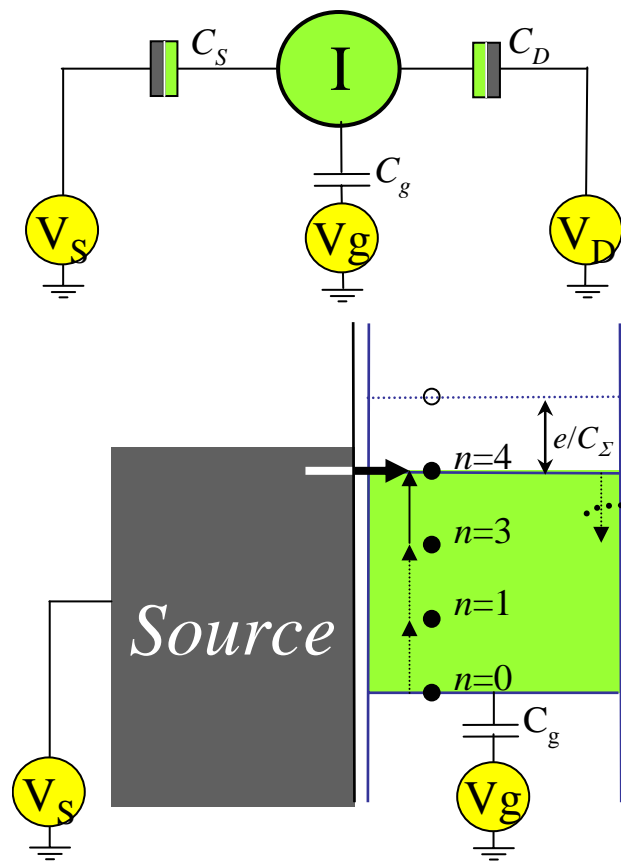
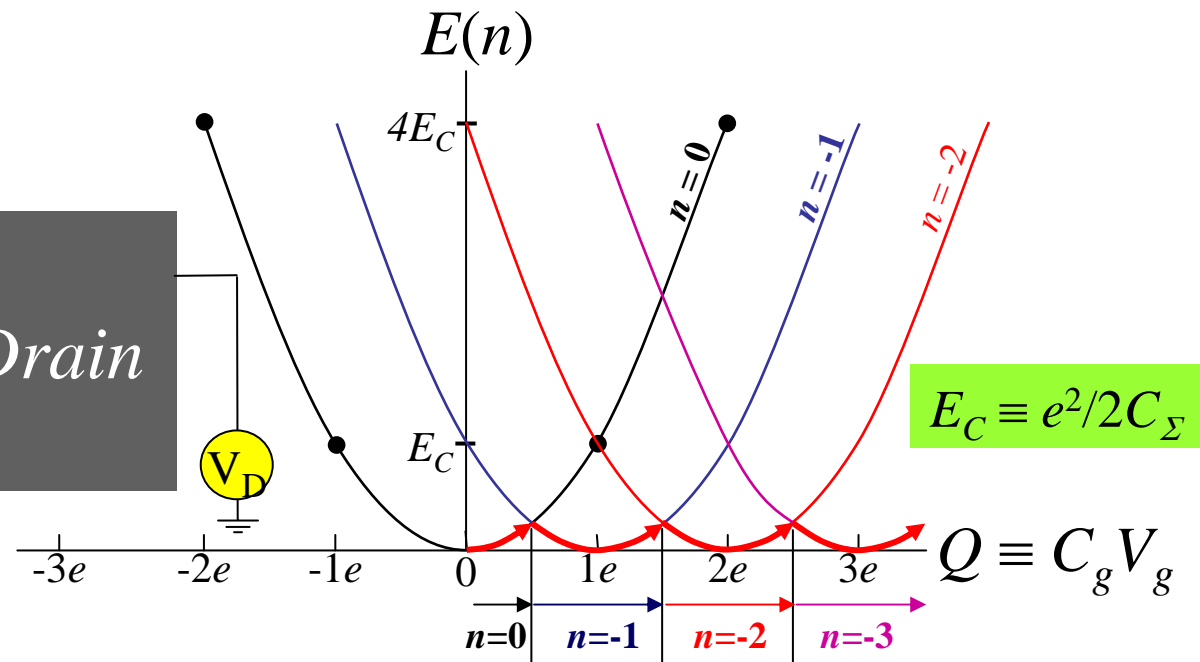


