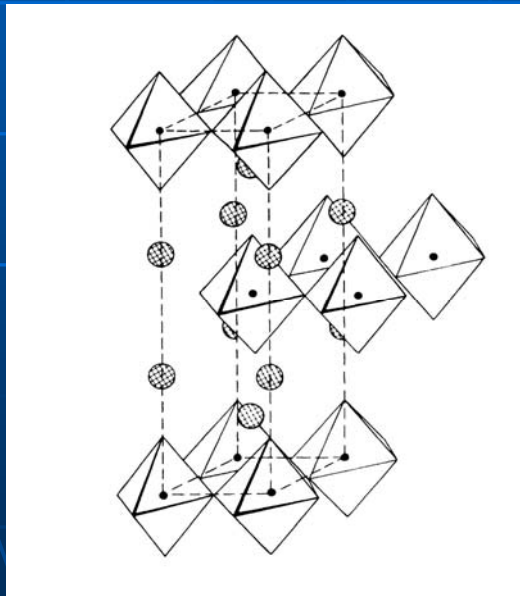
Natural Sciences: interfacial properties (surface physics, chemistry), nanoparticles, self-assembled monolayers, high T_c superconductors bond-breaking, conformation changes, etc.

Life Sciences: membrane protein, DNA, molecular recognition, hydration, charge transport, photosynthesis, etc.

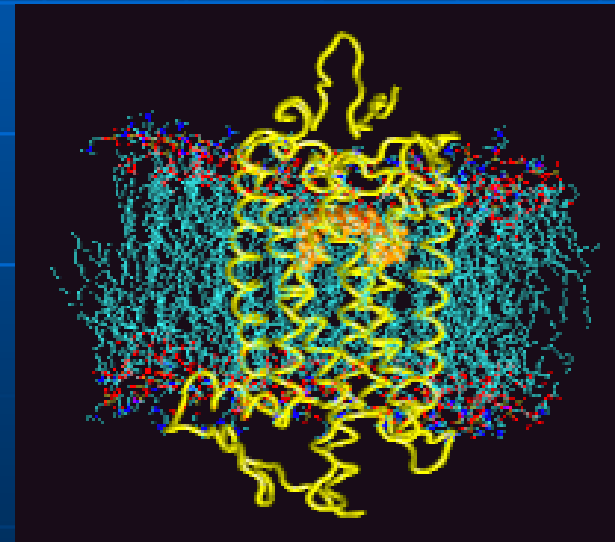




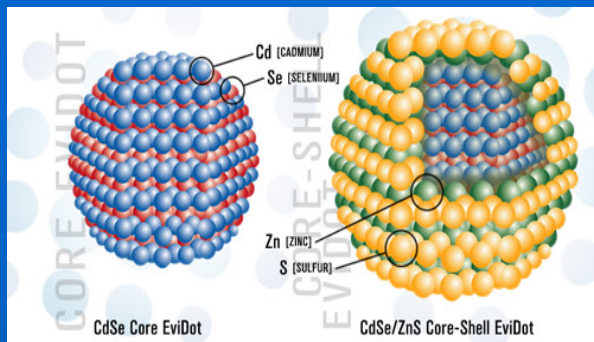
Photosynthetic RC



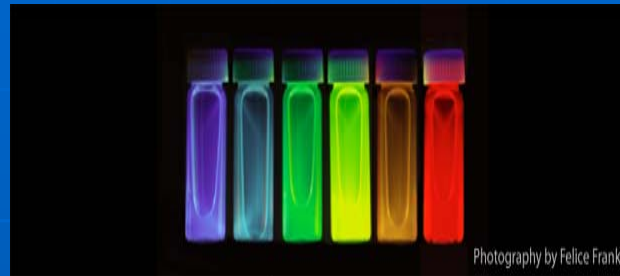
High T_c superconductor



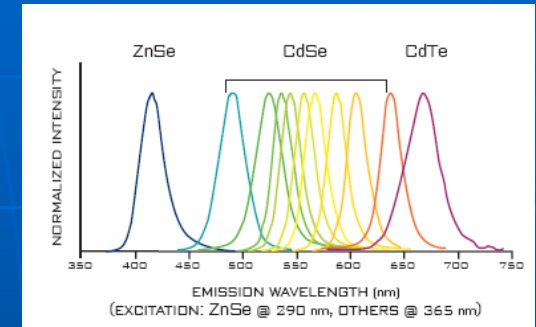
Rhodopsin



Quantum dots

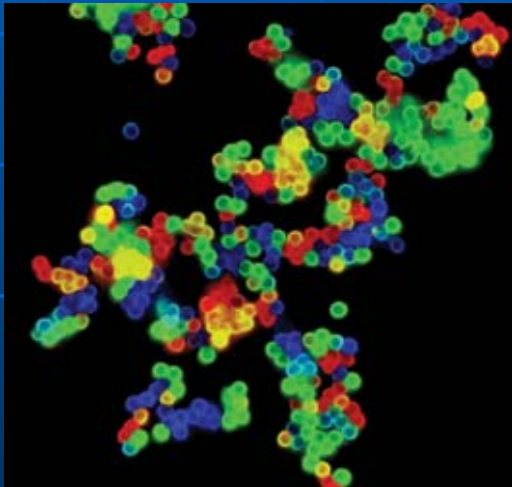


Emission spectra of QDs

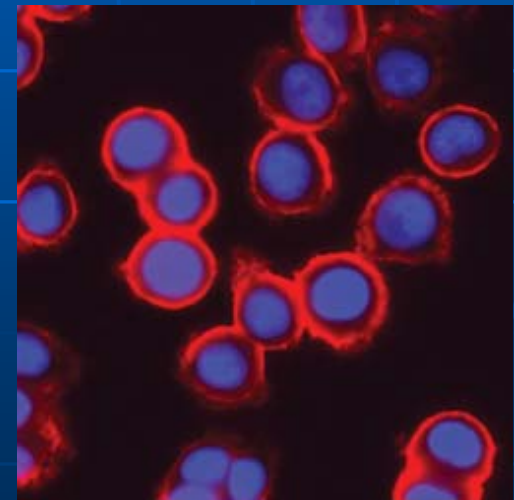
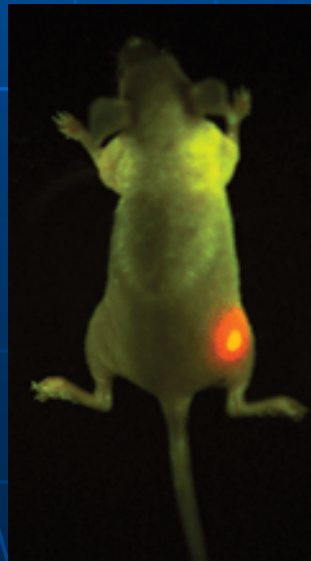


Polymer beads embedded with quantum dots fluoresce in five different colors.

labeling tags



Red QDs injected into a live mouse mark the location of a tumor.
C. Seydel, *Science*, **2003**, 300, 5616.



Red light-emitting quantum dots tag proteins on the surfaces of breast cancer cells.

Potential applications

Labeling tags

Solar cells

White LEDs

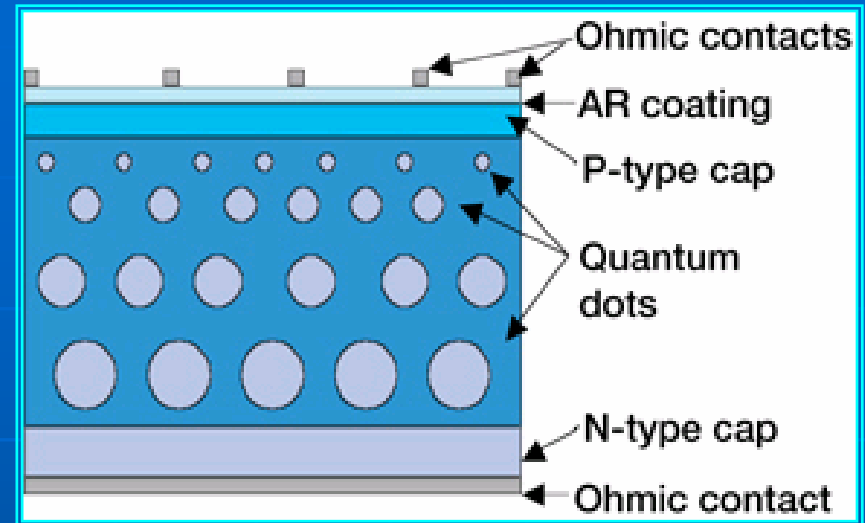
(LED Wavelength Converter)

Military Applications (Infrared Paints)

InAs/GaAs-based quantum dot laser

High density memory

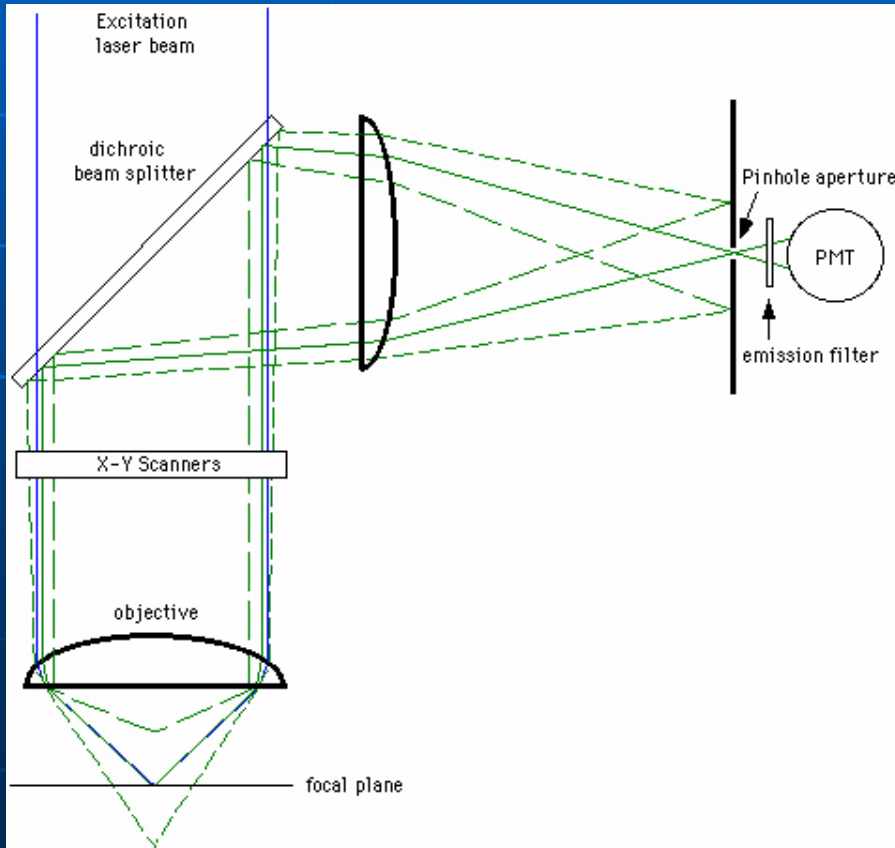
Quantum computing



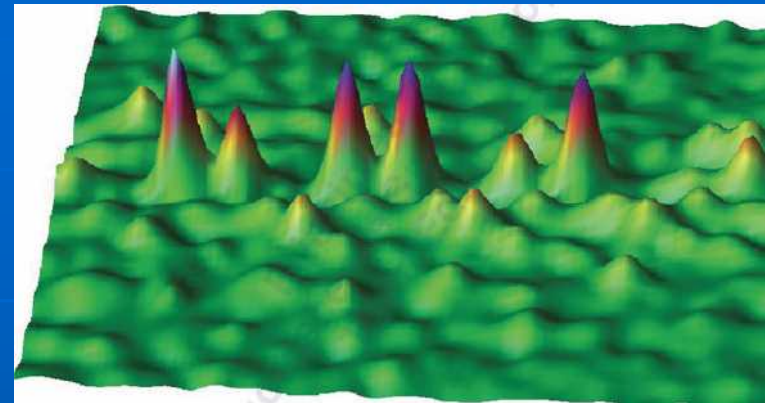
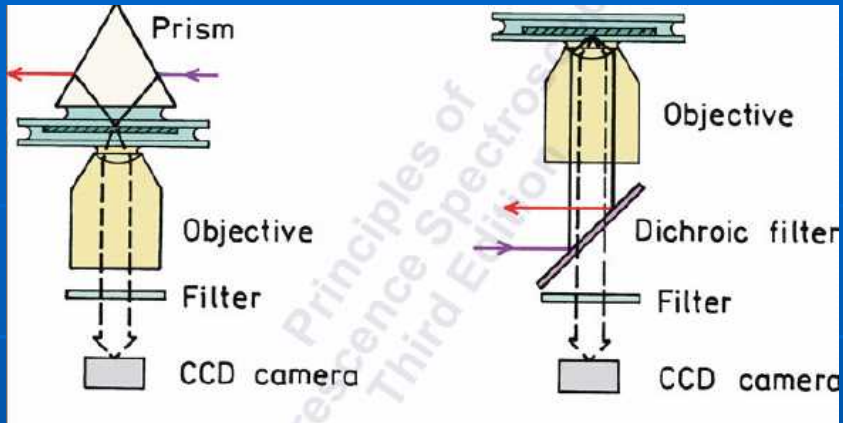
Intermediate-bandgap solar cell

