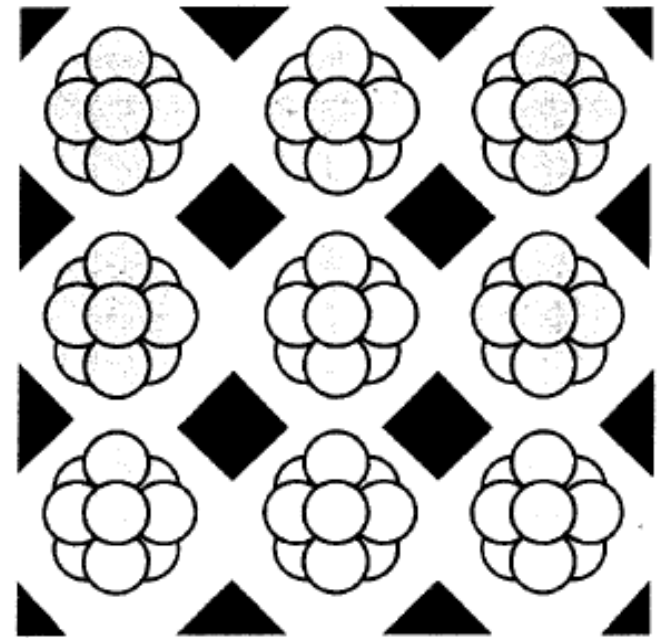


6.2.3 Arrays of nanoparticles in Zeolites

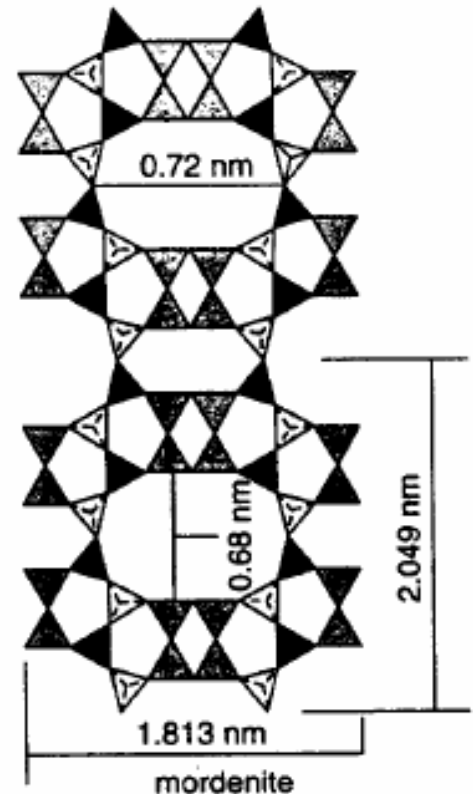
- Zeolites: Cubic mineral
- $(\text{Na}_2, \text{Ca})\text{Al}_2\text{Si}_4\text{O}_{12} \cdot 8\text{H}_2\text{O}$
- Porous materials



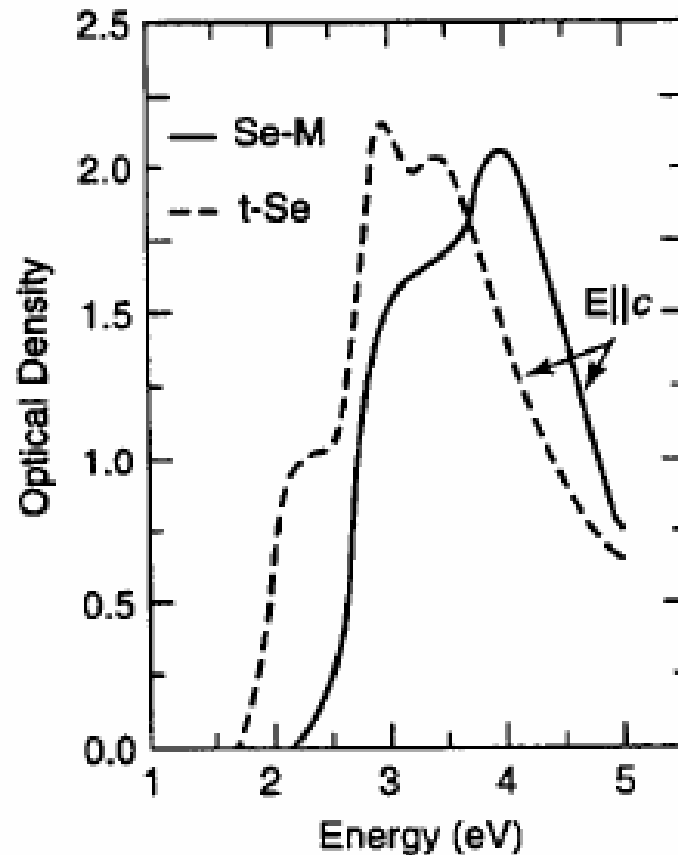
Schematic of cluster assemblies in zeolite pores. (With
ov et al., in *Handbook of Nanostructured Materials and*
1., Academic Press, San Diego, 2000, Vol. 4, Chapter 4, p. 23

Zeolite--Mordenite

- $(\text{Ca,Na,K}_2)(\text{Al}_2\text{Si}_{10})\text{O}_{24}\cdot 7\text{H}_2\text{O}$
- Orthorhombic
- Long parallel channels
- With $d = 0.6 \text{ nm}$
- Se atoms fill into the channel
- to form chains of single atom
-



Optical absorption



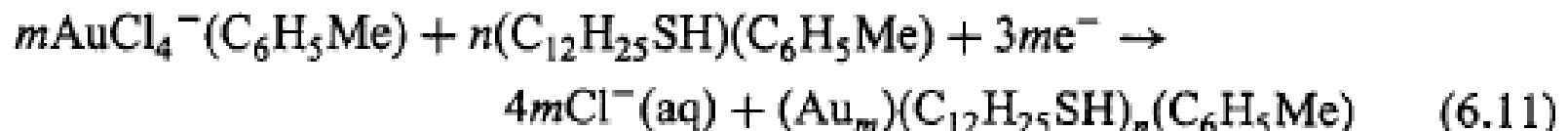
The electronic structure of Se chains is different from Se crystal

6.2.4 Crystals of metal nanoparticles

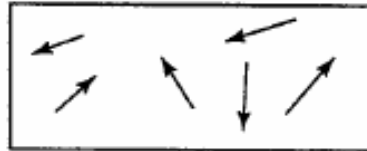
Chemical reduction !

6.2.4. Crystals of Metal Nanoparticles

A two-phase water toluene reduction of AuCl_4^- by sodium borohydride in the presence of an alkanethiol ($\text{C}_{12}\text{H}_{25}\text{SH}$) solution produces gold nanoparticles Au_m having a surface coating of thiol, and embedded in an organic compound. The overall reaction scheme is



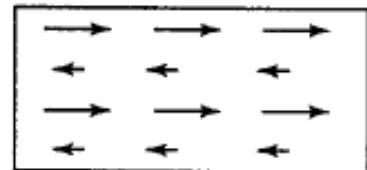
Chapter 7 Nanostructured Ferromagnetism



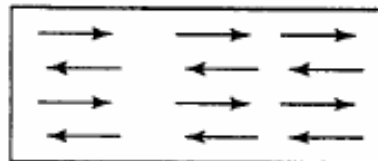
(a) PARAMAGNETIC



(b) FERROMAGNETIC



(c) FERRIMAGNETIC



(d) ANTIFERROMAGNETIC

ious arrangements of individual atomic

The interaction between atomic magnetic moments is of two types:

- 1. Exchange interaction

$$E = \int \left[\frac{1}{2} \frac{e^2}{r_{12}} \right] [\Psi_A(1)\Psi_B(2) - \Psi_A(2)\Psi_B(1)]^2 dV_1 dV_2 \quad (7.1)$$

which involves carrying out a mathematical operation from the calculus called *integration*. Expanding the square of the wavefunctions gives two terms:

$$E = \int \left[\frac{e^2}{r_{12}} \right] [\Psi_A(1)\Psi_B(2)]^2 dV_1 dV_2 - \int \left[\frac{e^2}{r_{12}} \right] \Psi_A(1)\Psi_B(1)\Psi_A(2)\Psi_B(2) dV_1 dV_2 \quad (7.2)$$

The first term is the normal Coulomb interaction between the two charged particles. The second term, called the *exchange interaction*, represents the difference in the

2. Dipolar interaction

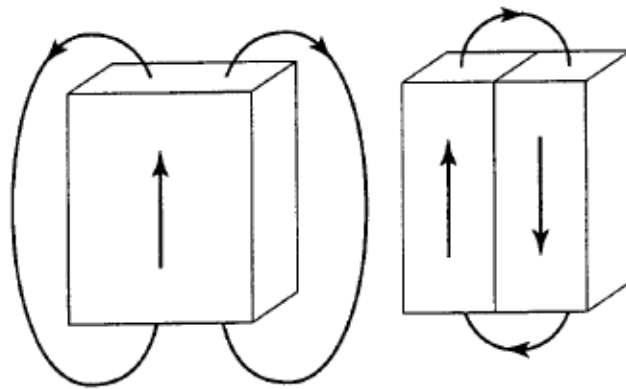
$$\frac{\mu_1 \cdot \mu_2}{r^3} - 3(\mu_1 \cdot \mathbf{r}) \frac{\mu_2 \cdot \mathbf{r}}{r^5} \quad (7.3)$$

Coulomb interaction and exchange interaction for $m_s = -1/2$

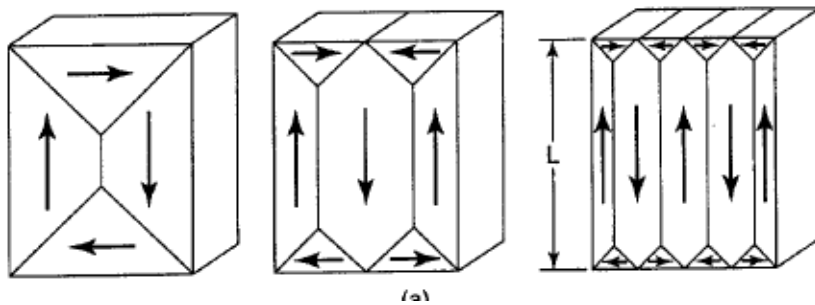
Magnetization M of a bulk

$$M(T) = M(0)(1 - cT^{3/2}) \quad (7.4) \quad \text{For } T \ll T_c$$

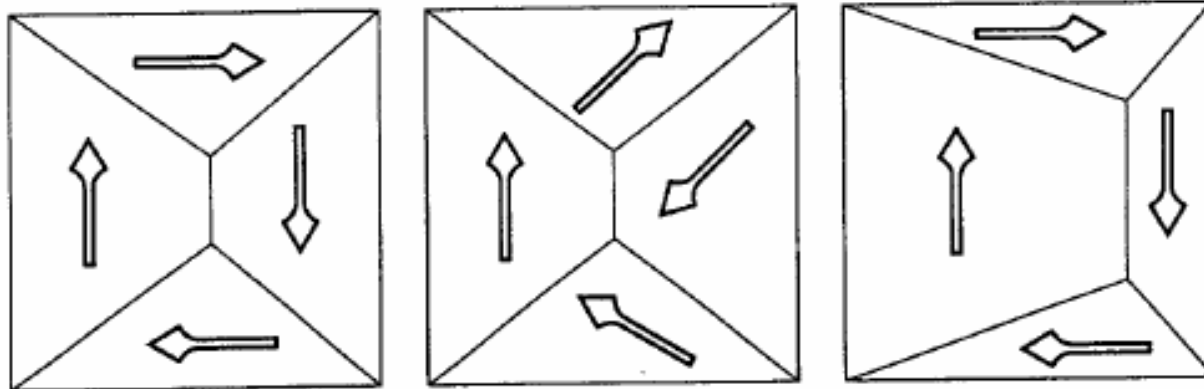
$$\chi = M/H.$$



$H=0$



(a)