Single Electron Transistor



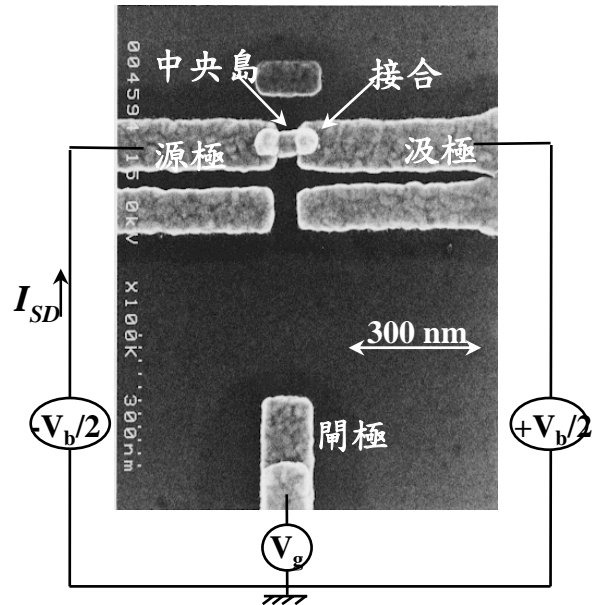
$$E(n) = (ne + C_g V_g)^2 / 2C_\Sigma$$

$$C_\Sigma = C_S + C_D + C_g$$



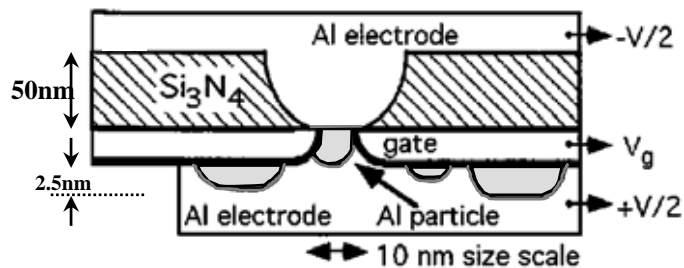
Various Single electron transistors:

(1) All metal transistors: Al island



our work (and many other groups)

(2) Sandwich structure: Al particle

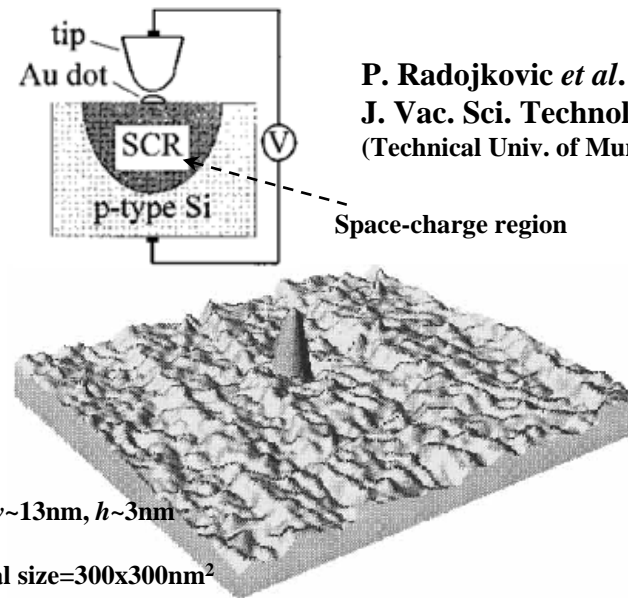


Al particles are isolated by thin Al_2O_3 layer

D.C.Ralph, C.T.Black and M. Tinkham
PRL, 78, 21 p. 4087 (1997) (Harvard Univ)

(3) STM: on Au particle

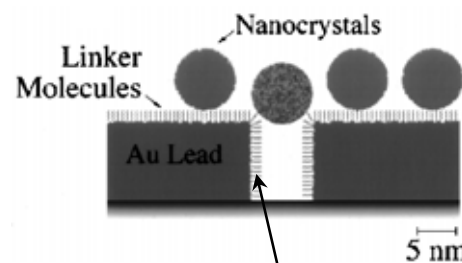
(STM=Scanning Tunneling Microscope)