Confocal microscopy

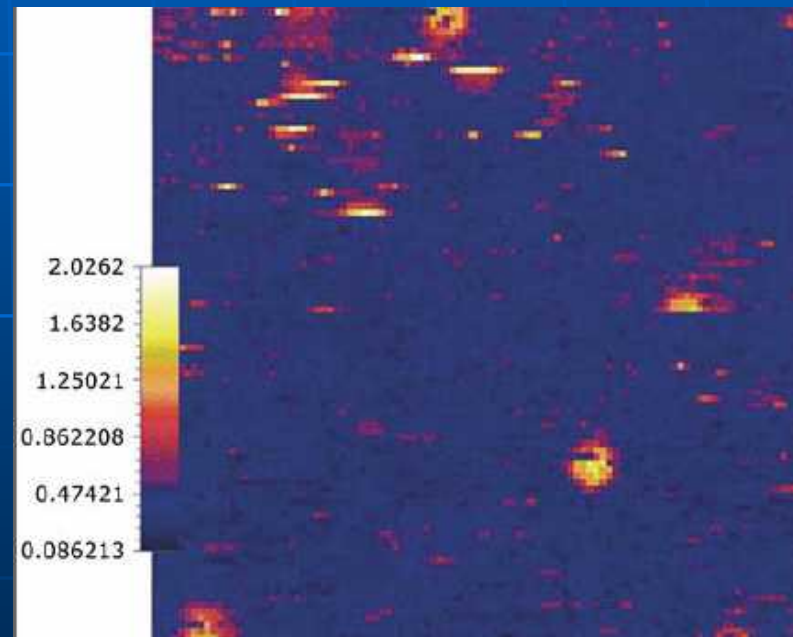
The microscope is able to filter out the out-of-focus light from above and below the point of focus in the object. A "pinhole" situated in front of the image plane which acts as a spatial filter and allows only the in-focus portion of the light to be imaged



Inverted microscope



Green fluorescent protein



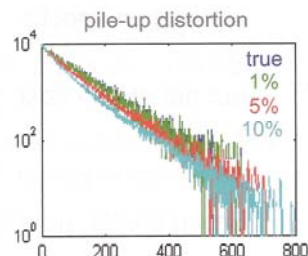
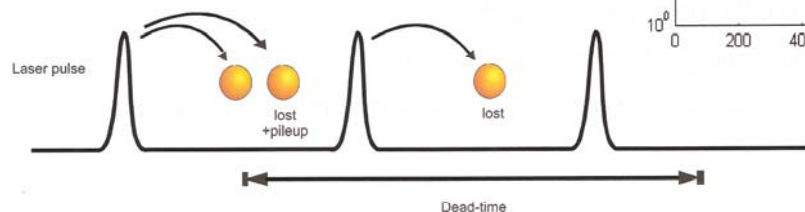
10 μm x 10 μm

TCSPC Disadvantages and Difficulties

- Need pulsed light source and fast electronics (no problem today)
- Need fast single photon detector (can be difficult in the IR)
- Considered complicated & expensive (not true today)

Specific Difficulties:

- Pile-up → keep count rate $< 1..2\%$, use high laser pulse rate
- Dead-time → keep count rate $< 1..2\%$, correct intensities

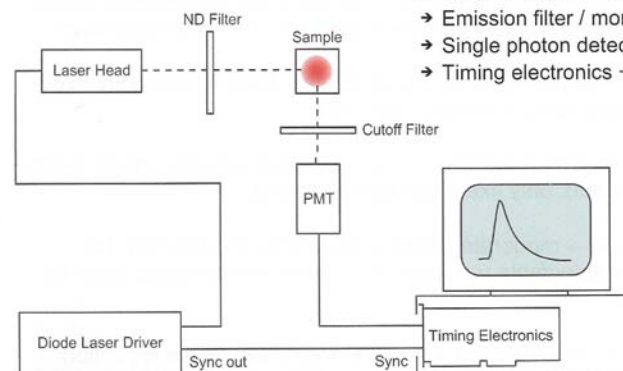


9

Basic TCSPC / TR-Fluorescence Setup

Typical system components:

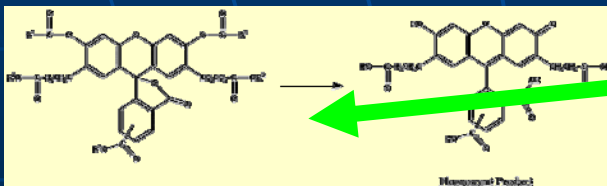
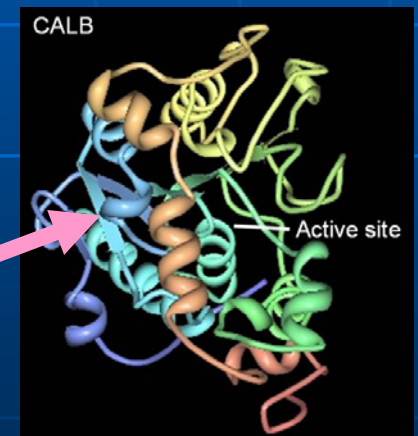
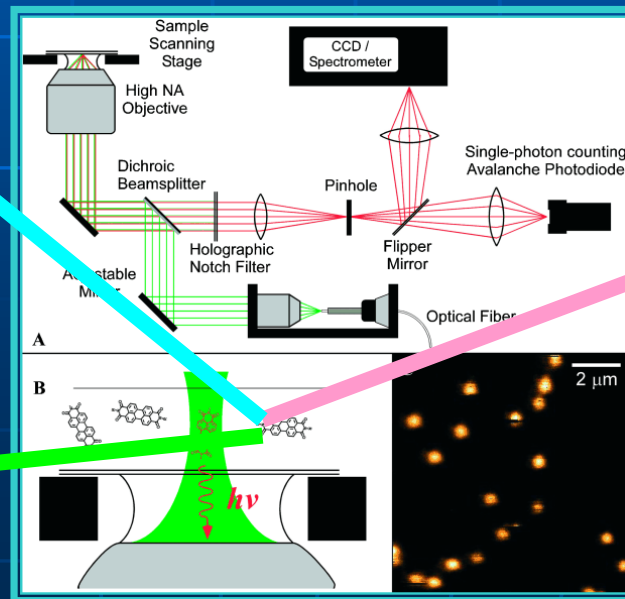
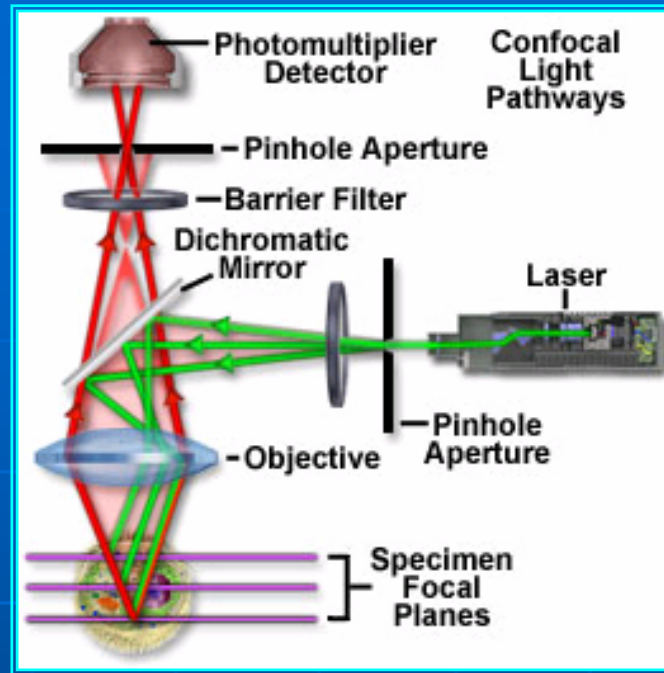
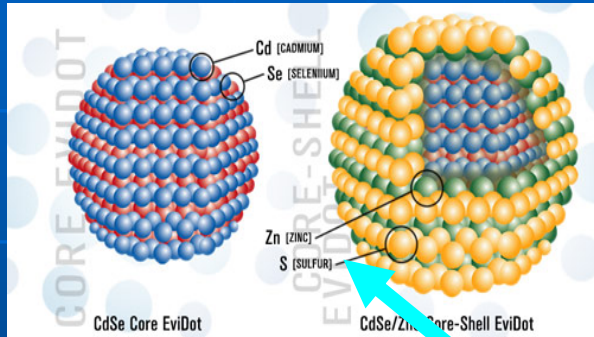
- Pulsed excitation light source
- Excitation filter / monochromator
- Emission filter / monochromator
- Single photon detector
- Timing electronics + computer



10

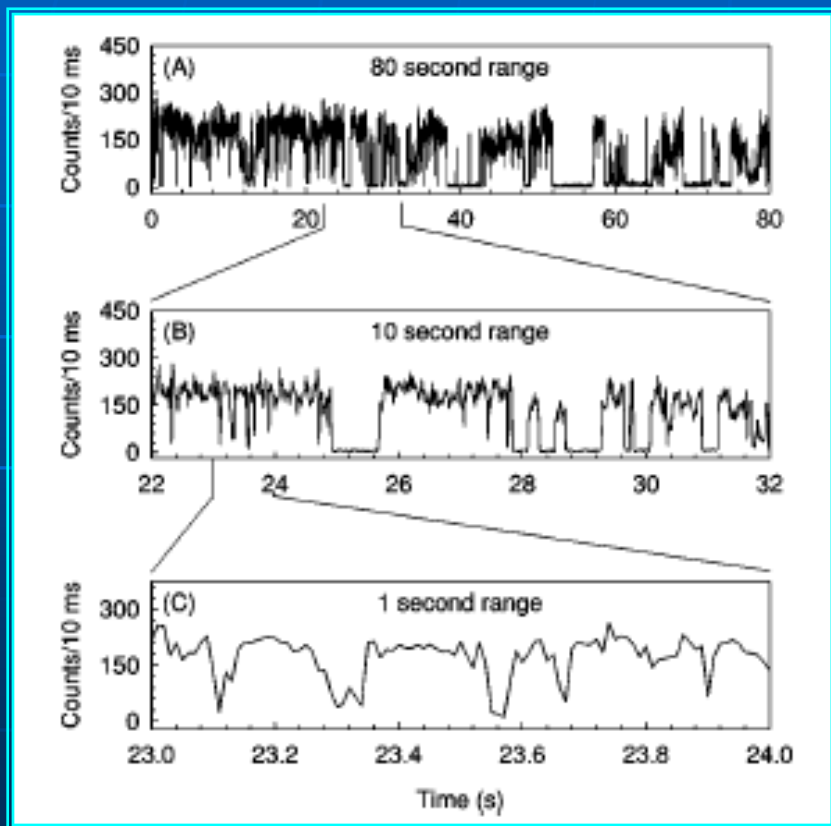


Confocal microscopy & Single-Particle Intermittency

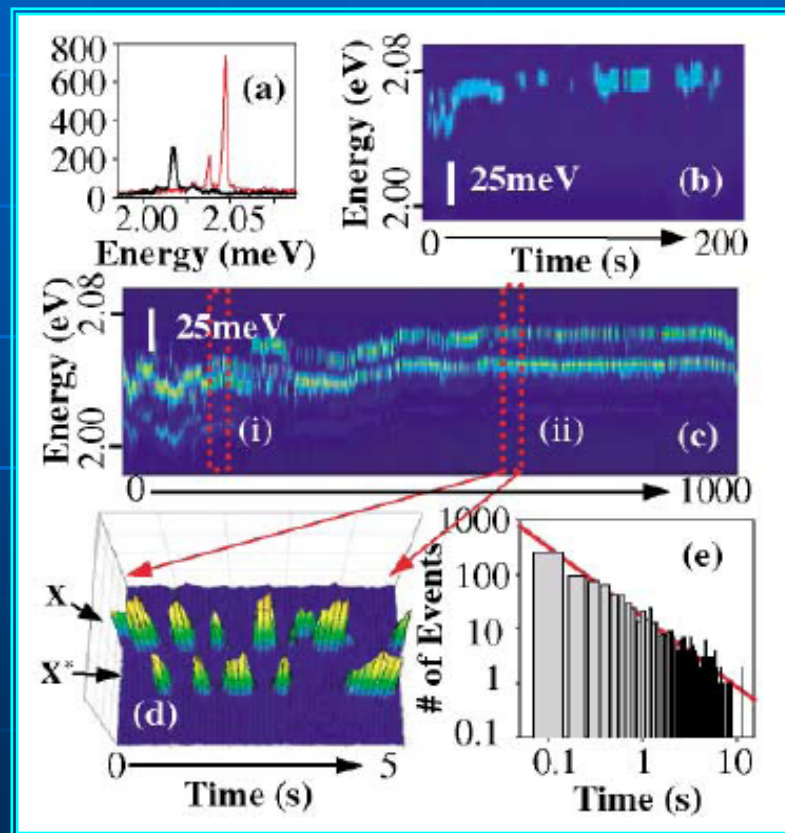


Fluorescence intermittency (blinking)