INITIAL
CONFIGURATION

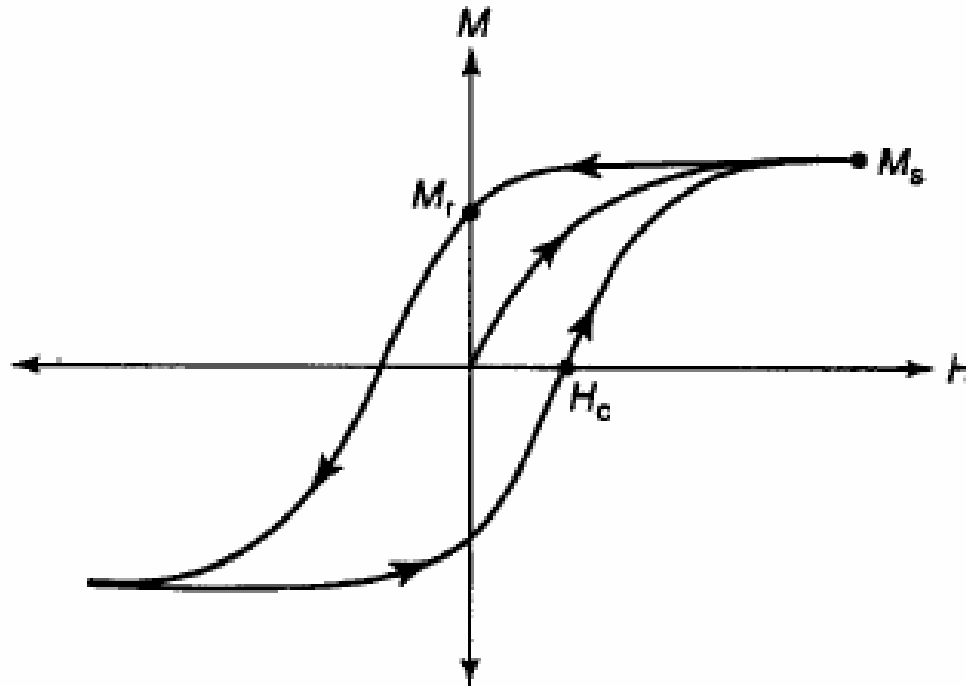
DOMAIN
ROTATION

DOMAIN
GROWTH

(b)

$H>0$

Remnant magnetization, Coercive field, Saturation magnetization



7.3. Plot of the magnetization M versus an applied magnetic field H for a hard magnetic material, showing the hysteresis loop with the coercive field H_c , the remnant magnetization M_r , and the saturation magnetization M_s , as indicated.

7.2 Effect of bulk nanostructuring on magnetic properties

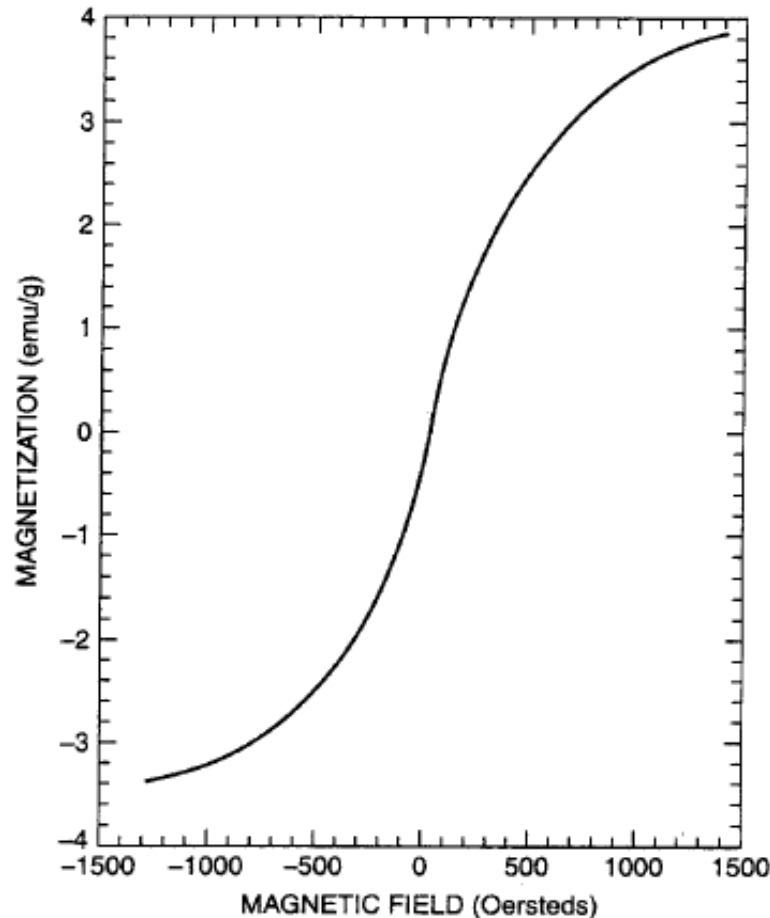


Figure 7.4. Reversible magnetization curve for nanosized powders of a Ni-Fe-Co alloy exhibits no hysteresis. An oersted corresponds to 10^{-4} T (tesla). [Adapted from K. Shafi et al. *Appl. Mater. Res.* 15, 332 (2000).]

Superparamagnetism

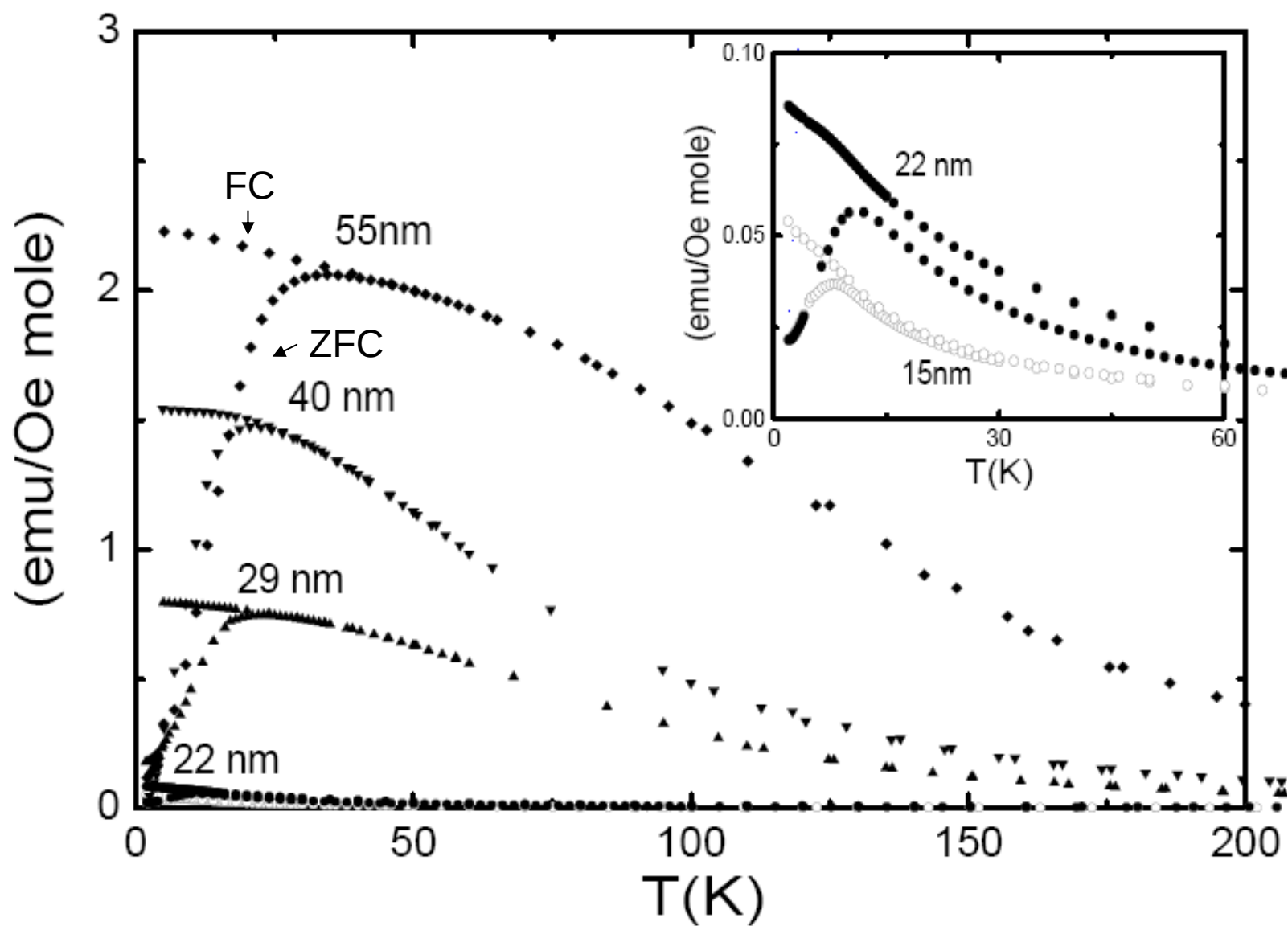
⇒ Monodomain particles

- Below 100 nm

⇒ Fluctuation of the magnetic moment like in a paramagnet

⇒ Ferromagnetic particles with moment (T_c is high)

⇒ Moment dependent on particle volume



此圖為FeSi₂奈米粉末的DC磁化率-溫度曲線

Temperature dependence of χ

- 1. The temperature of peak value of χ in ZFC is defined as the Blocking temperature T_{B0}
- 2. χ of ZFC and χ of FC deviate at T_B
- 3. Above T_B , χ of ZFC and χ of FC are overlap.

Blocking Temperature

$$T_B = \frac{KV}{25k_B}$$

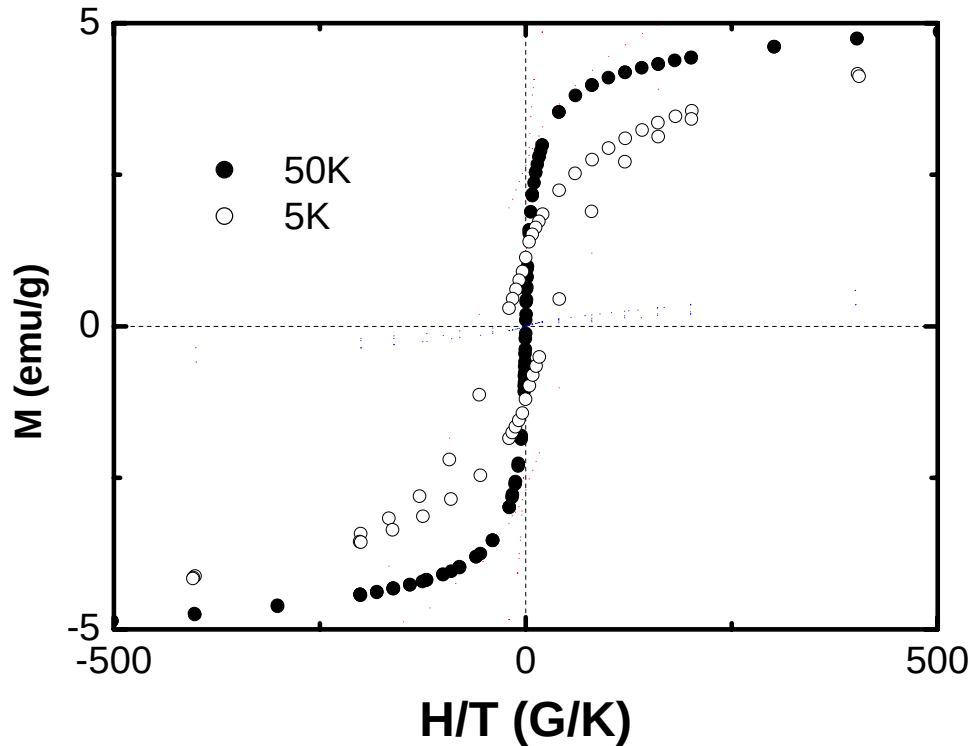
k_B is the Boltzmann constant

K is the anisotropic constant

V is the volume of nanoparticle

Analysis of size-dependent blocking temperature

M-H曲線



- 1. $T < T_B$, Hysteresis appears in M-H. Due to thermal energy is less than the interactions among particles
- 2. $T > T_B$, No hysteresis appears in M-H. Since thermal energy is larger than the interactions among particles