P. Radojkovic *et al.*

J. Vac. Sci. Technol. B 14(2), 1229 (1996)
(Technical Univ. of Munich)

(4) Colloidal: Au or CdSe particles



self-assembled molecular
of dithiol molecules

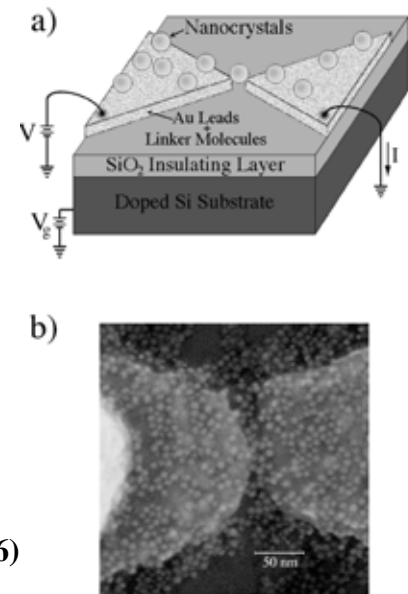
David L. Klein *et al.*

Nature, 389, p.699 (1997)

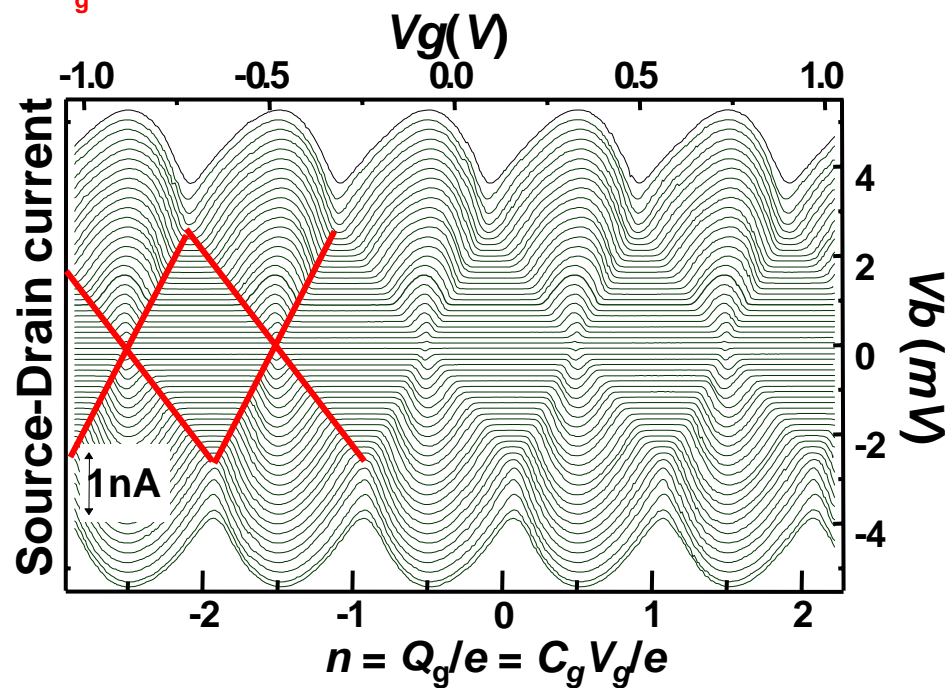
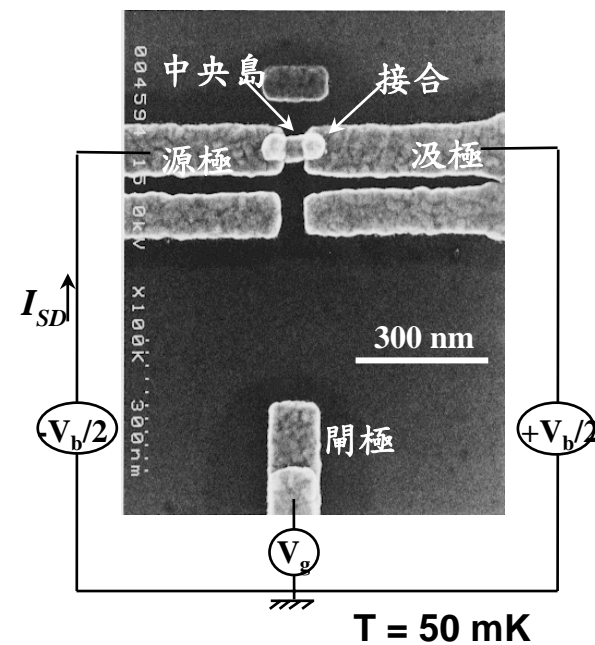
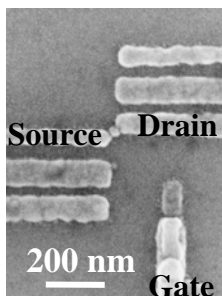
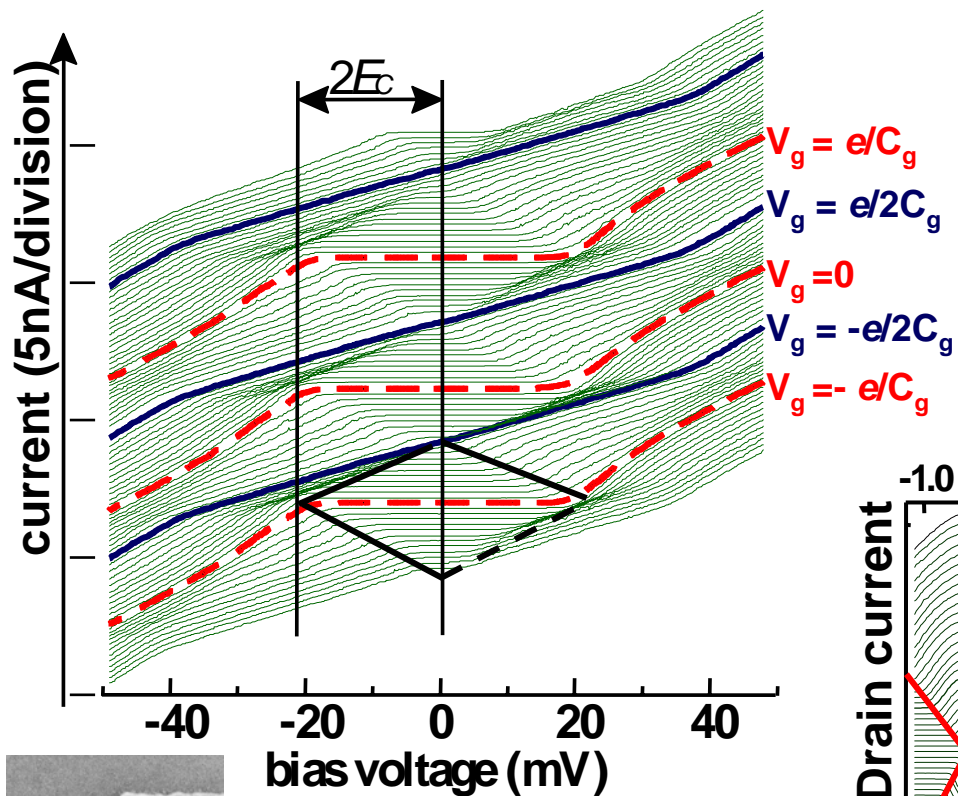
(Lawrence Berkeley National Lab.)