histogram

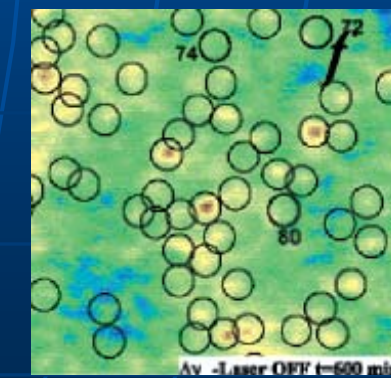
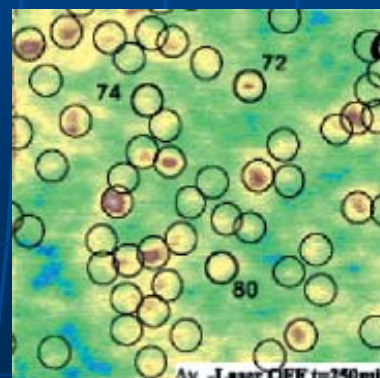
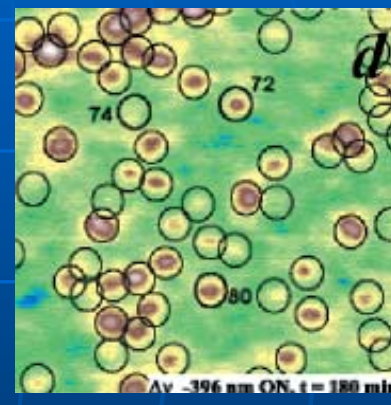
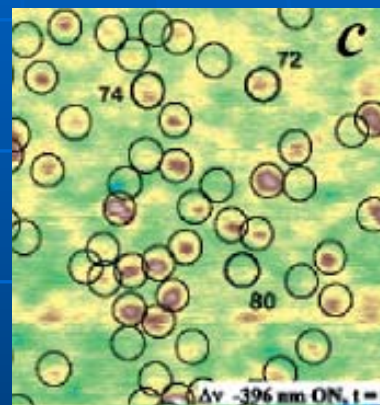
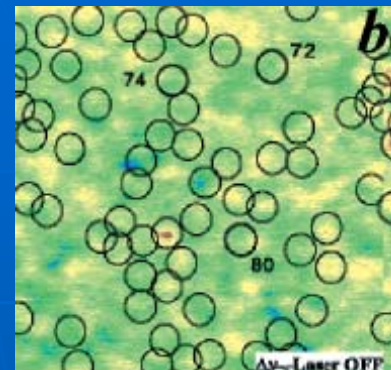
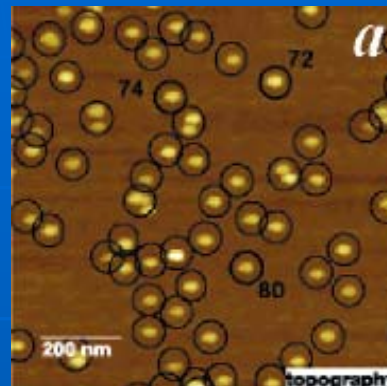
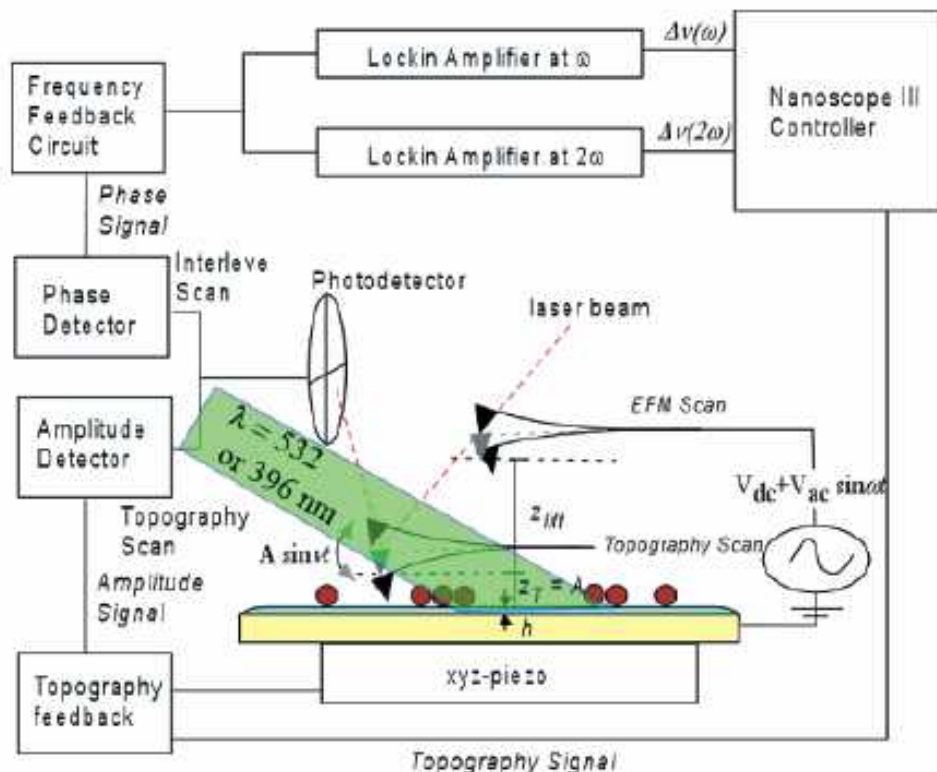


D. Nesbitt et al; *J. Chem. Phys.* 2000



M. G. Bawendi., et al. *PRL* 2002

Electrostatic force microscopy of CdSe



a) Topography of CdSe QDs

Charge map $10^3 \times 10^3 \text{ nm}^2$

b) Before illumination

C-d) *light-on*

e) 4 hrs f) 10 hrs after light -off

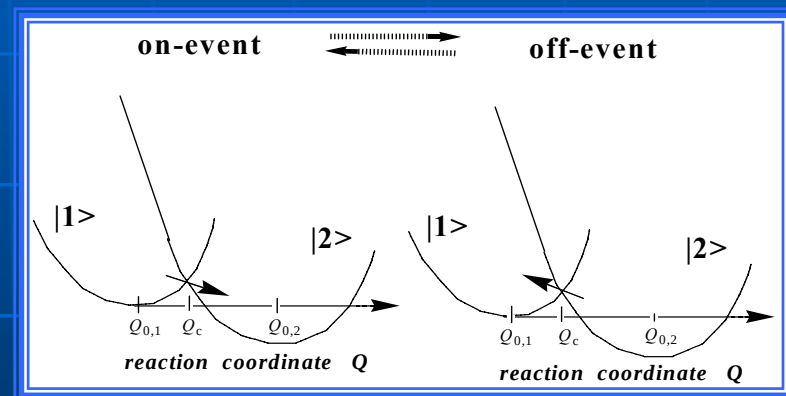
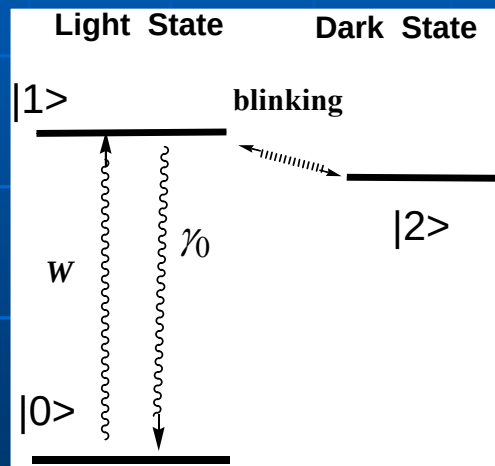
L. Brus, *et al.* *J. Phys. Chem. B*, **2004**, 108, 4946

- Why does a QD blink ?
- What causes a power law distribution (t^{-m}) ?
- Why is the exponent m close to $3/2$?

Diffusion-controlled electron transfer (DCET) model

(Tang & Marcus, PRL and JCP , 2005)

Neutral light state $\longleftrightarrow$ Positively charged dark state



1-D diffusion in energy space

$$\frac{\partial}{\partial t} \rho_1(Q, t) = L_1 \rho_1(Q, t) - C_1 \rho_1(Q, t)$$

$$\frac{\partial}{\partial t} \rho_2(Q, t) = L_2 \rho_2(Q, t) - C_2 \rho_2(Q, t)$$

$$L_k = D_k(t) \frac{\partial}{\partial Q} \left(\frac{\partial}{\partial Q} + \frac{1}{k_B T} \frac{\partial}{\partial Q} U_k(Q) \right)$$

$$C_k = \frac{2\pi |V_k|^2}{\hbar} \delta(U_{12}(Q))$$

DCET model for power-law intermittency

Blinking statistics $P(t)$: waiting time distribution for “on” or “off” events

$$\overline{P}_k(s) = -\int_0^\infty dt e^{-st} \frac{d}{dt} \left(\int_{-\infty}^\infty dQ \rho_k(Q, t) \right) = \frac{A_k \overline{G}_k(Q_c, Q_c; s)}{1 + A_k \overline{G}_k(Q_c, Q_c; s)}$$

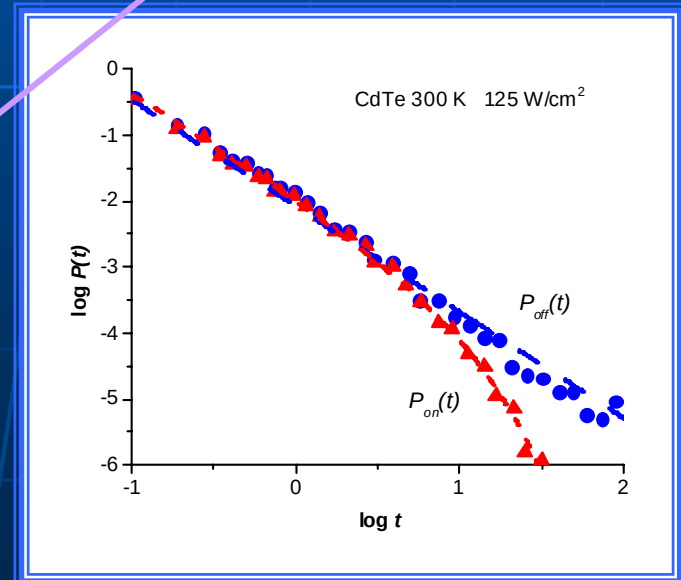
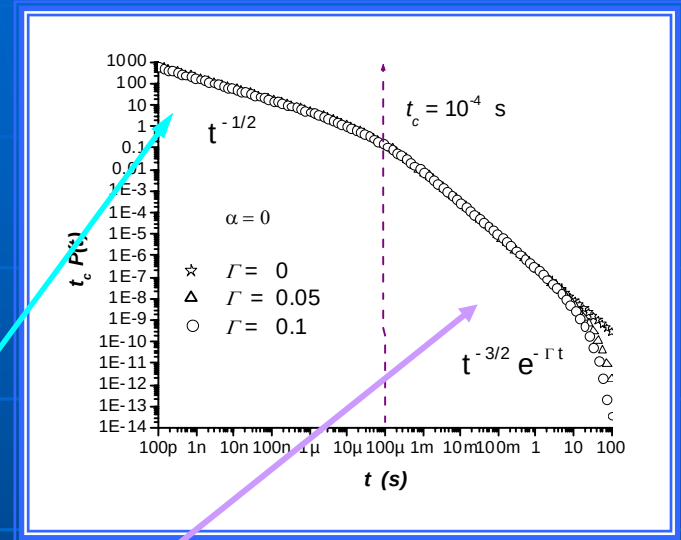
$$A_k = \frac{2\pi}{\hbar} |V_k|^2 \left| \partial(U_1(Q) - U_2(Q)) / \partial Q \right|_{Q=Q_c}$$