FeSi_2 40nm particles

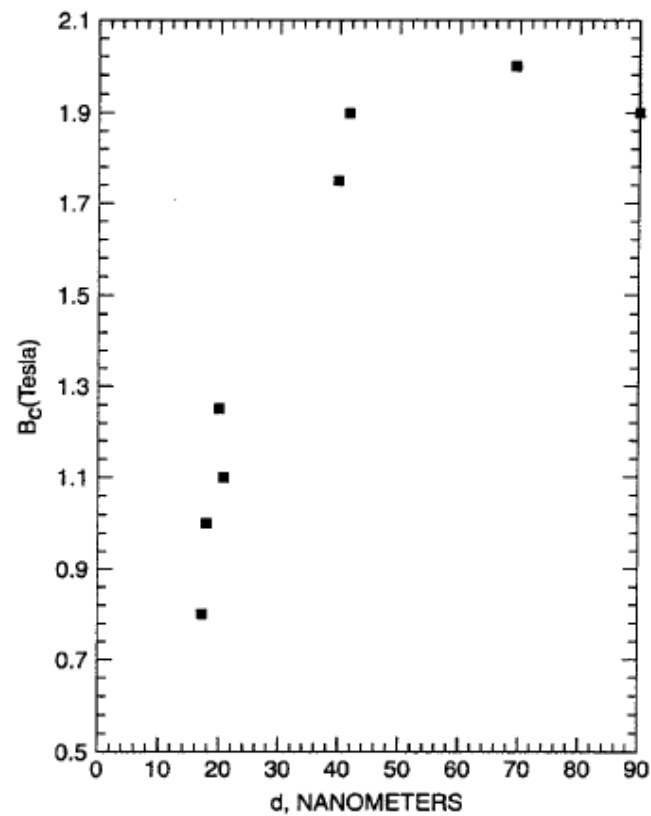
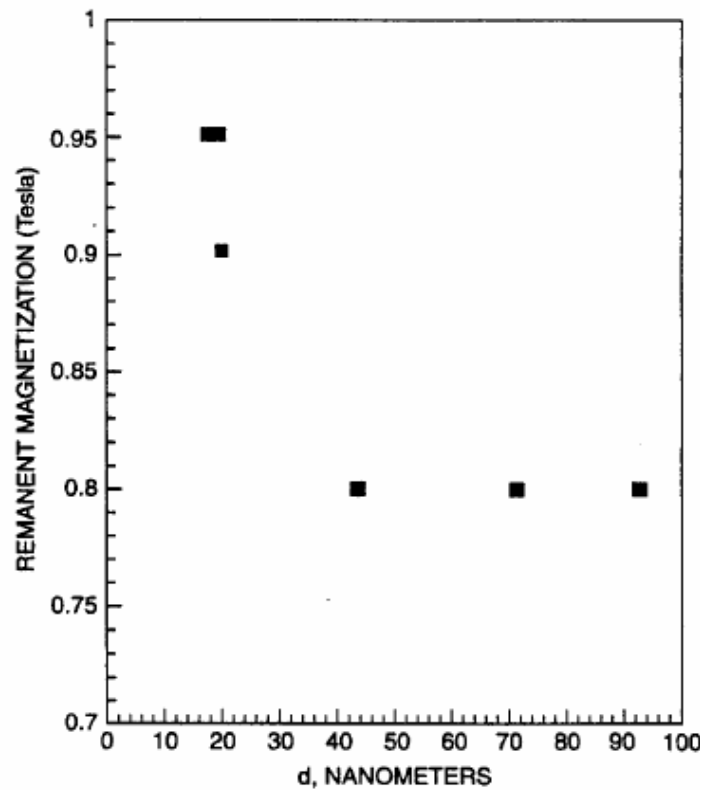


Figure 7.6. Dependence of the coercive field B_c (i.e., H_c) on the granular particle size d of Nd-B-Fe permanent magnet. [Adapted from A. Manaf et al., *J. Magn. Magn. Mater.* 1

$\text{Nd}_2\text{Fe}_{14}\text{B}$ Grain d

7.4 nanopore containment of magnetic particles

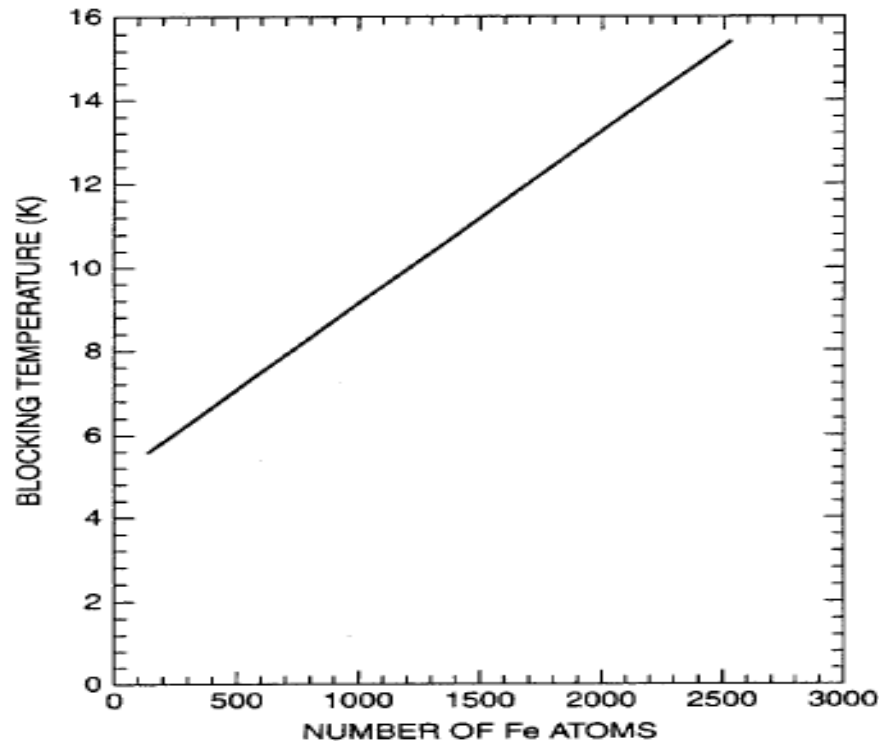


Figure 7.9. Plot of blocking temperature T_B versus the number of iron atoms in the cavity of ferritin. [Adapted from D. D. Awschalom and D. P. DiVincenzo, *Phys. Today* (April 1995).]

7.5 Nanocarbon ferromagnets

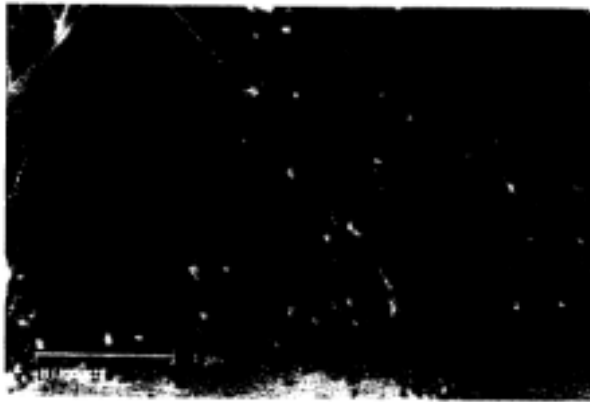
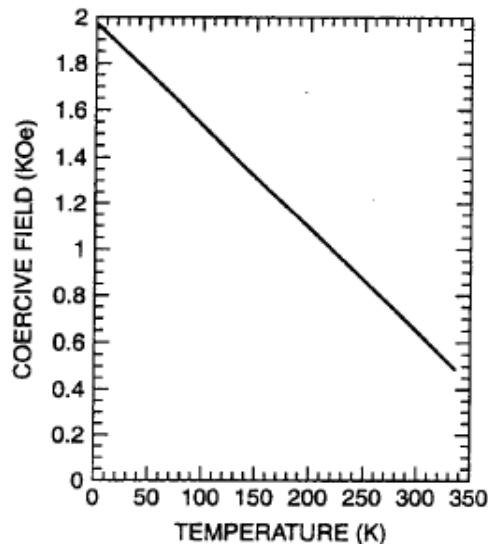
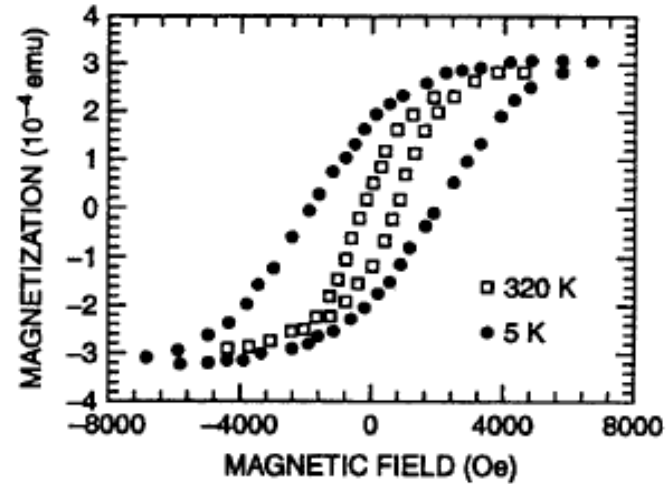
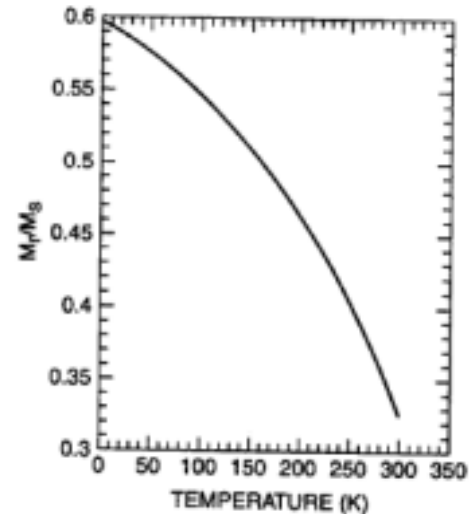


Figure 7.11. Scanning electron microscope image of iron particles (light spots) on the tips of ignited carbon nanotubes. [With permission from Z. Zhang et al., *J. Magn. Mater.* 231, 1 (2001).]



Plot of coercive field H_c versus temperature T for iron particles on the tips of carbon nanotubes.

Magnetization curve hysteresis loops for iron particles on the tips of carbon nanotubes.



Plot of remnant magnetization M_r to saturation magnetization M_s versus temperature T for iron particles on the tips of carbon nanotubes.