Ronald P. Andres *et al.* Science 272,1323 (1996)

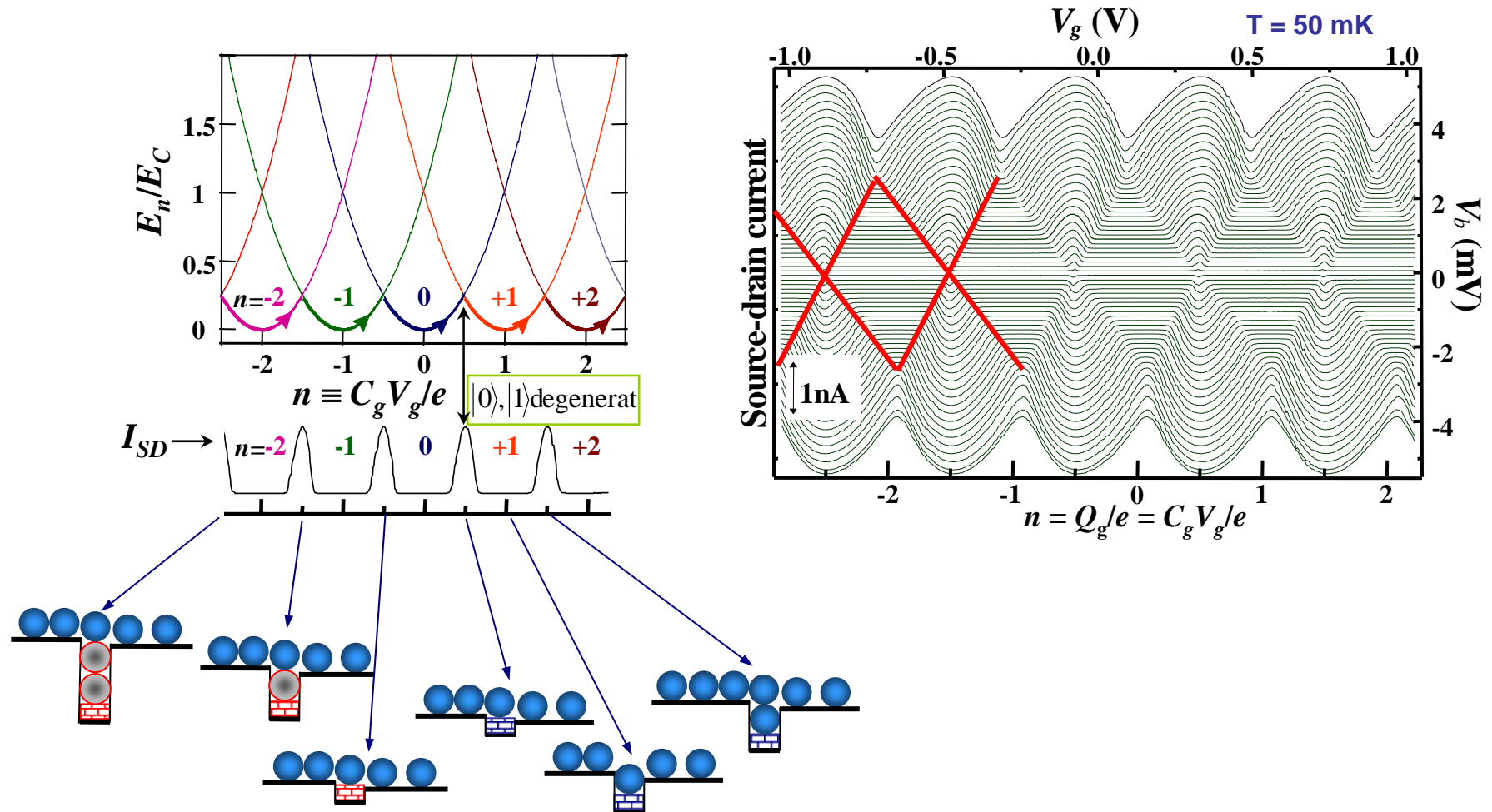
(Purdue Univ.)