Debye medium (D is a constant)

$$P_k(t) \sim \frac{1}{\sqrt{\pi t_{c,k}}} t^{-1/2} \quad \text{if } t \ll t_{c,k}$$

$$P_k(t) \sim \frac{\sqrt{t_{c,k}}}{2\sqrt{\pi}} t^{-3/2} \exp(-\Gamma_k t) \quad \text{if } t_{c,k} \ll t \ll \tau_{L,k}$$

CdTe

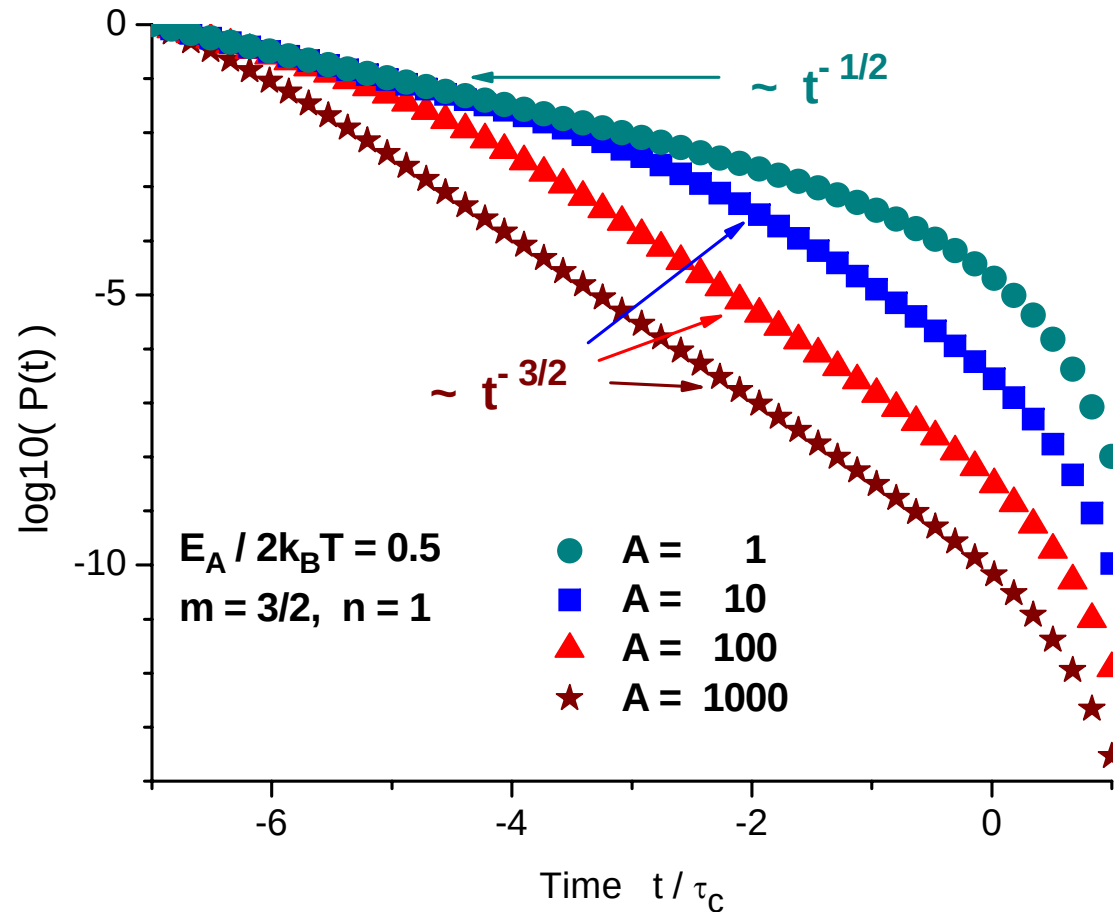


$$P_k(t) = \frac{1}{\sqrt{\pi t_{c,k}}} \left[1 - \sqrt{\pi t / t_{c,k}} \exp(t / t_{c,k}) \operatorname{erfc}(\sqrt{t / t_{c,k}}) \right] \exp(-\Gamma_k t),$$

(DCET model)

Small A:
weak coupling

Large A:
strong coupling

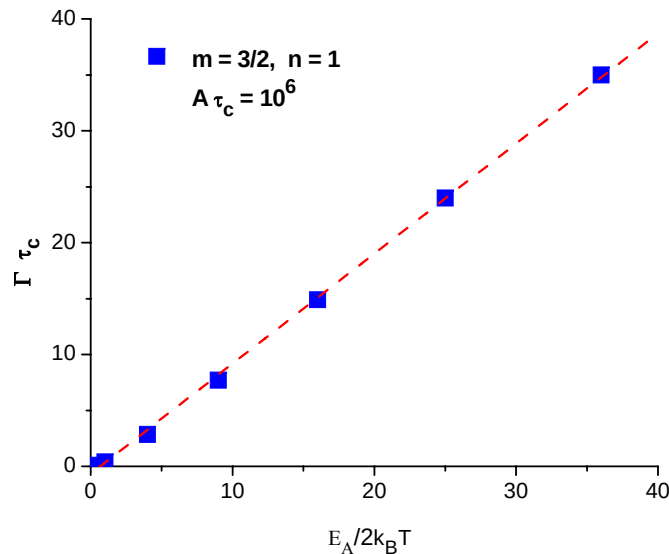


Bending rate Γ at longer times

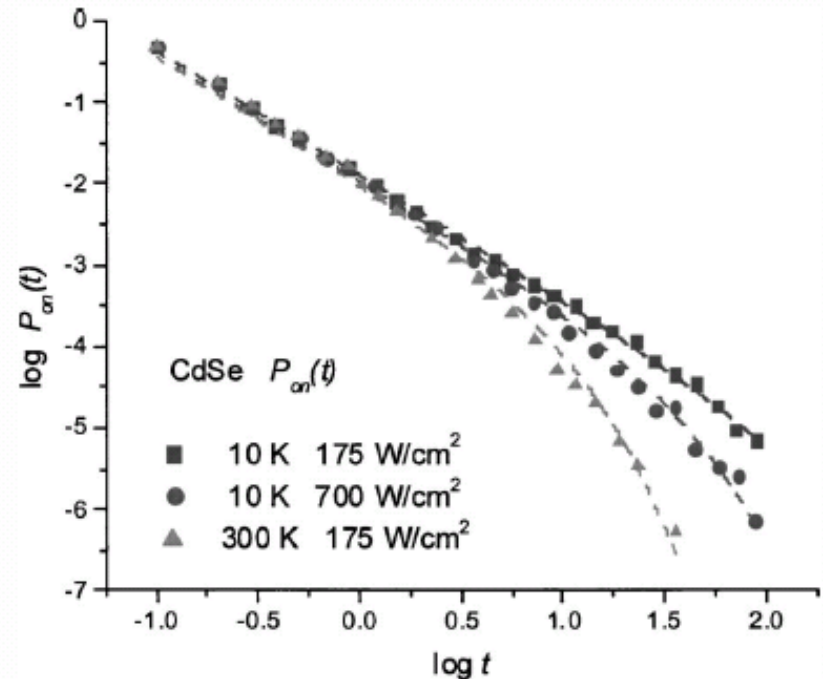
DCET model

$$\Gamma_k \tau_k \equiv \frac{E_{A,k}}{2k_B T},$$

$$E_{A,1} \equiv \frac{(\lambda + \Delta G^0)^2}{4\lambda}, \quad E_{A,2} \equiv \frac{(\lambda - \Delta G^0)^2}{4\lambda}$$



Bending rate and diffusion rate increase with:
 Light intensity and QD size
 Temperature (faster diffusion)



Normal vs. anomalous diffusion

$$\sigma^2(t) \sim t^{\beta_{CD}}$$

$\beta_{CD} < 1$ (subdiffusion)

$\beta_{CD} = 1$ (normal diffusion)

$\beta_{CD} > 1$ (superdiffusion)

Debye dielectric medium

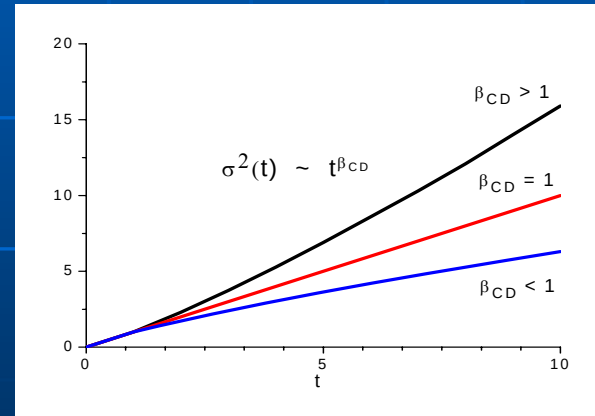
$$\bar{\chi}(s) = 1/(1 + s \tau_D) \quad \bar{\tau}_{L,k}(s) = \tau_{L,k} \equiv \tau_{D,k} \varepsilon_\infty / \varepsilon_0,$$

$$\Theta_k(t) = \exp(-t / \tau_{L,k}) \quad D_k(t) = \Delta_k^2 / \tau_{L,k}$$

$$\sigma^2(t) = \langle (Q(t) - \langle Q(t) \rangle)^2 \rangle \approx 2(\Delta^2 \tau_D / \tau_L) t$$

$$D_k(t) = -\Delta_k^2 \frac{d}{dt} \ln \Theta_k(t), \quad \Theta_k(t) \equiv L^{-1} \left(\frac{1}{s + \frac{1}{\tau_{L,k}(s)}} \right)$$

$$\bar{\tau}_{L,k}(s) = \frac{\varepsilon_\infty}{\varepsilon_0} \frac{1 - \bar{\chi}_k(s)}{s \bar{\chi}_k(s)}, \quad \bar{\chi}_k(s) = \frac{\bar{\varepsilon}_k(s) - \varepsilon_\infty}{\varepsilon_0 - \varepsilon_\infty}.$$



Cole-Davison medium

$$\bar{\chi}(s) = 1/(1 + s \tau_D)^{\beta_{CD}}$$

$D(t)$ is time-dependent

More general anomalous diffusion

J. Tang, PR B (2007)

$$\frac{\partial}{\partial t} G(Q, Q'; t) = \frac{\partial}{\partial Q} \left(D_2(t) \frac{\partial}{\partial Q} + D_1(t) Q \right) G(Q, Q'; t).$$

$$D_1(t) = - \frac{d}{dt} \ln \langle Q(t) \rangle$$

$$G(Q, Q'; t) = \frac{1}{\sqrt{2\pi \Delta_2(t)}} \exp \left[- \frac{(Q - Q' \Delta_1(t))^2}{2\Delta_2(t)} \right]$$

$$D_2(t) = D_1(t) \Delta_2(t) + \frac{1}{2} \frac{d}{dt} \Delta_2(t).$$

$$\Delta_1(t) \approx 1 - (t / \tau_c)^\nu$$

Kohlrausch-Williams-Watts mode
(Havriliak-Negami dielectric function)

$$\Delta_2(t) \approx \Delta_2(\infty) (t / \tau_c)^\mu$$

$$\varepsilon(\omega) = \varepsilon_\infty + \Delta\varepsilon / \left[1 + (i \omega \tau_{HN})^\alpha \right]^\gamma$$