Quantum size effects 量子尺寸效應：

Phonon quantum size effect： 聲子與表面積之改變。

Electronic quantum size effect： 電子結構改變與能階劈裂。

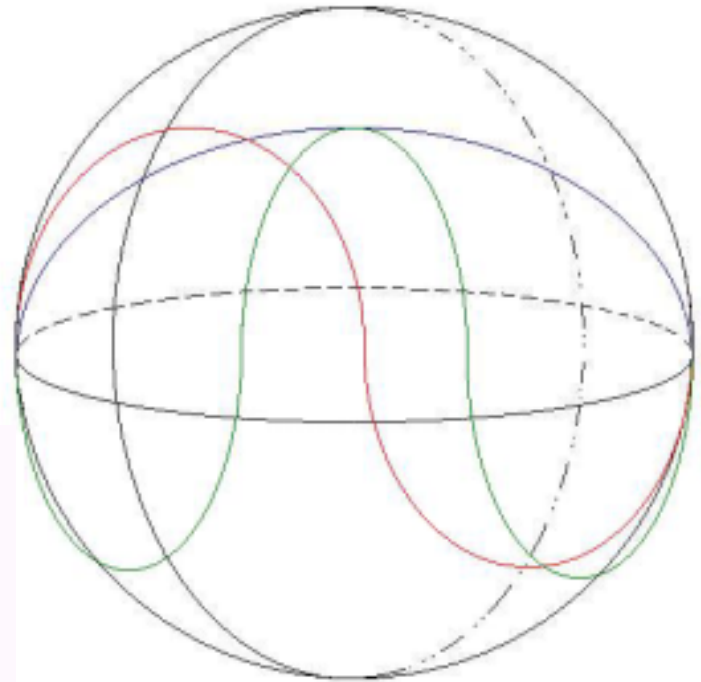
Phonon quantum size effect : 聲子與表面積之改變。

Wave mode:
Spherical Bessel Function

$$C_p = \gamma T + \sum_l \sum_s \frac{V_m (2l+1) K_B z^2 e^z}{4R^3 (e^z - 1)^2}$$

with

$$z = \frac{\hbar c a'_{l,s}}{K_B R T}$$



Nanoparticle :

聲子量子尺寸效應：聲子與表面積之改變。

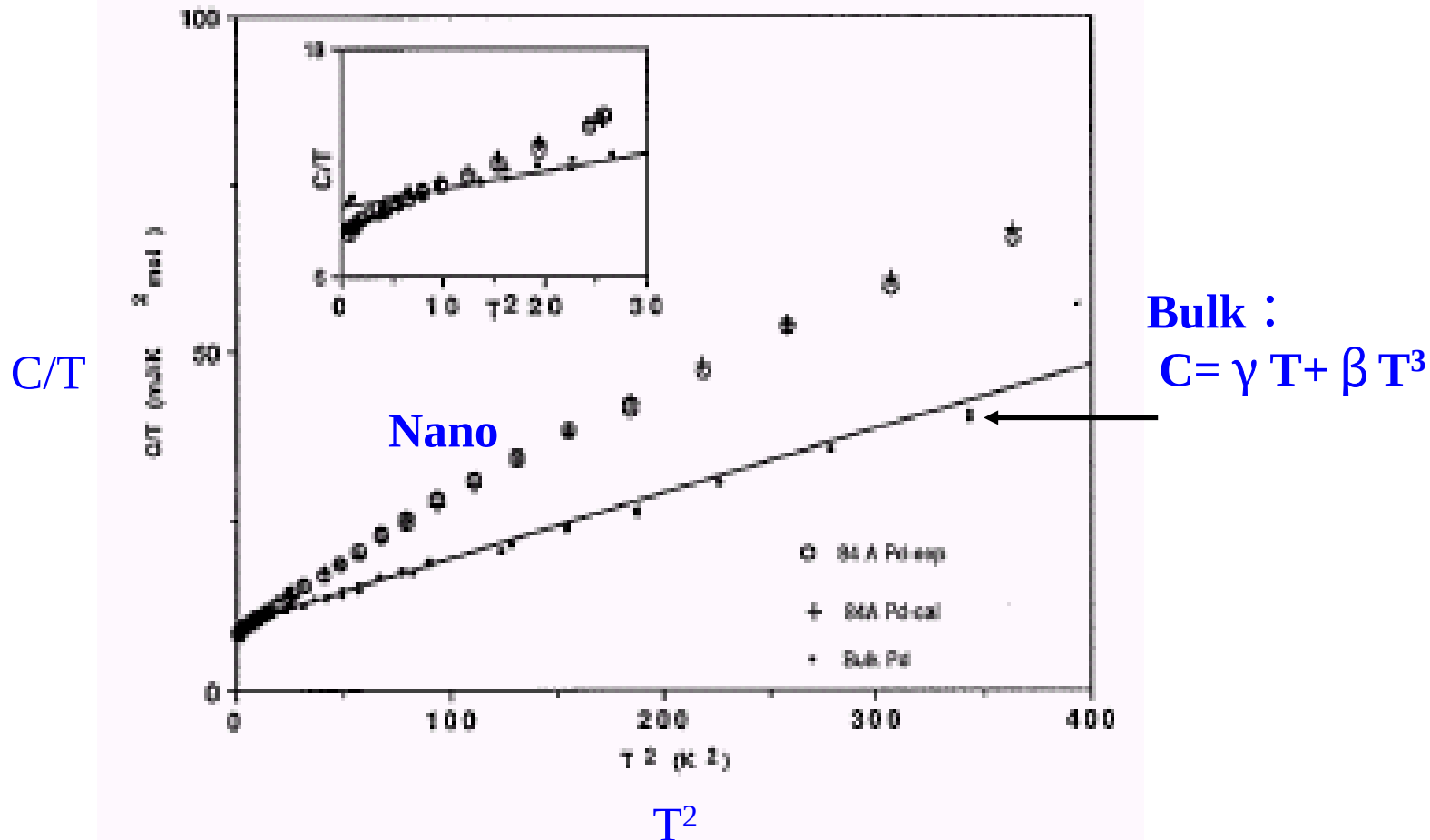
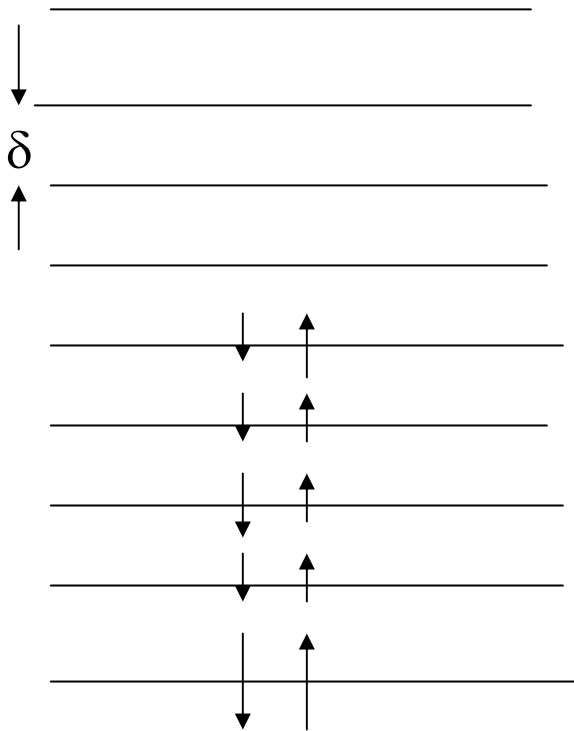


FIG. 3. The specific heat of bulk palladium and 84-Å Pd nanocrystals. The small dots represent bulk Pd and the solid

Y.Y. Chen etc, PRB 52, 9364, 1995

Electronic quantum size effect : 電子結構改變與能階劈裂。



$$\epsilon_F \sim 10^4 \text{ K}$$

$$\delta \approx E_F / N \propto V^{-1}$$

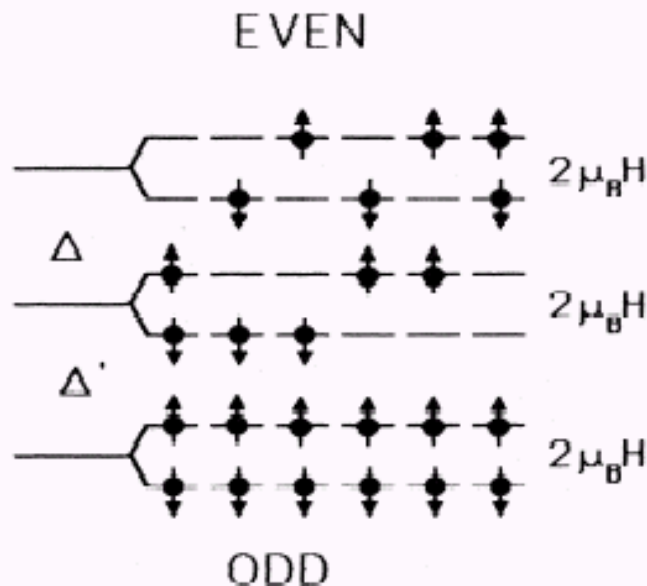
Bulk Pd $\delta \sim 10e^{-30}\text{K}$

80 Å Pd $\delta \sim 0.5\text{K}$

No. of atoms ~ 8000

Magnetic field

$H=0$



$H>0$

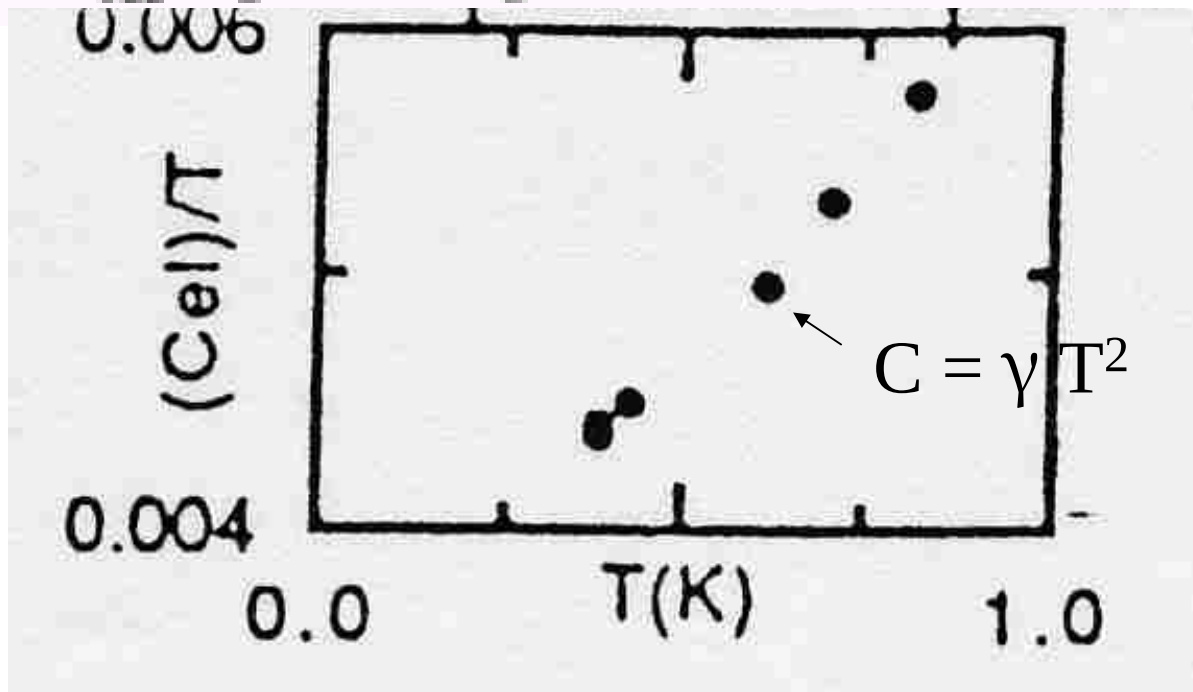
FIG. 2. Electron-level structure diagrams for the two cases of particles with an even number of electrons, and for particles with an odd number of electrons. Only the lowest-energy configurations are shown with the ground state at the left and progressively higher excited states to the right for each of the even and odd cases.

$$Z_{\text{even}} \approx 1 + 2(1 + \cosh 2\beta\mu_B H)(e^{-\beta\Delta} + e^{-2\beta\Delta}) ,$$