Ornstein-Uhlenbeck process: $\mu = \nu = 1$; D_1, D_2 : constant

Non-Debye medium (anomalous diffusion)

J. Tang and R. A. Marcus, PRL and JCP (2005)

$$\overline{P}_k(s) \approx \frac{1}{1 + \left((s + \Gamma_k) t_{c,k} \right)^{1 - \beta_{CD}/2}}$$

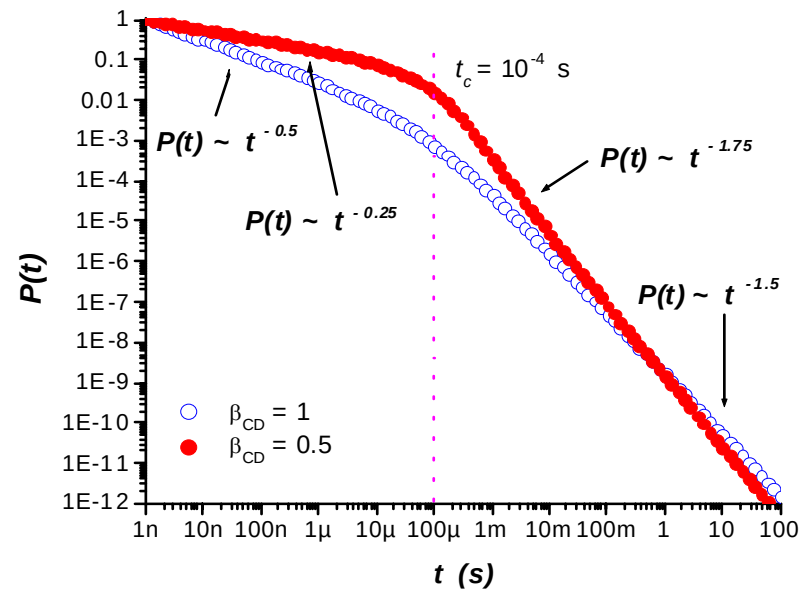
Mittag-Leffler function

$$E_a(z) = \sum_{n=0}^{\infty} z^n / \Gamma(na + 1).$$

$$P_k(t) \sim \frac{d}{dt} \left[E_{-1+\beta_{CD}/2} \left(-(t/t_{c,k})^{-1+\beta_{CD}/2} \right) \right] \exp(-\Gamma_k t) \quad \text{if } t \leq 1/\Gamma_k < \tau_{L,k}$$

$$P_k(t) \sim (t/t_{c,k})^{-\beta_{CD}/2} / \Gamma(1 - \beta_{CD}/2) t_{c,k}, \quad \text{if } t < t_{c,k}$$

$$P_k(t) \sim (t/t_{c,k})^{-2+\beta_{CD}/2} \exp(-\Gamma_k t) / \left| \Gamma(\beta_{CD}/2 - 1) \right| t_{c,k} \quad \text{if } t_{c,k} < t \leq 1/\Gamma_k < \tau_{L,k}.$$



Blinking statistics of nanorods and size dependence

J. Tang, JPC (2007)

$$P(t) \approx P_0 (t/\tau_c)^{-2+\mu/2} \exp\left(-(\Gamma t)^{2\nu-\mu}\right)$$

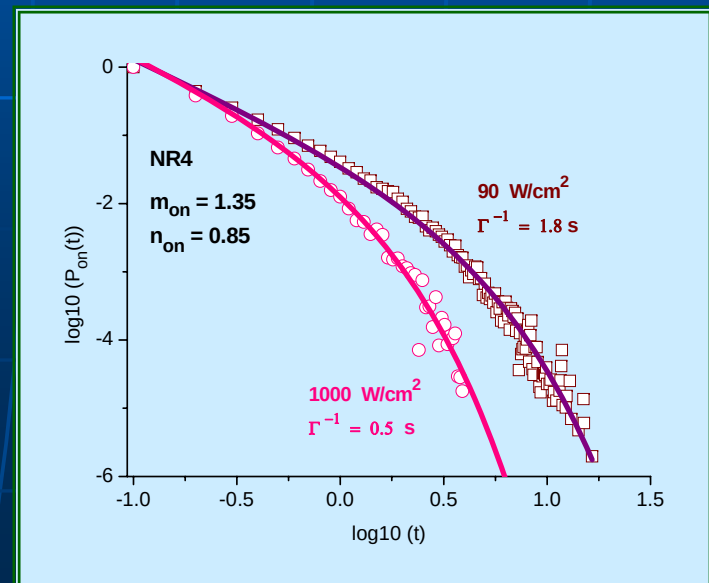
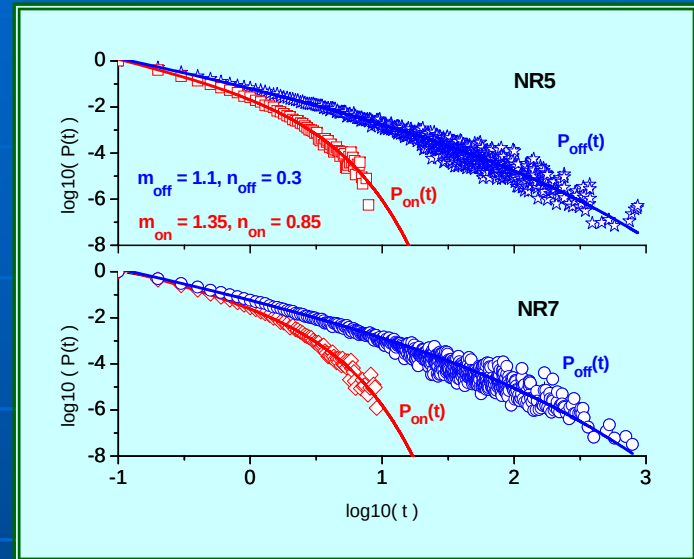
$$\propto t^{-m} \exp\left(-(\Gamma t)^n\right),$$

$$m = 2 - \mu/2$$

Size dependence

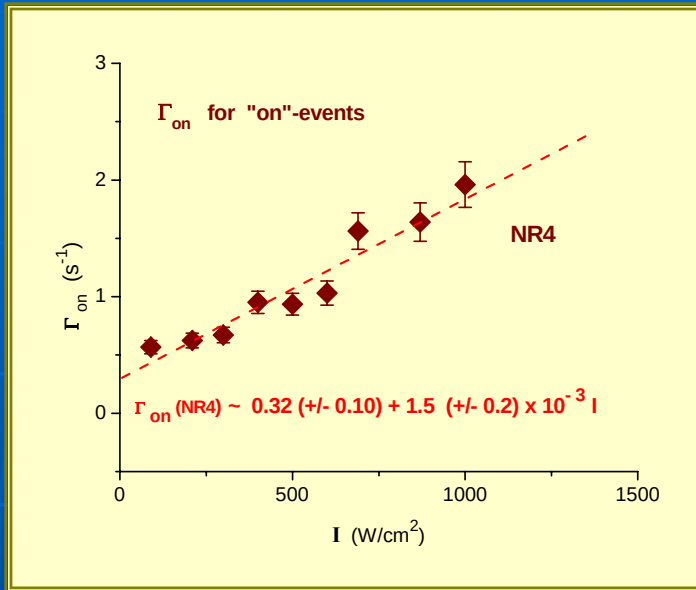
$$n = 2\nu - \mu$$

Intensity dependence

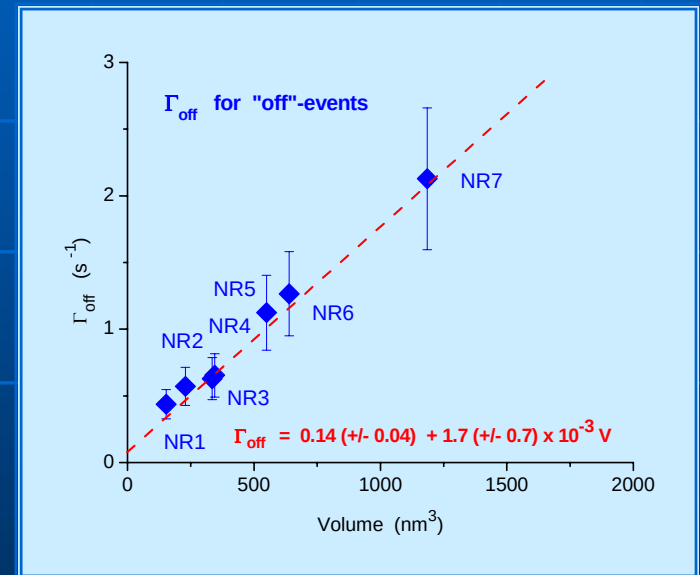


Light NRs

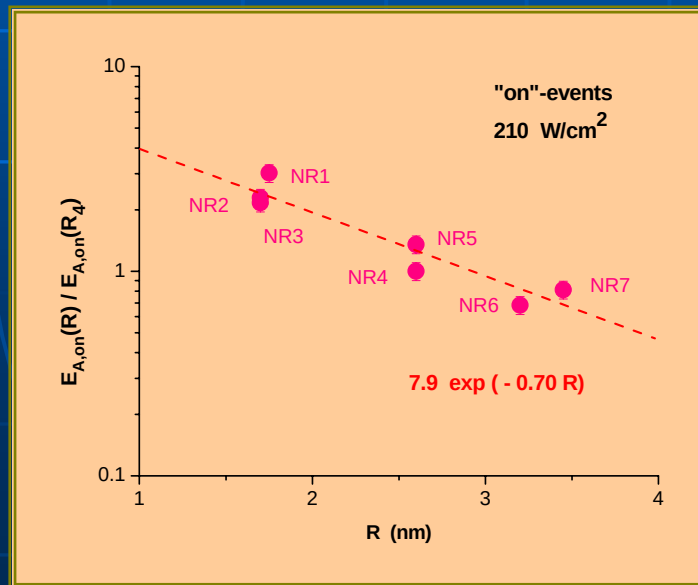
Intensity dependence



Dark NRs



Volume dependence



Radius dependence

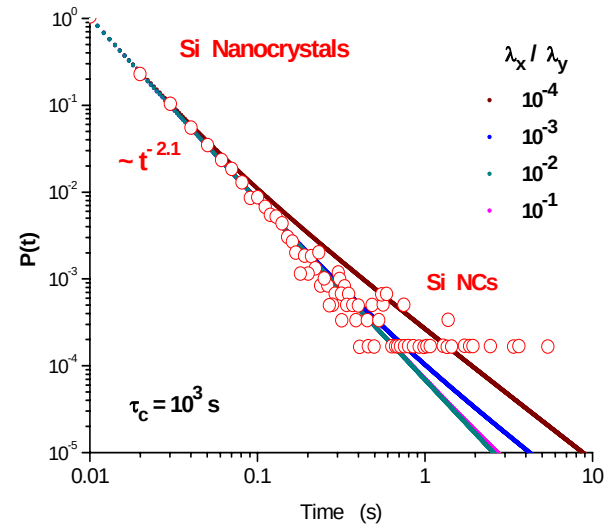
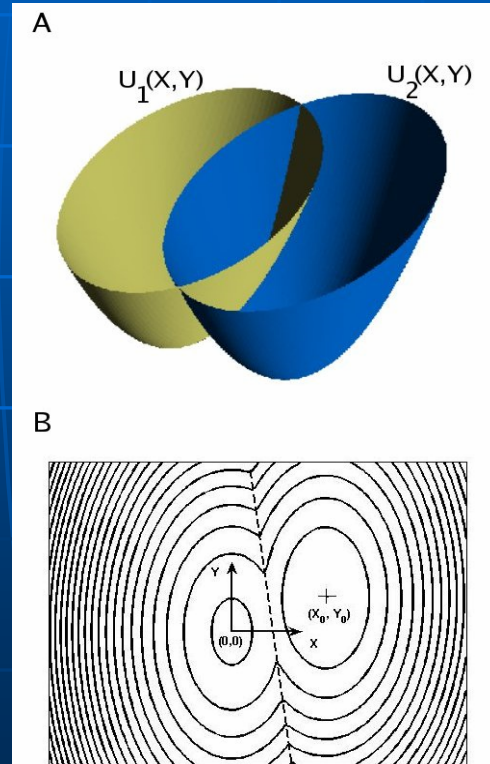
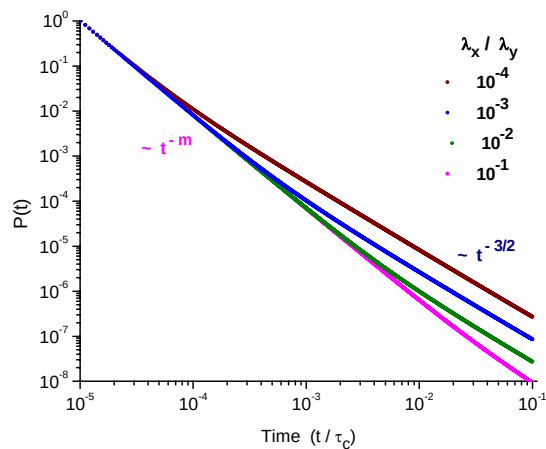
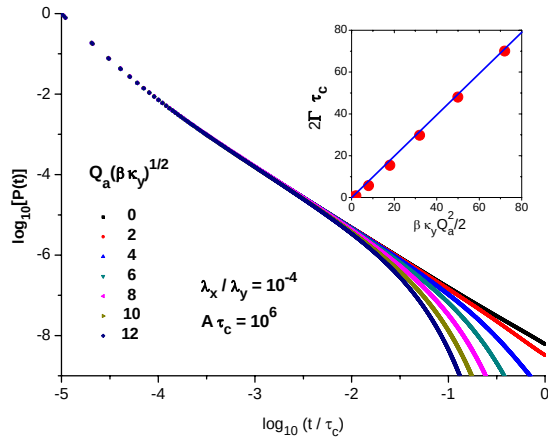
Exotic power law for Si blinking