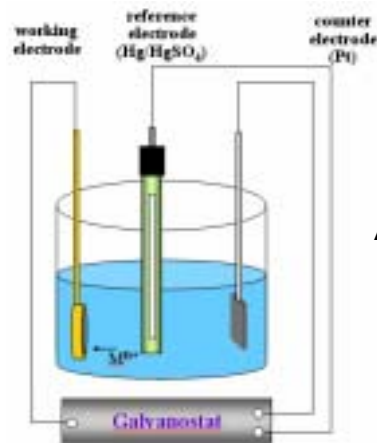
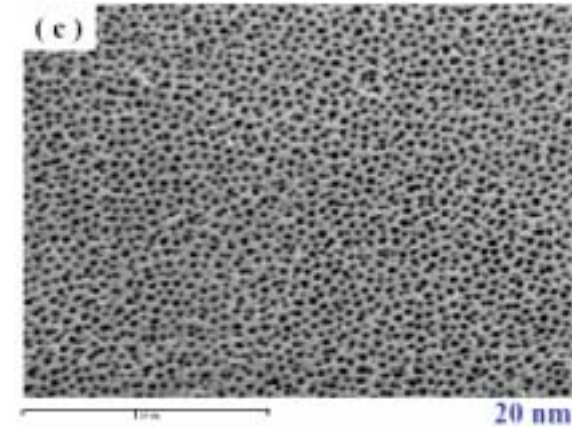
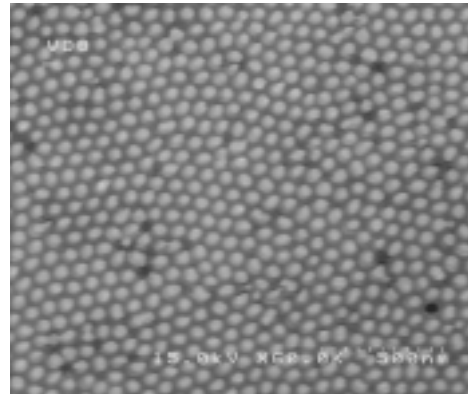
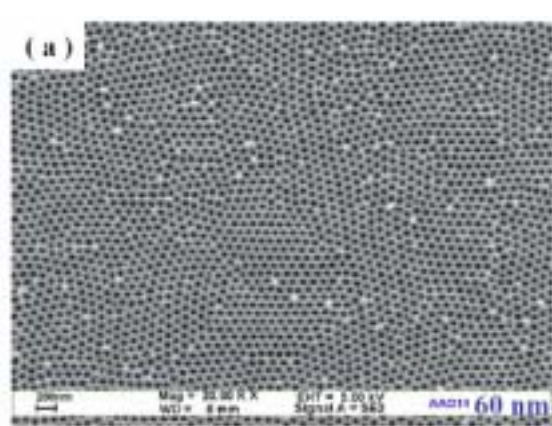
$$Z_{\text{odd}} \approx 2(\cosh \beta\mu_B H)(1 + e^{-\beta\Delta} + e^{-\beta\Delta'}) .$$

$$C_{\text{even}}^1/k_B = 30.2(k_B T/8)^2 ,$$

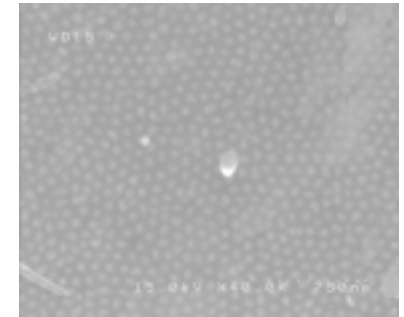
$$C_{\text{odd}}^1/k_B = 17.8(k_B T/8)^2 , \quad \text{orthogonal}$$



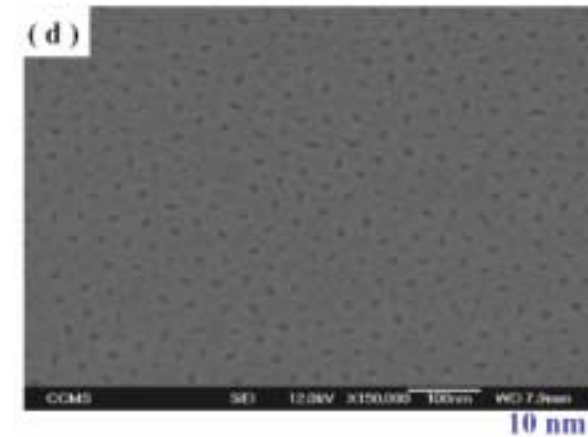
Size dependence of magnetization in Fe nanowires



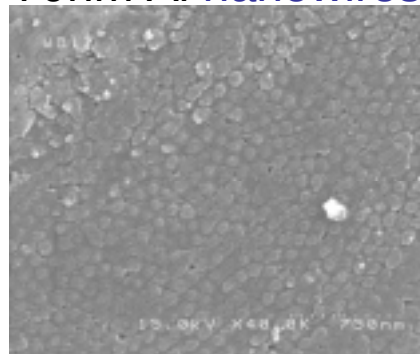
AAO method



Cu nanowire

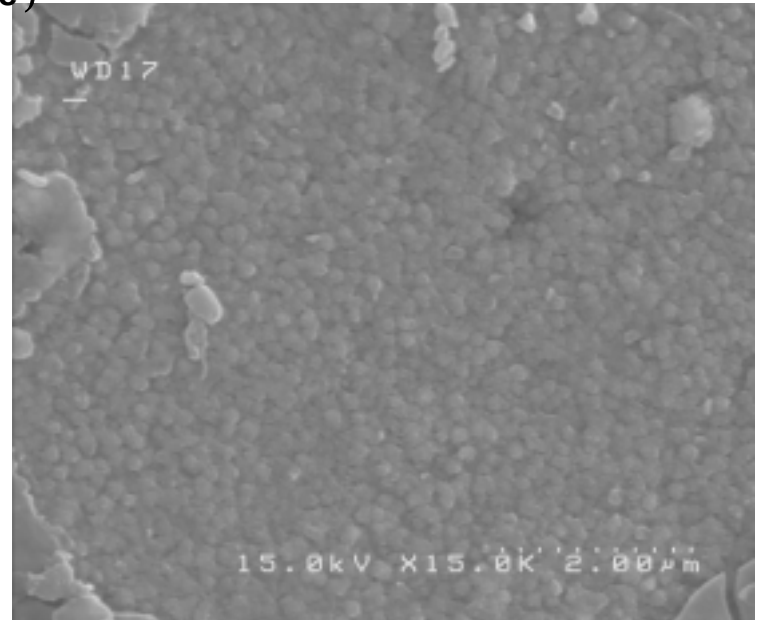
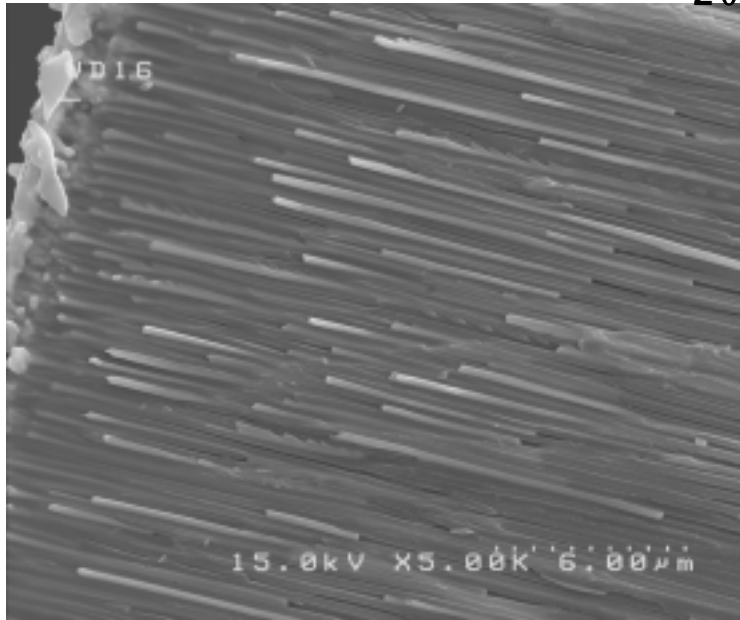
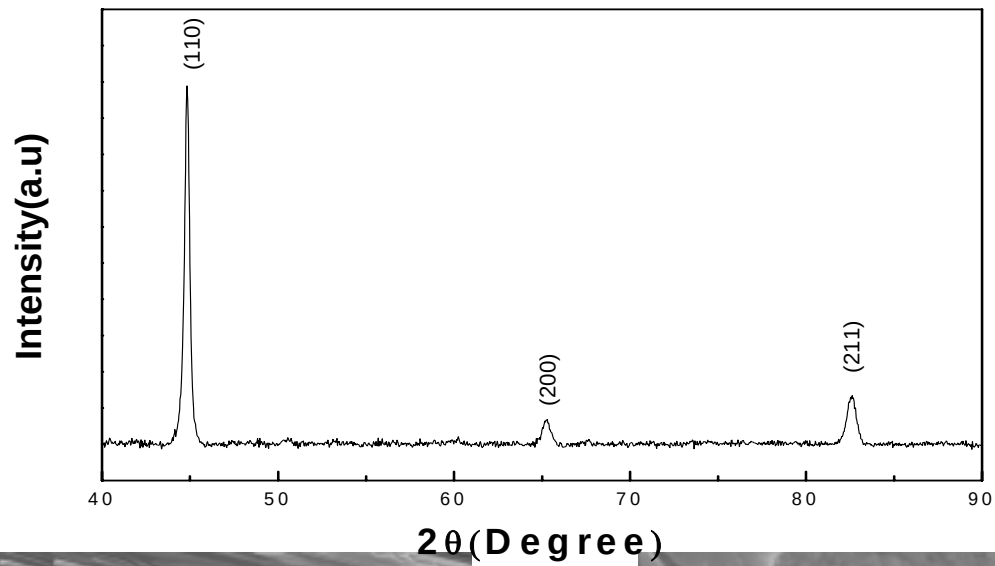