J. Tang, JCP (2007)

2-D model
Fast and slow
reaction coordinates

Crossing point is not like
A δ -function

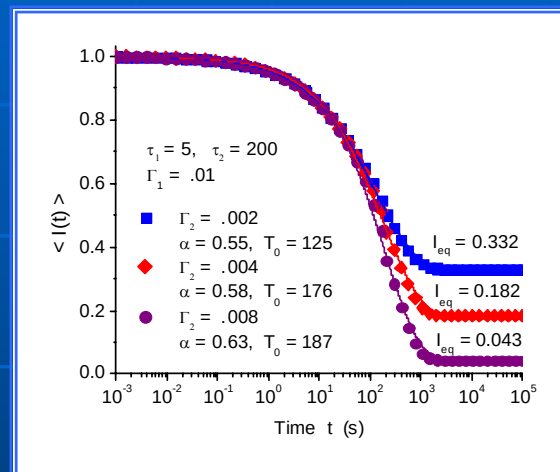
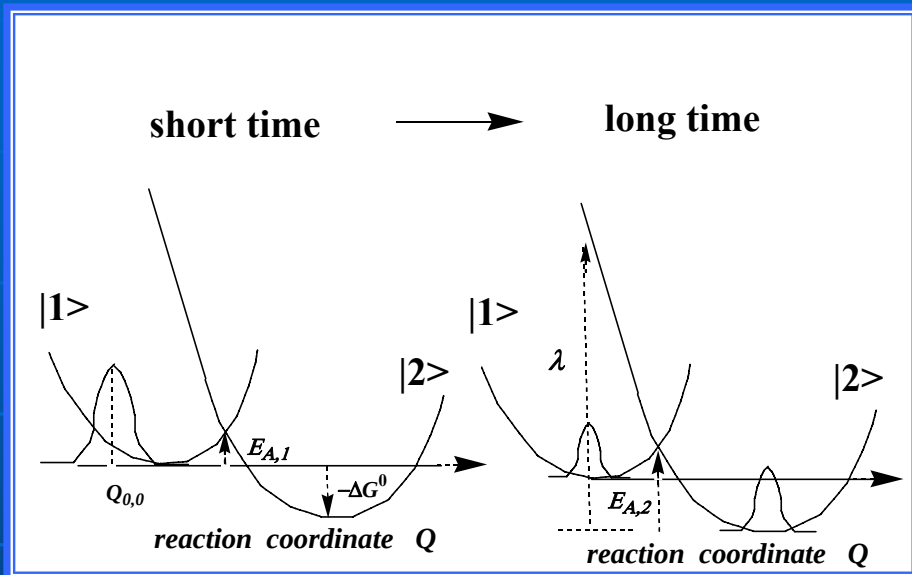
Power-law exponent
Could be greater than 2



Ensemble fluorescence intensity decay

stretched exponential decay

$$I_{eq} + (1 - I_{eq}) \exp\left(-(t/T_0)^\alpha\right)$$



Single particle

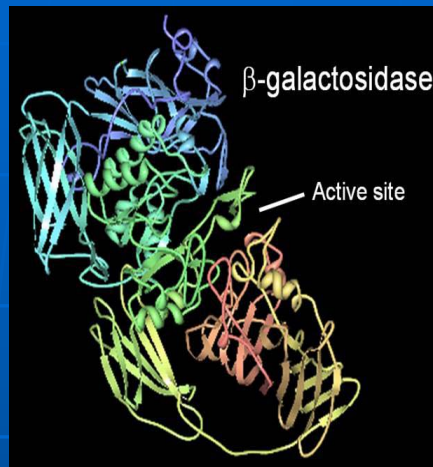
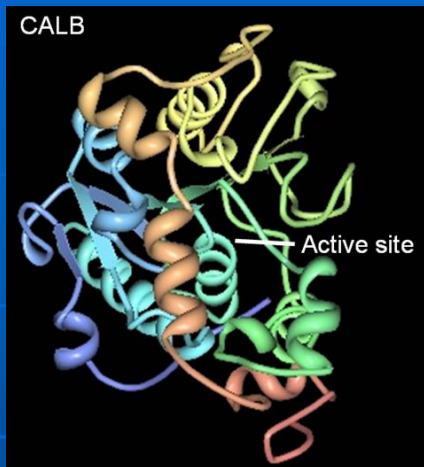
vs.

ensemble

- Power-law decay
- distinguishable
- at the crossing initially
- nonergodic

- Stretched exponential decay
- indistinguishable
- in the ground state initially
- ergodic

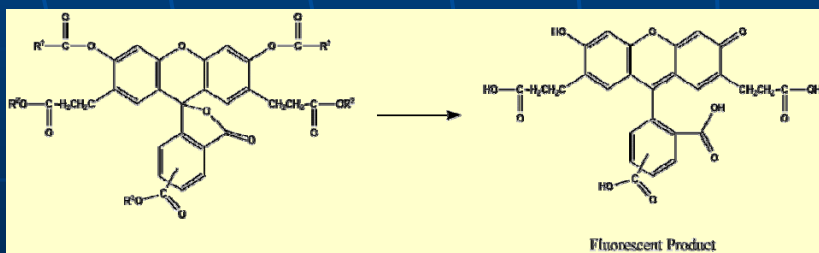
Intermittency of single-enzyme activity



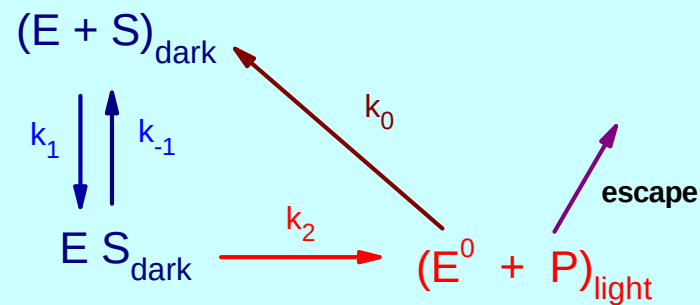
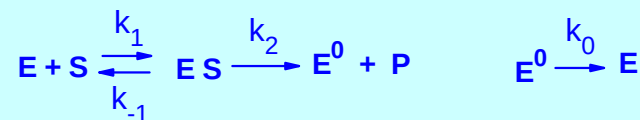
X. S. Xie, et al. Nature Chem. Bio. 2, 87 (2005)

O. Flomenbom, et al., PNAS 102, 2368 (2005)

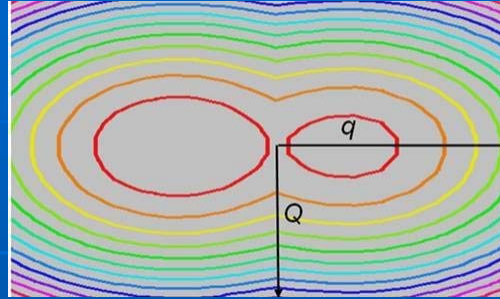
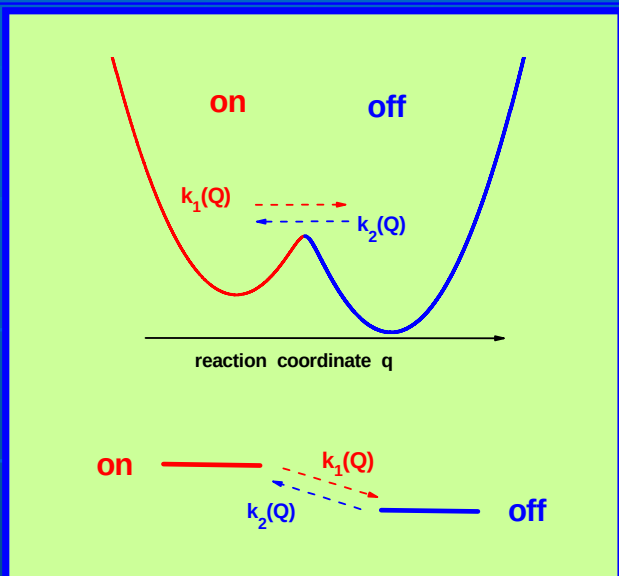
J. Tang, JPC (submitted)



Michaelis-Menten Scheme

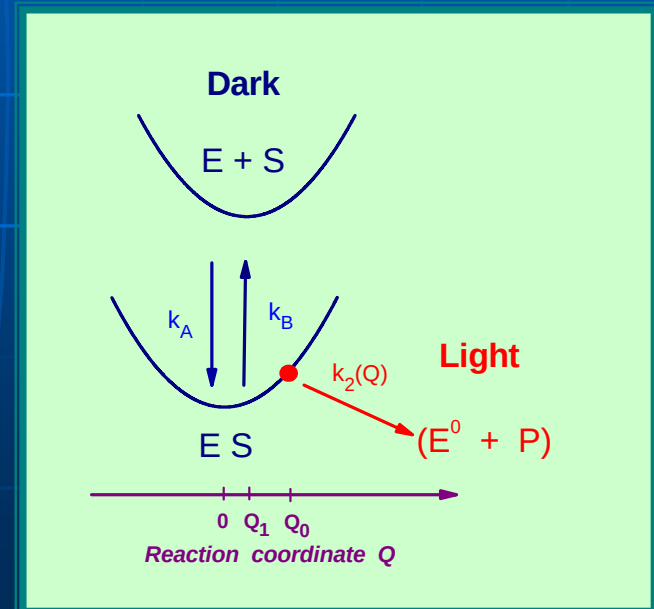
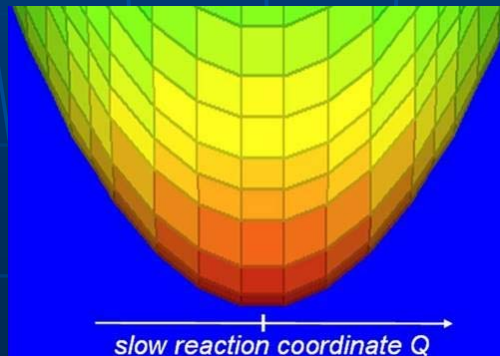
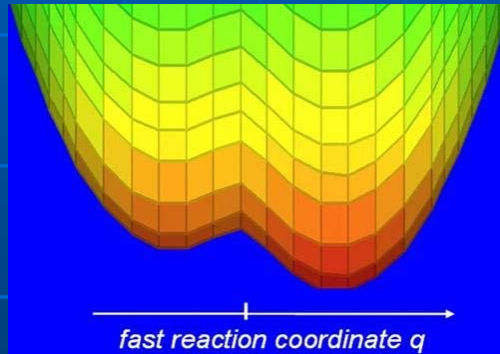