70nm Al nanowires



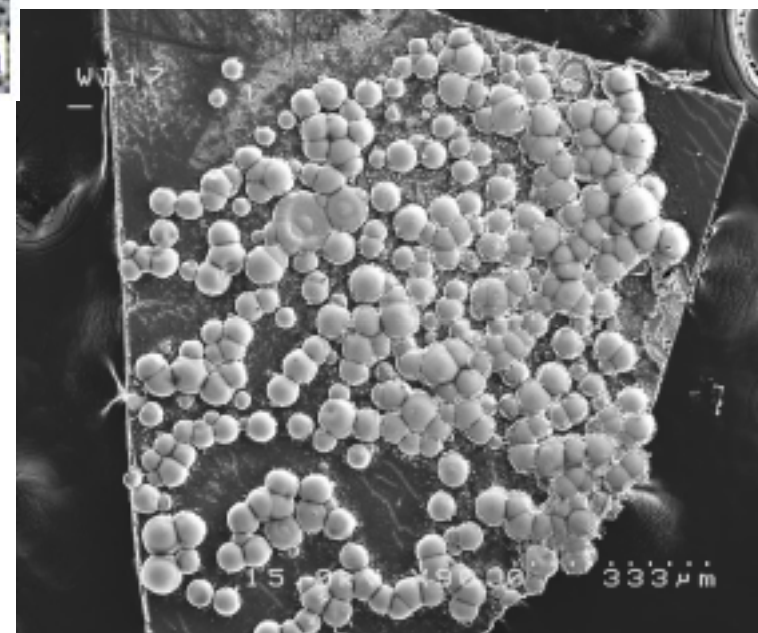
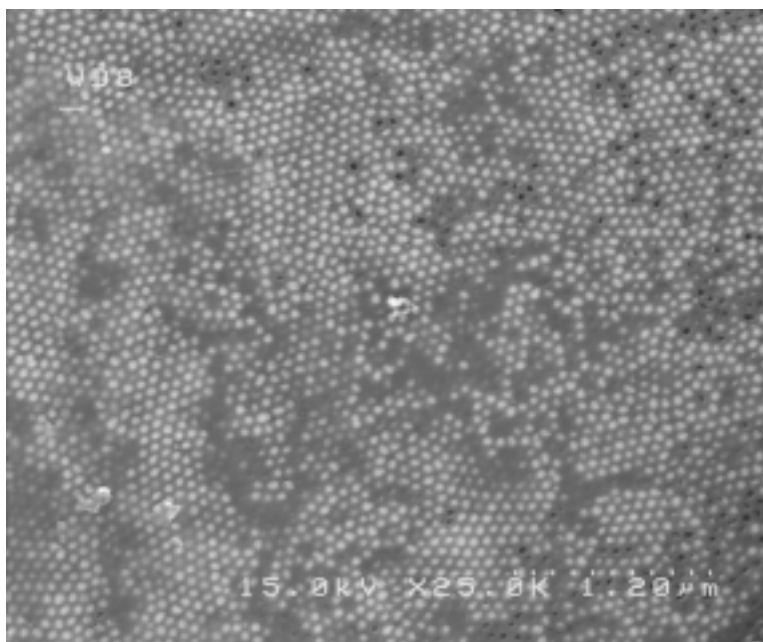
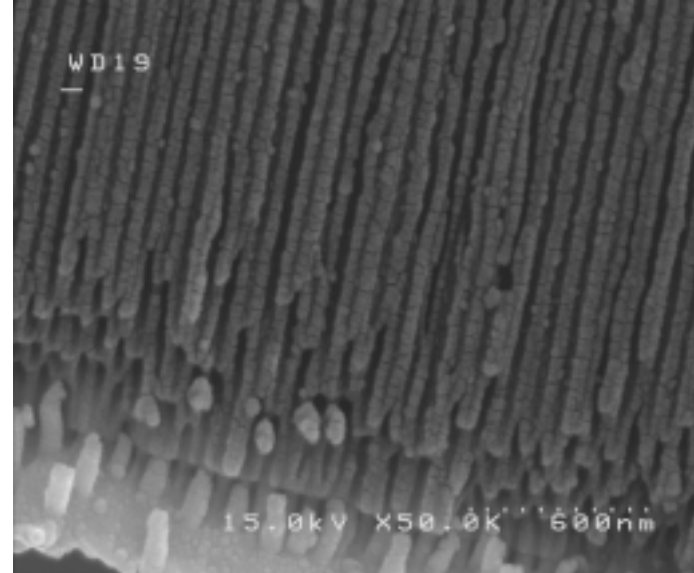
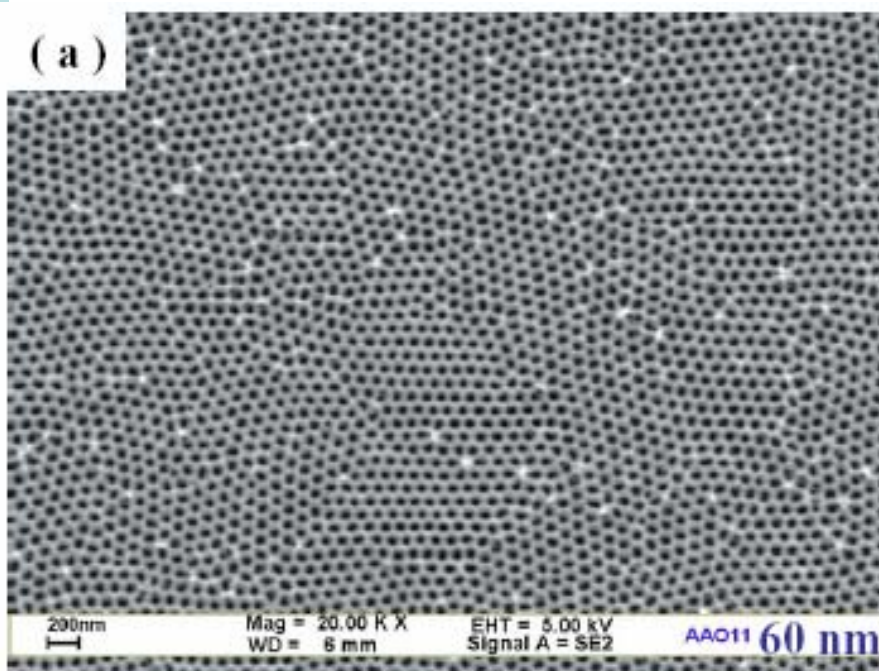
$\text{CuSO}_4 \cdot 5\text{H}_2\text{O}$, pH~2

Fe



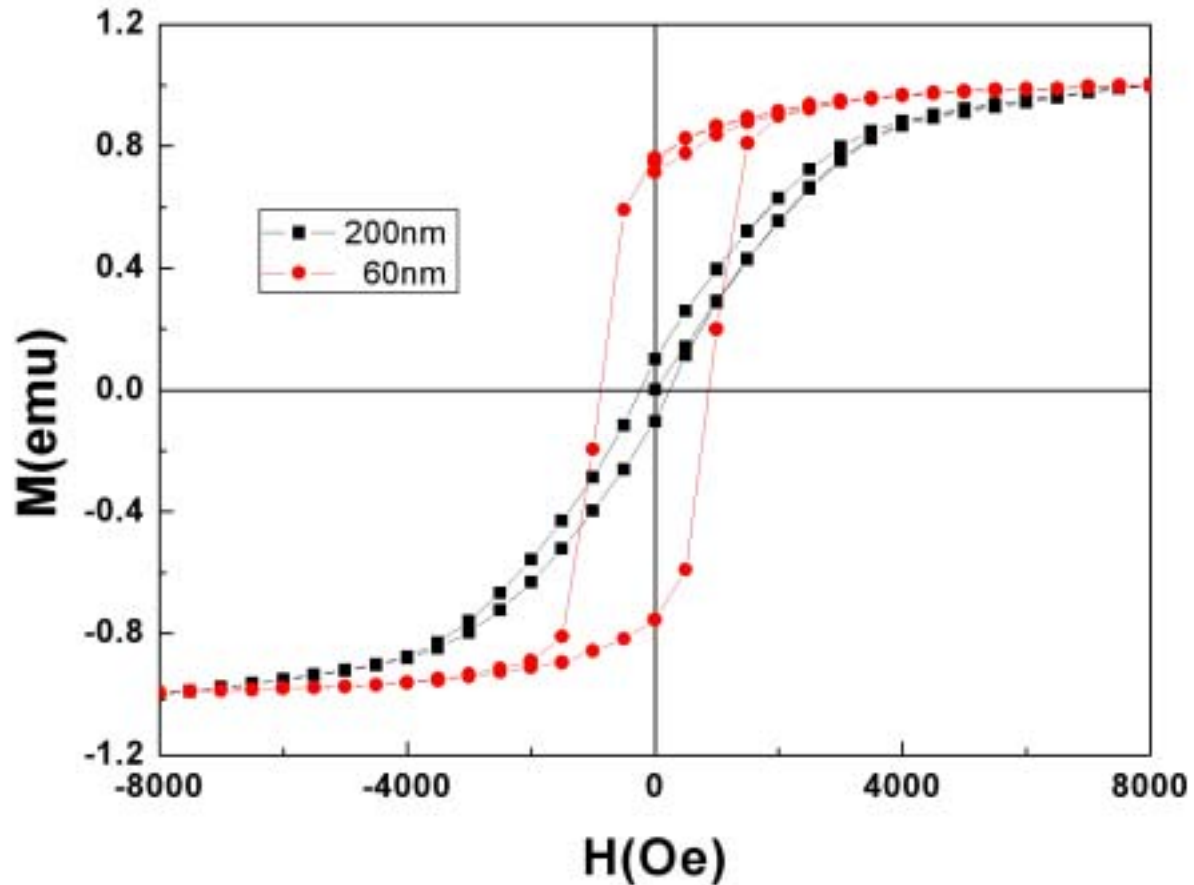
Fe nanowire ~200 nm

(a)

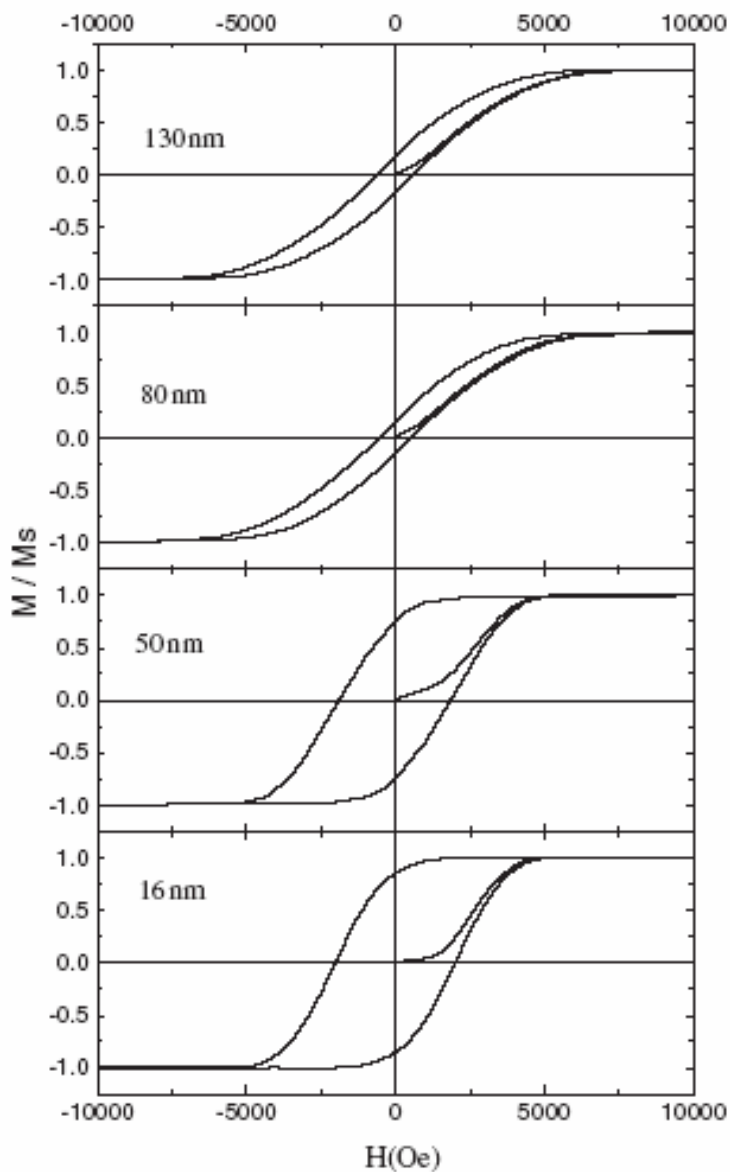


Fe nanowire ~60 nm

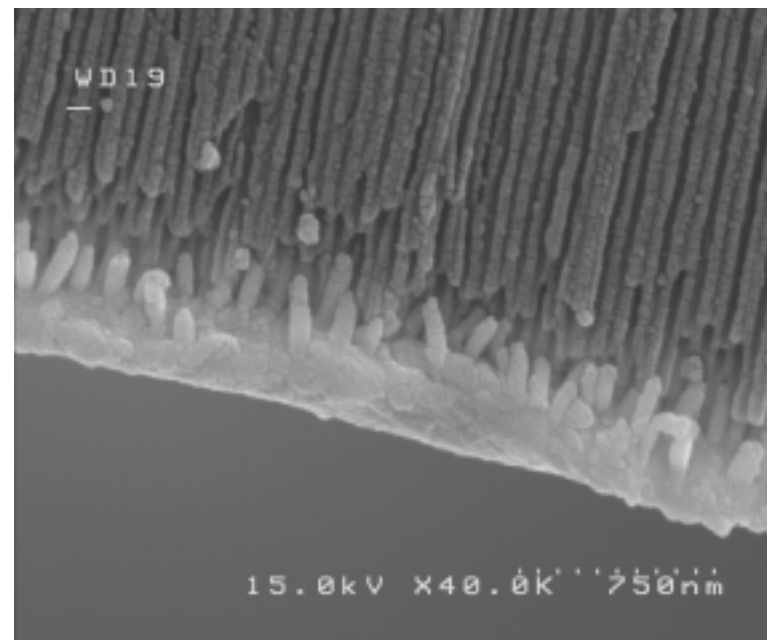
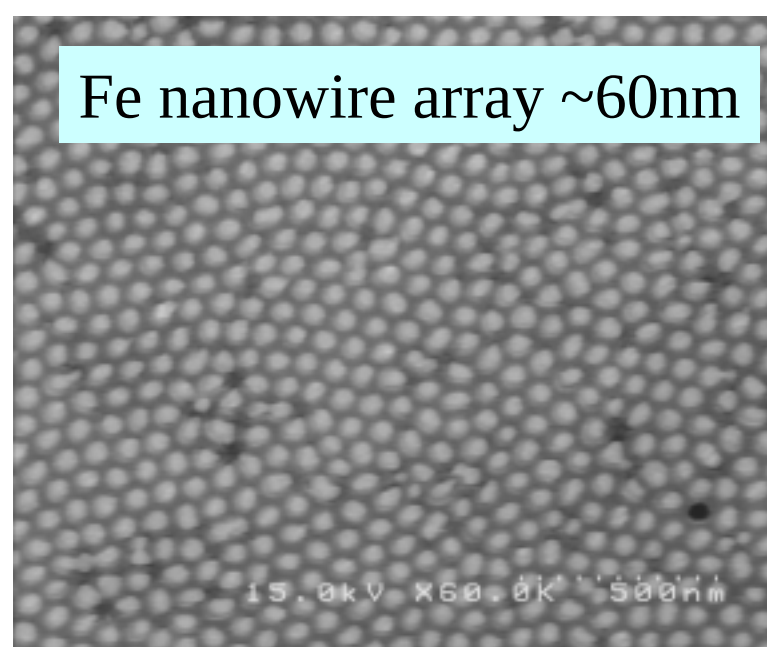
Normalized (M/M_s) -H



Magnetic texture of nanowire arrays



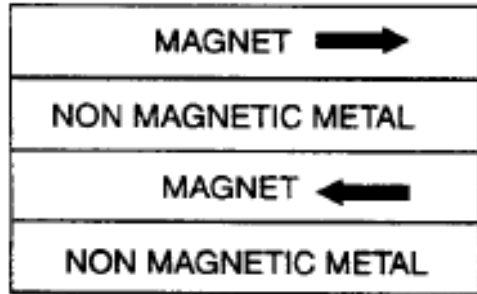
Fe nanowire array $\sim 60\text{nm}$



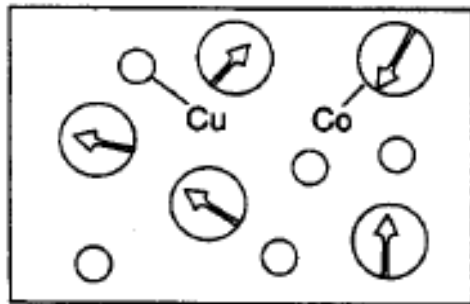
7.6 Giant and colossal magnetoresistance

- Metal:
- 1. The conduction electrons being forced to move in helical trajectories about an applied magnetic field.
- 2. Field curves the electron trajectory within a length of its mean free path
- Cu at 4 K with $H=10$ T
- R increases 10 times

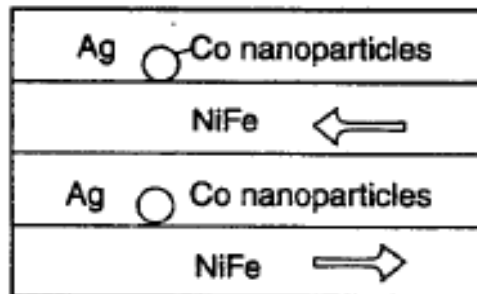
Giant magnetoresistance (GMR)



(a)

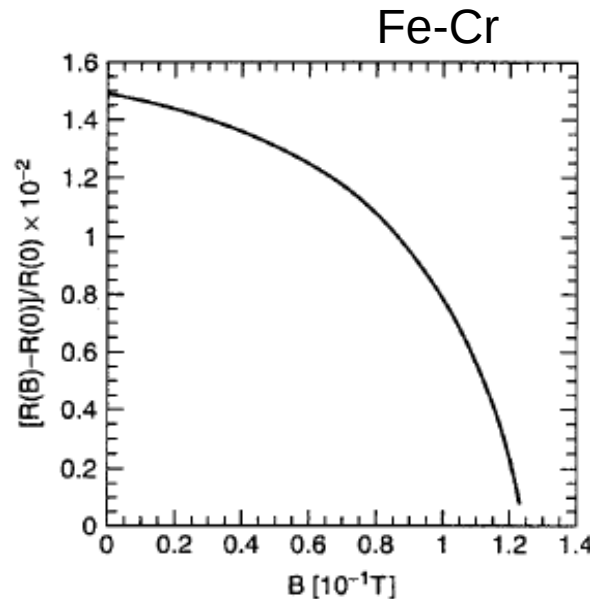


(b)

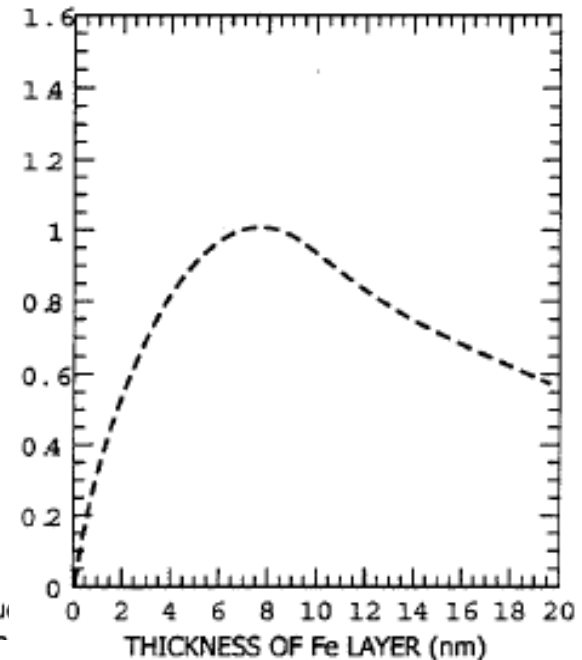


(c)

Alternate layers of FM and Non-FM



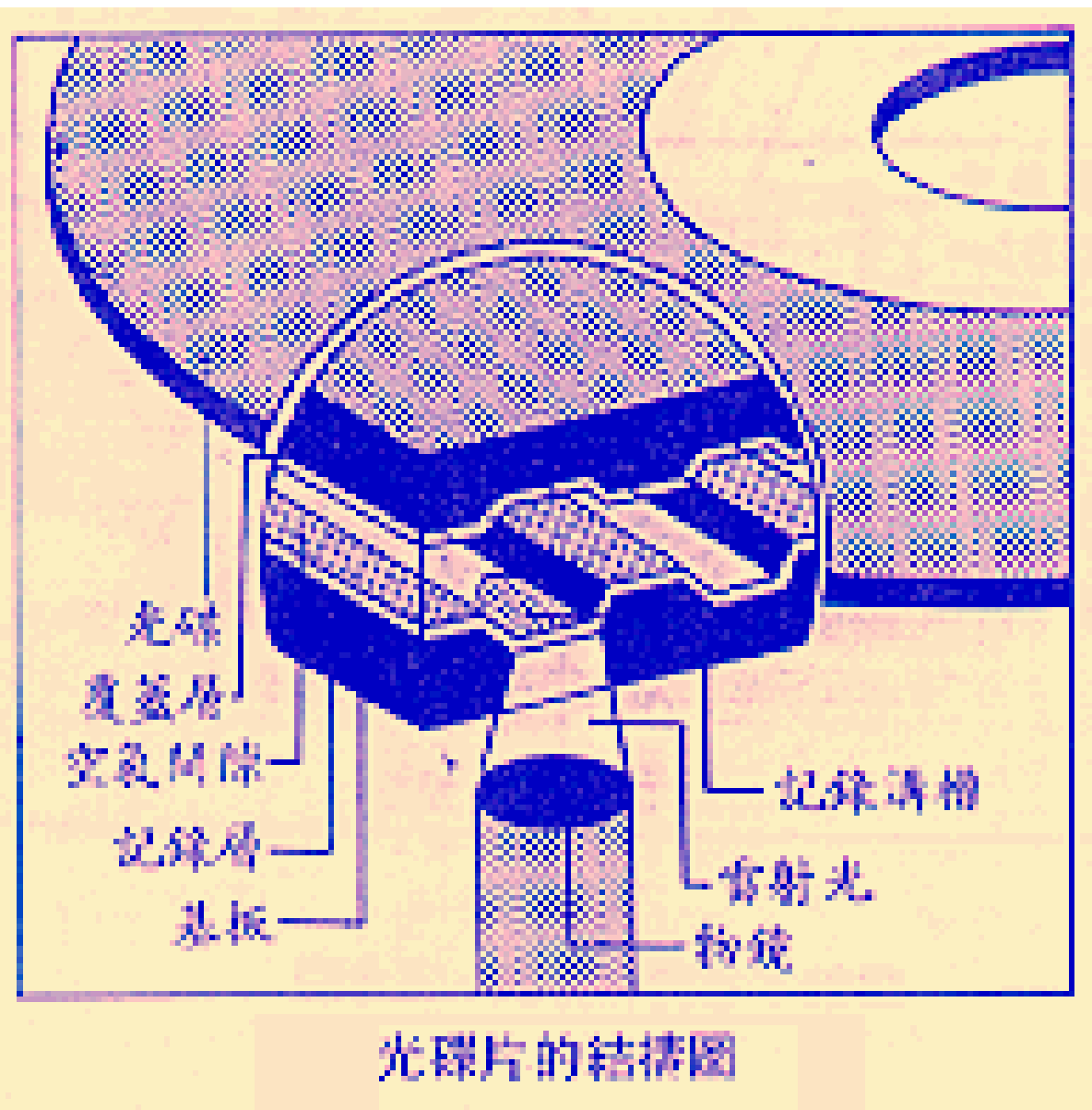
dependence of the electrical resistance $R(B)$, relative to its value at zero field, on a magnetic field B applied parallel to the current.

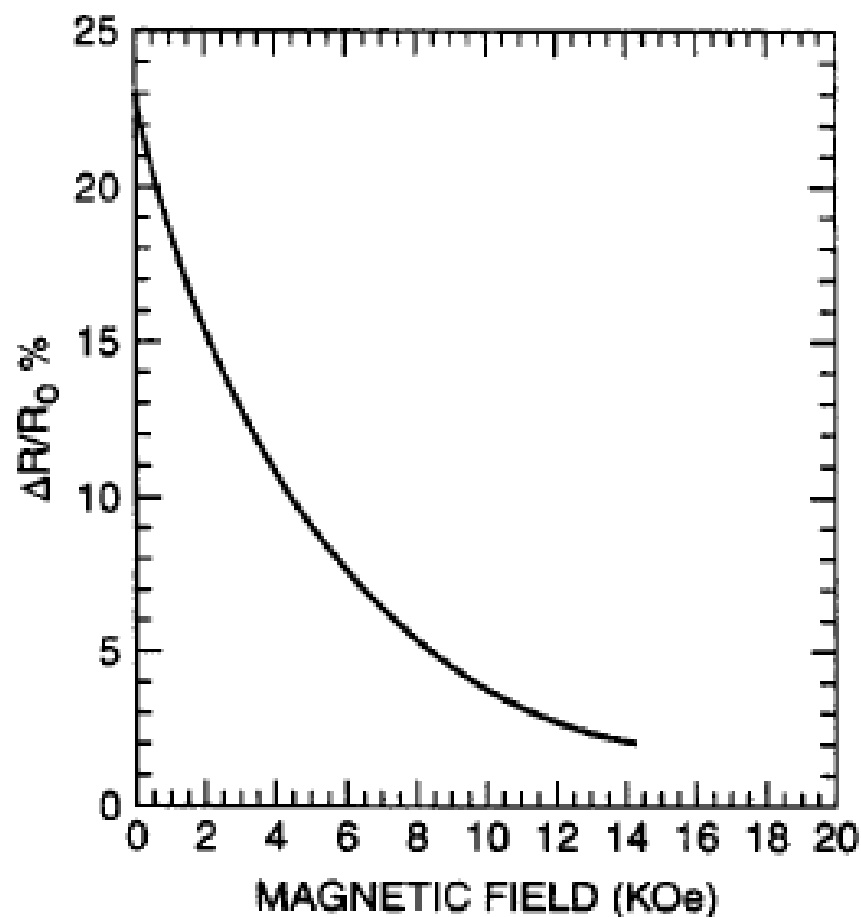


Electron // H has less scattering

Magnetic storage

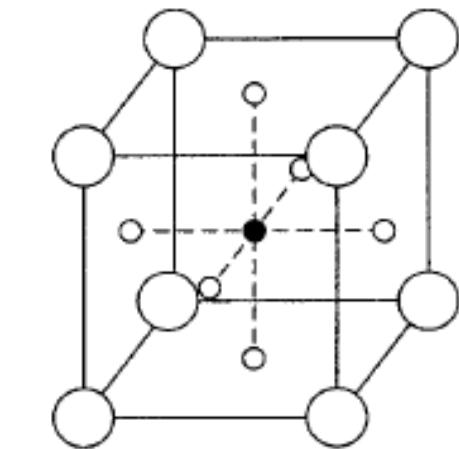
- 1. **Induction coils** to induce and read the magnetization
- 2. The **magnetoresistive reading** is more sensitive than **Induction coils**