Michaelis-Menten scheme for enzyme reactions

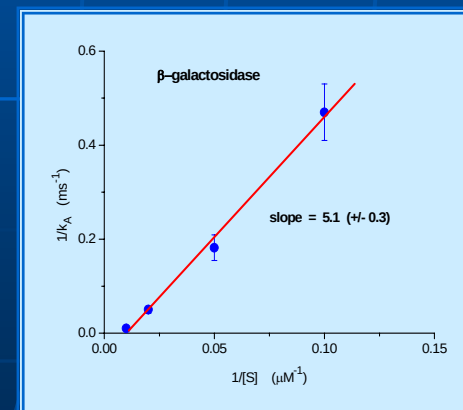
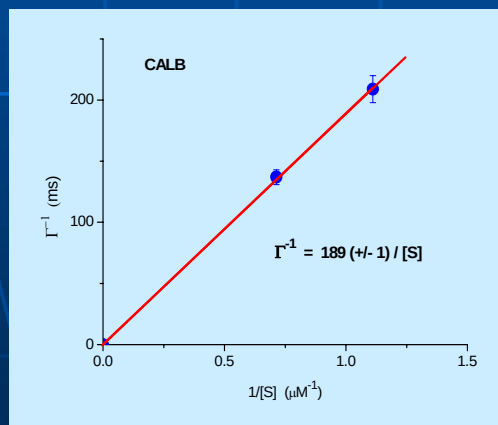
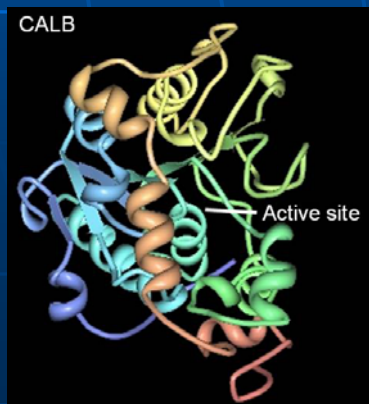
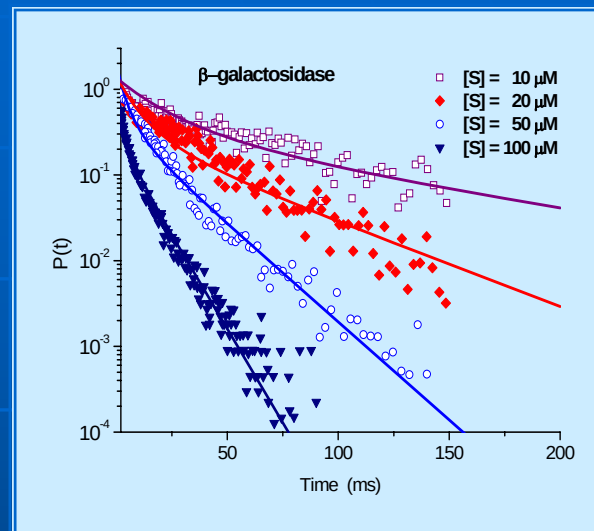
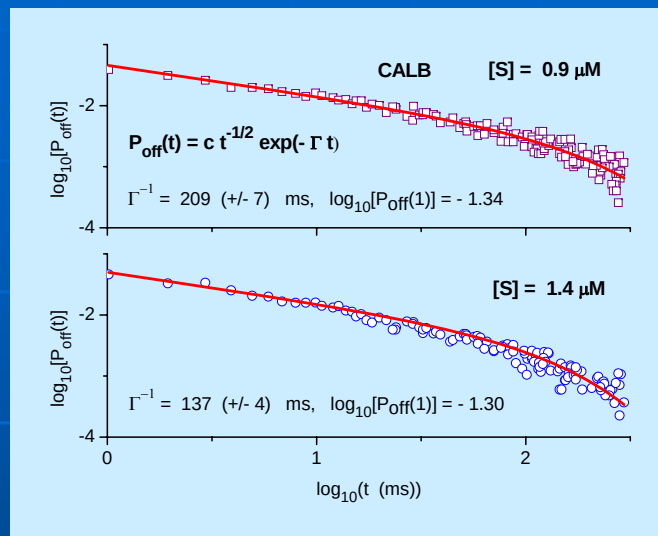


$$P(t) = \frac{k_A k_2}{\lambda_2 - \lambda_1} (e^{-\lambda_1 t} - e^{-\lambda_2 t})$$

$$\lambda_{1,2} = k_A + k_B + k_2 \mp \sqrt{(k_A + k_B + k_2)^2 / 4 - k_A k_B}.$$



Diffusion-controlled Michael-Menten reaction

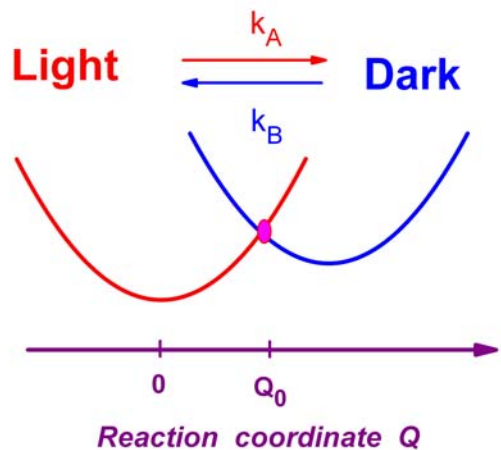


Fluctuations of potential barrier height

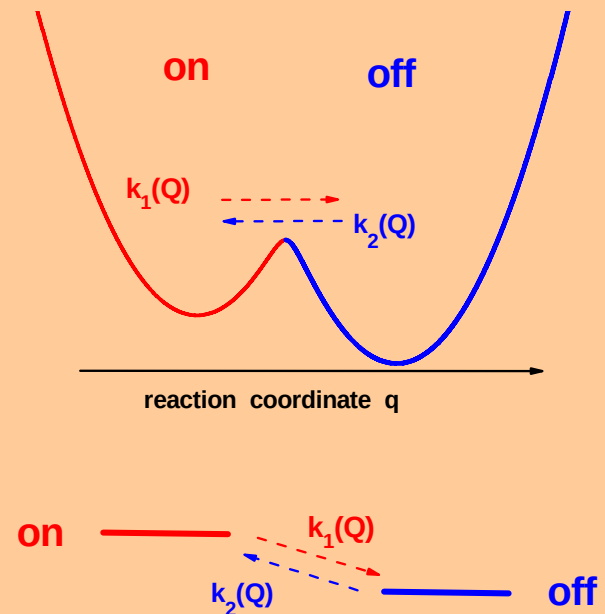


Non-exponential blinking statistics

Charge transfer reactions



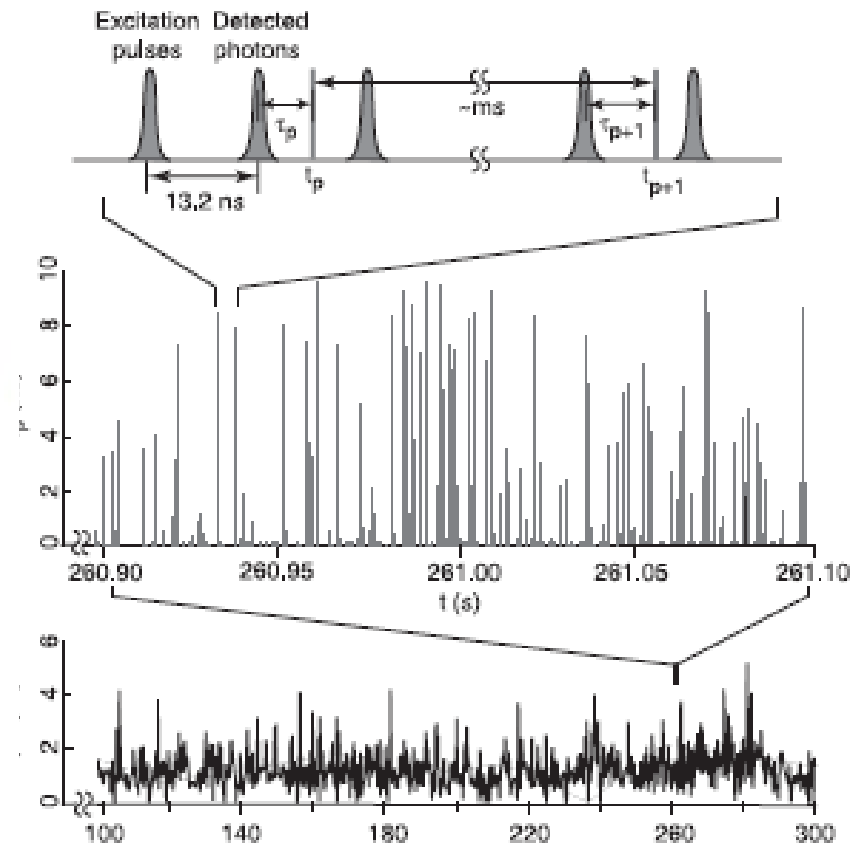
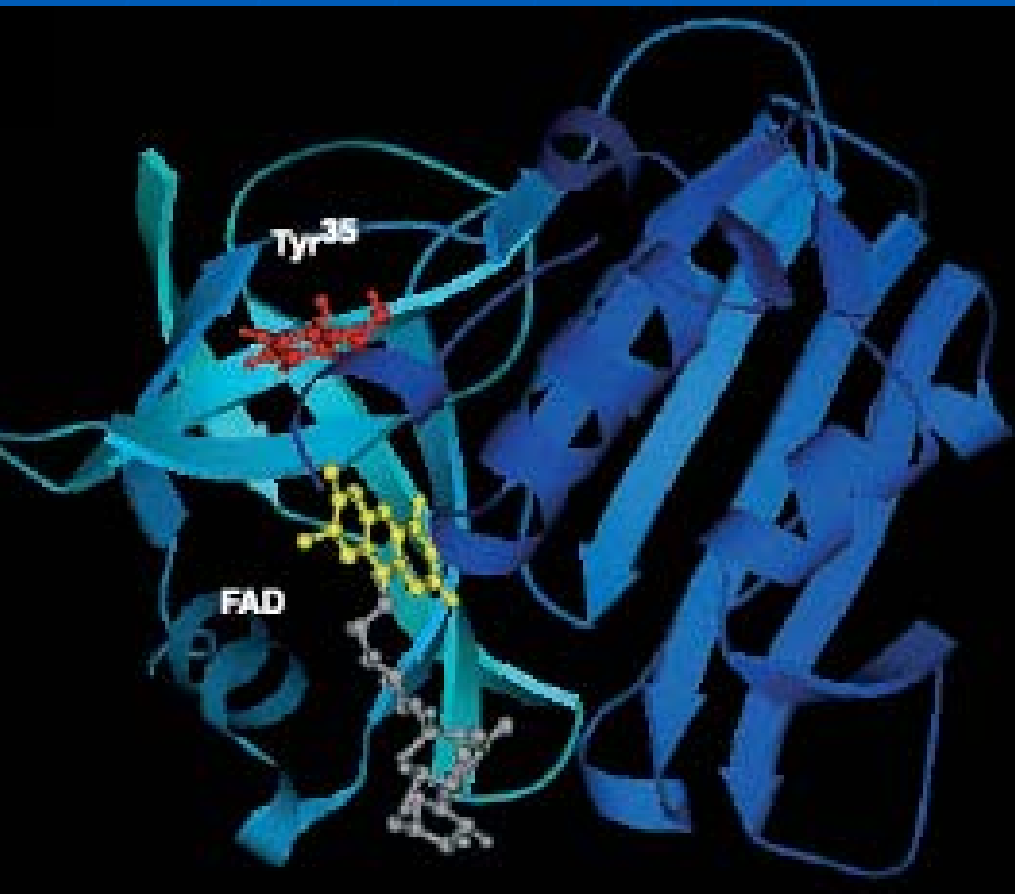
Catalytic reactions



Fluctuating electron transfer rates in single protein molecules

Fre / FAD complex

H. Yang *et al.* *Science* 302, 262 (2003)



Distance-fluctuation model

$$\gamma = \frac{2\pi |V_{ex}|^2}{\hbar \sqrt{4\pi \lambda k_B T}} \exp\left(-(\lambda + \Delta G^0)^2 / 4\lambda k_B T\right)$$