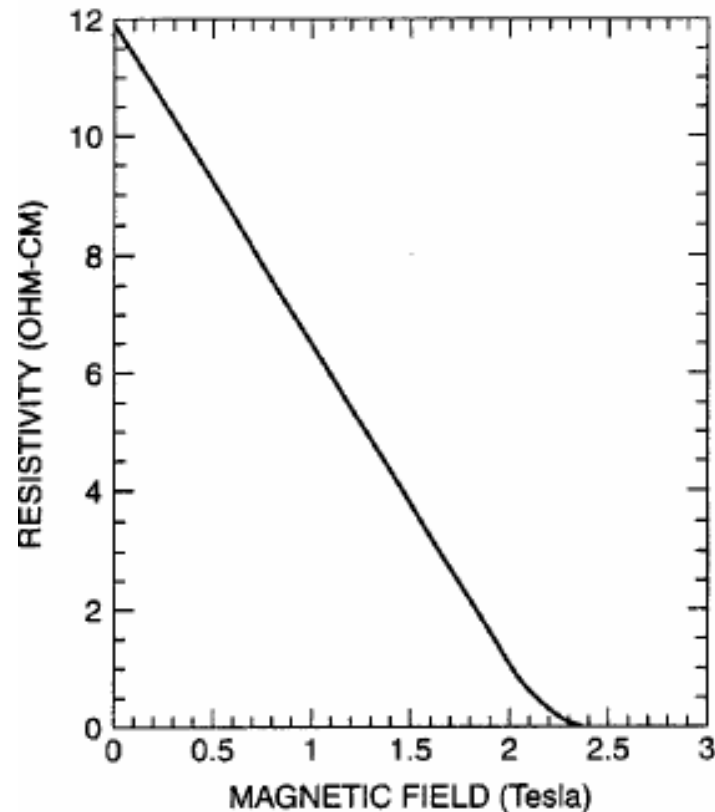


Dependence of the change of magnetoresistance ΔR versus the film of Co nanoparticles in a copper matrix. A kilooersted coil
 A. F. Radtsig, Phys. Rev. Lett. 88, 0745 (2002).

Colossal magnetoresistance

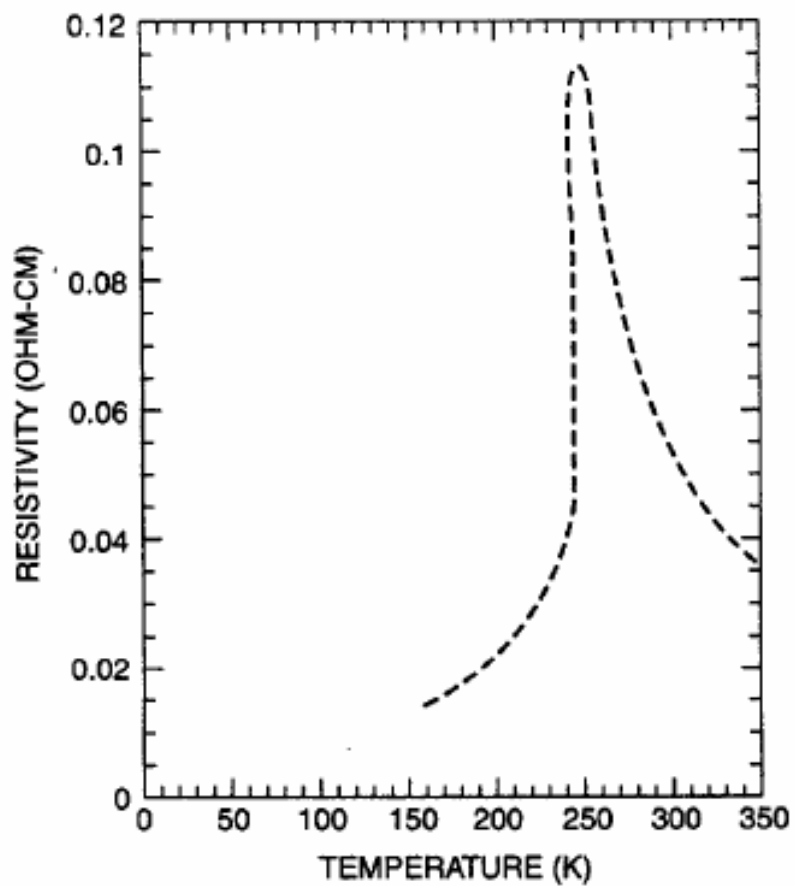


- (La,Ca Sr)
- (Mn)
- (O)



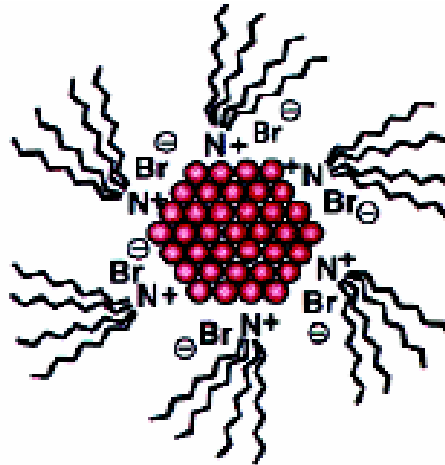
ence of the resistivity (normalized magnetoresistance) of La-Ca-Mn-O on field in the neighborhood of the Curie temperature at 250 K. (With

La-Ca-Mn-O



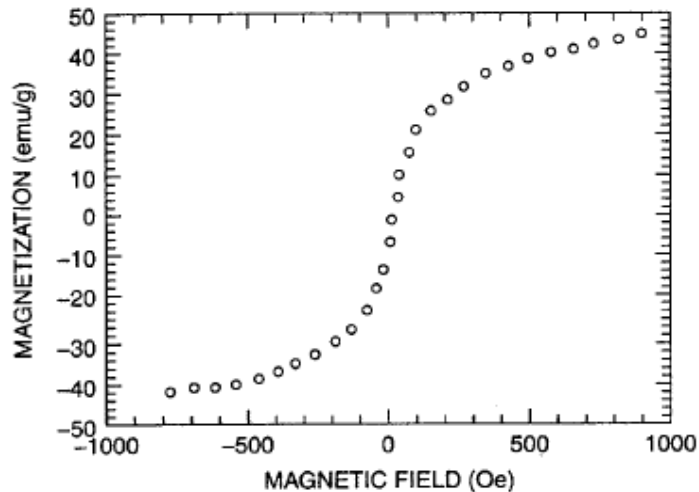
7.7 Ferrofluids

- 1. Nanoparticles are single domain ~ 10 nm
- 2. Coated with a surfactant to prevent aggregation
- 3. Suspended in a liquid of transformer oil or Kersene

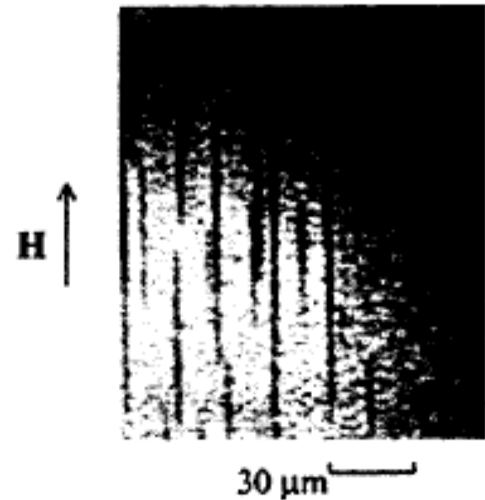


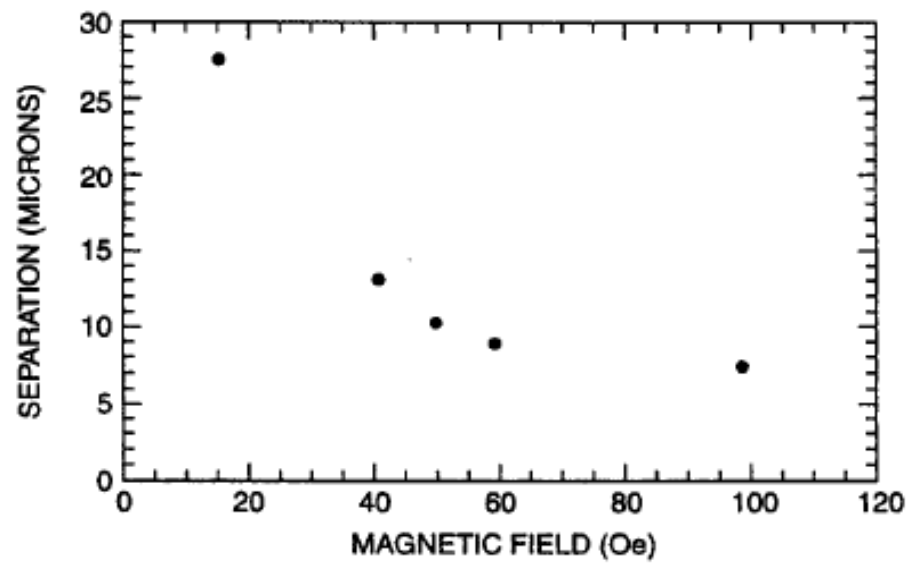
Magnetic-field-dependent anisotropic optical properties

7.7. FERROFLUIDS 187



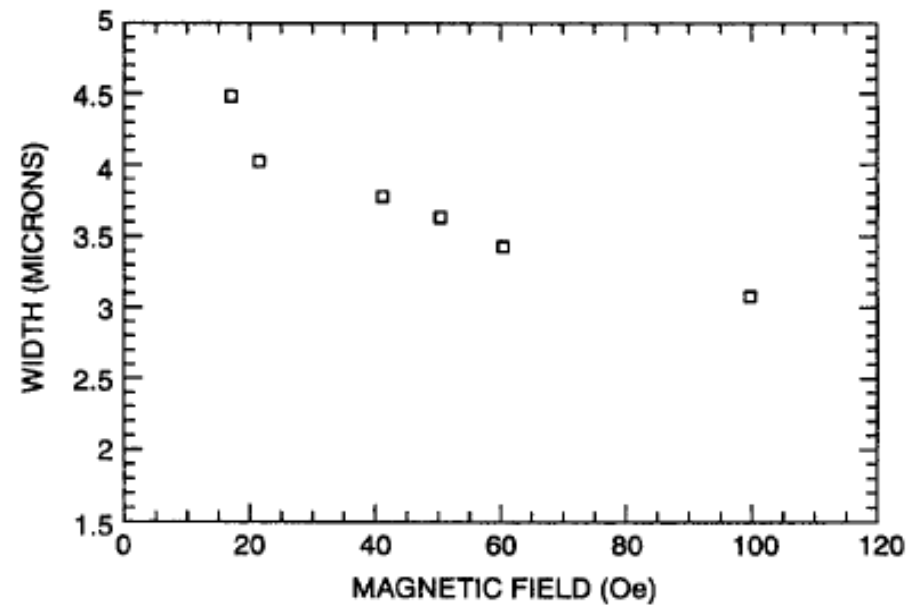
7.22. Magnetization curve for a ferrofluid made of magnetite, Fe_3O_4 , nanoparticles





(a)

H // Surface



(b)

H is perpendicular to the surface

7.7. FERROI