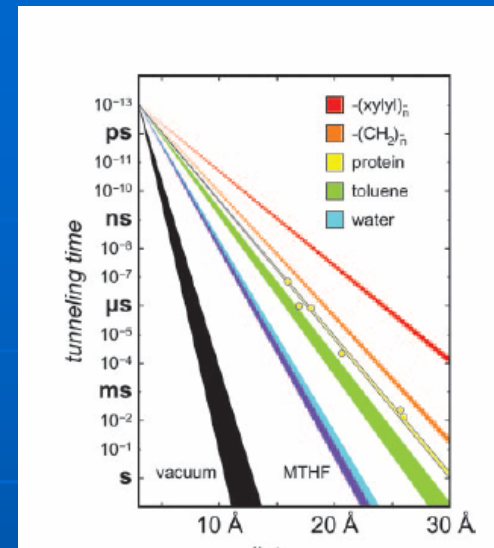
$$|V_{ex}|^2 \sim \exp(-\beta_{DA} \ell(t))$$

$$C_2(t) \equiv \frac{\langle \delta \gamma^{-1}(t) \delta \gamma^{-1}(0) \rangle}{\langle \gamma^{-1}(t) \rangle \langle \gamma^{-1}(0) \rangle}$$

$$C_2(t) \equiv \exp(\beta^2 C_Q(t)) - 1$$

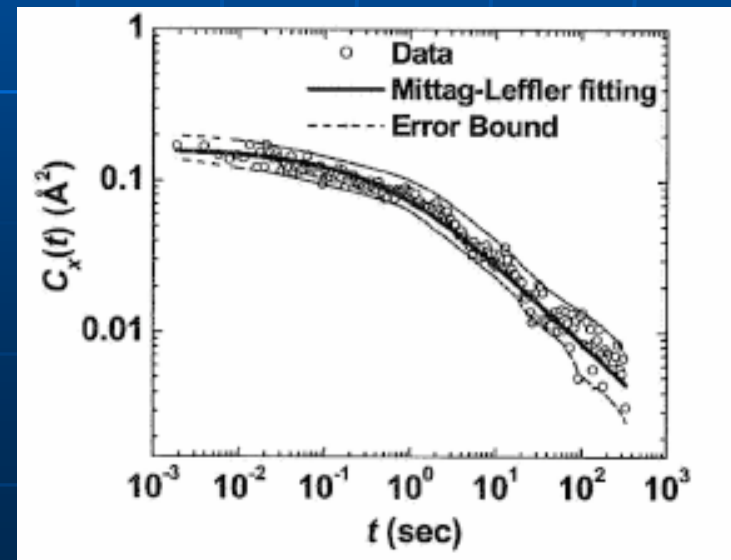
$$C_Q(t) \equiv \langle \delta \ell(t) \delta \ell(0) \rangle,$$

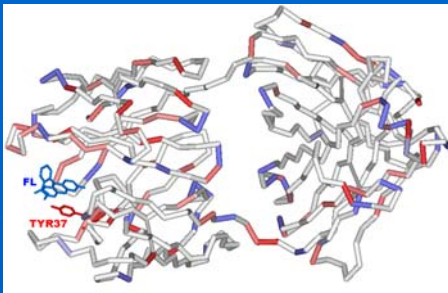
$$\delta \ell(t) = \ell(t) - \langle \ell(t) \rangle$$



If ET is activationless $\Delta G^0 = -\lambda$

ET rate $\sim \exp(-\beta_{DA} \ell)$



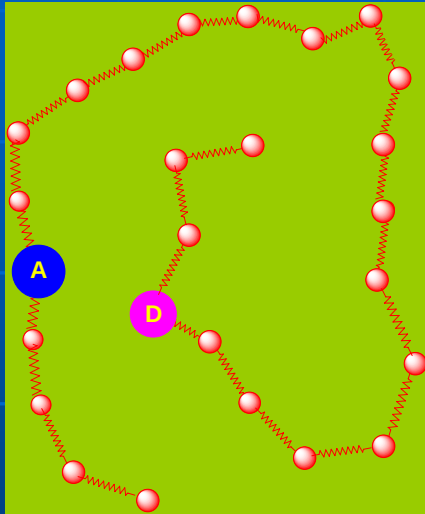


Distance-fluctuation model of a Rouse chain

J. Tang and R. A. Marcus, PR E (2006)

Polymer, DNA, protein

- Neutron scattering, NMR



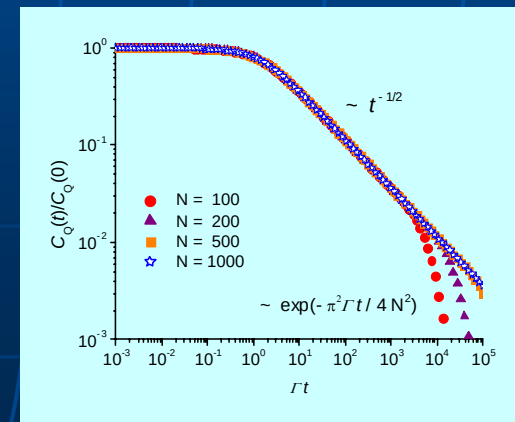
$$\frac{d}{dt} \begin{pmatrix} Q_0 \\ Q_1 \\ \vdots \\ Q_{N-1} \end{pmatrix} + \frac{\omega^2}{\zeta} \begin{pmatrix} 1 & -1 & 0 & 0 \\ -1 & 2 & -1 & 0 \\ 0 & -1 & 2 & -1 \\ 0 & 0 & -1 & 2 \end{pmatrix} \begin{pmatrix} Q_0 \\ Q_1 \\ \vdots \\ Q_{N-1} \end{pmatrix} = \begin{pmatrix} F_0(t)/m\gamma \\ F_1(t)/m\gamma \\ \vdots \\ F_{N-1}(t)/m\gamma \end{pmatrix}$$

$$\zeta \frac{d}{dt} \mathbf{Q}(t) + \omega^2 \mathbf{R} \mathbf{Q}(t) = \mathbf{F}(t)/m$$

$$\langle F_i^\mu(t) \rangle = 0$$

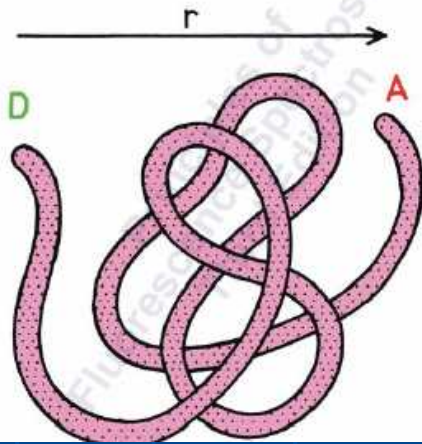
$$\langle F_i^\mu(t) q_j^\nu(t) \rangle = 0$$

$$\langle F_i^\mu(t) F_j^\nu(\tau) \rangle = 2mk_B T \zeta \delta(t - \tau) \delta_{ij} \delta^{\mu\nu}$$



Förster's Fluorescence resonance energy transfer (FRET)

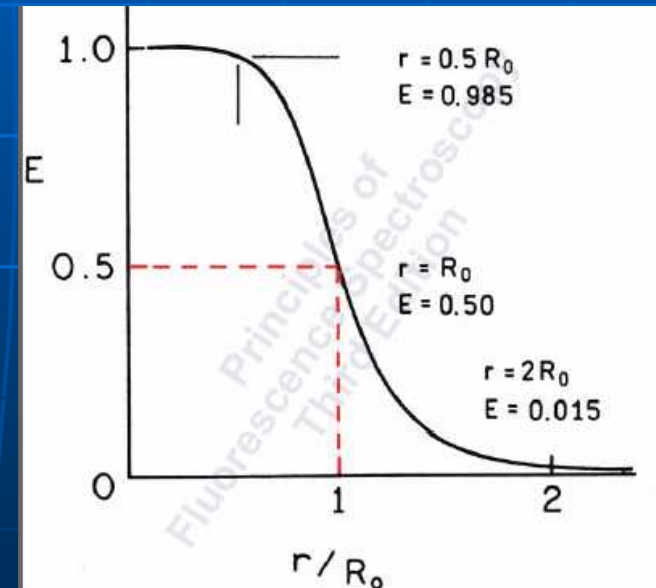
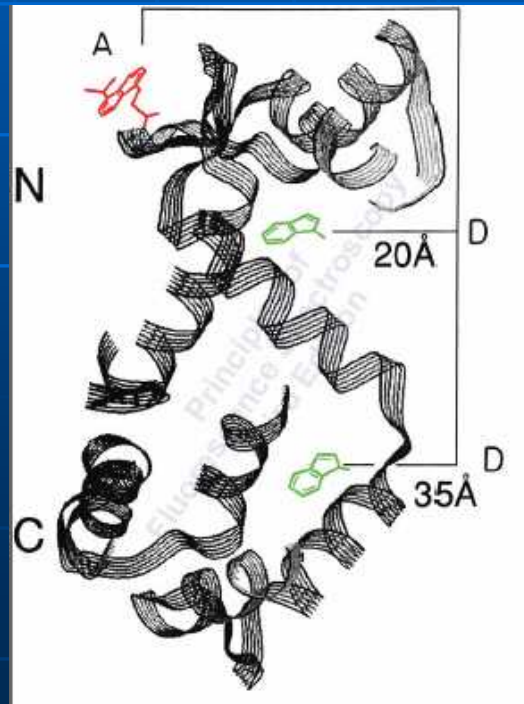
$$k_T(r) = \frac{1}{\tau_D} \left(\frac{R_0}{r} \right)^6 = \text{TRANSFER RATE}$$

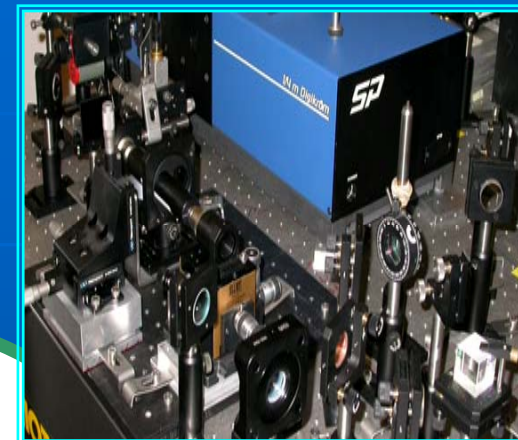
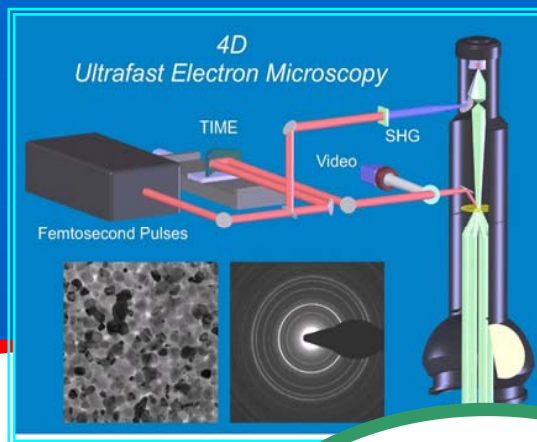
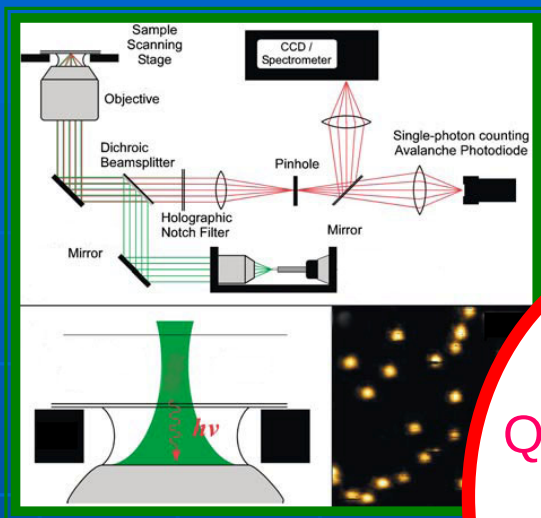


Efficient of energy transfer

$$E = \frac{k_T(r)}{\tau_D^{-1} + k_T(r)}$$

Spinach calmodulin
D: tryptophan
A: AEDANS





Quantum dots,
Nanorods,
Biomolecules

Photochemistry,
Photovoltaics

Light-induced
phase transitions

Protein dynamics &
Biomed applications