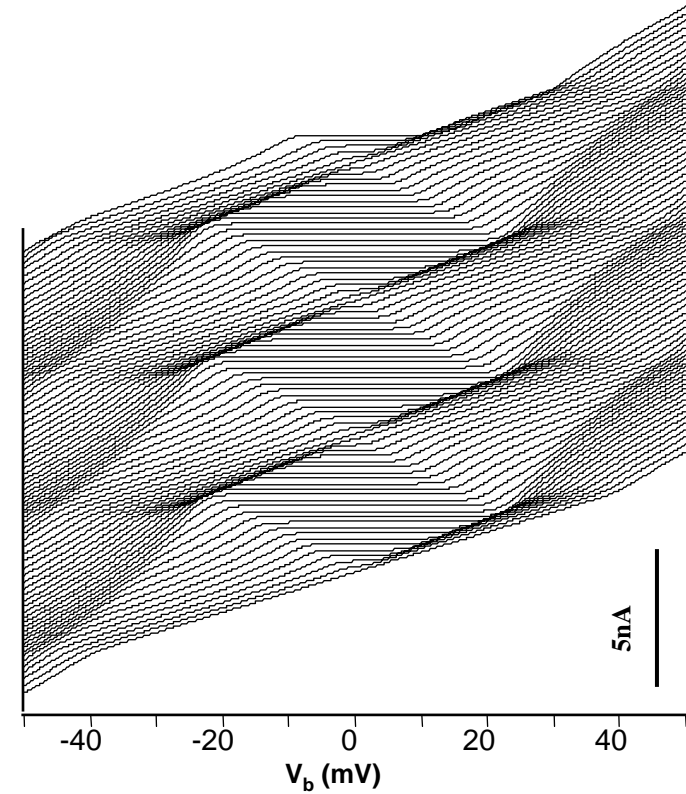
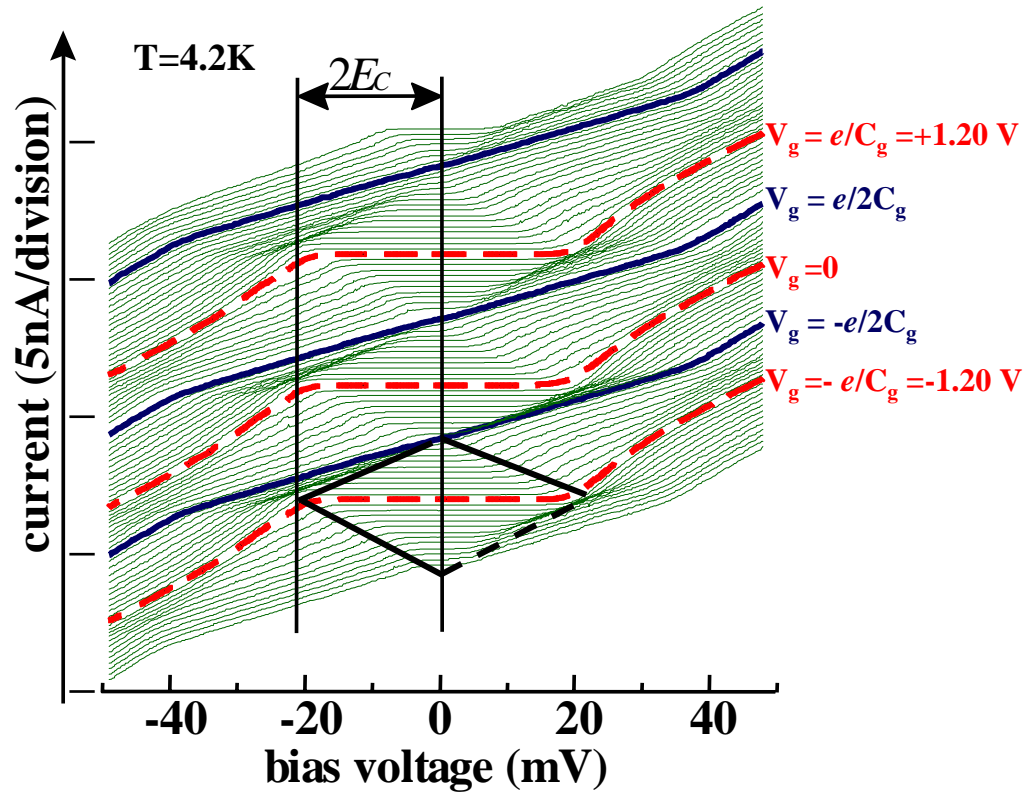
$I(V_b)$ vs. $I(V_g)$



Coulomb Oscillation



Measurement vs. Simulation



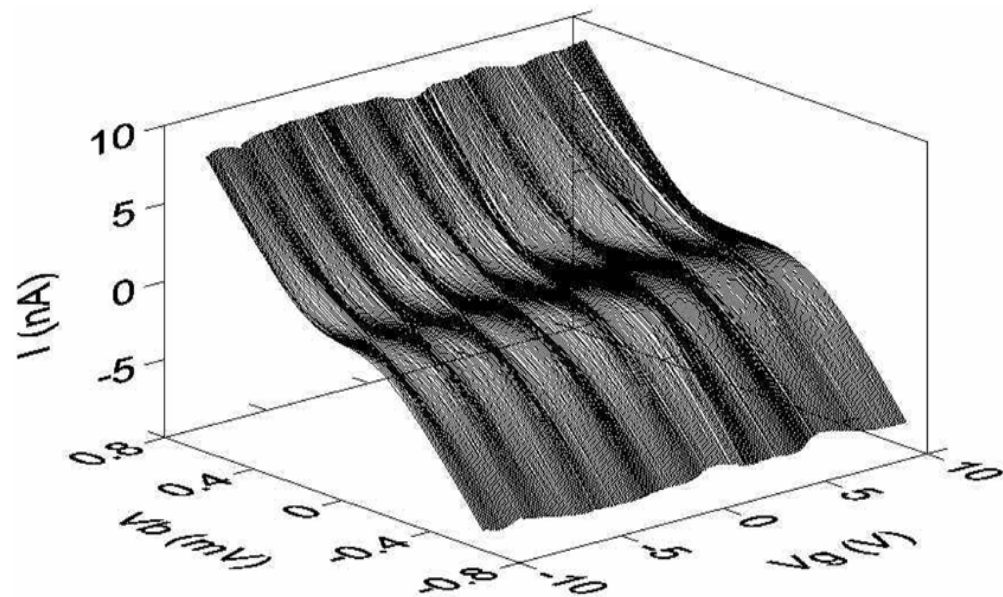
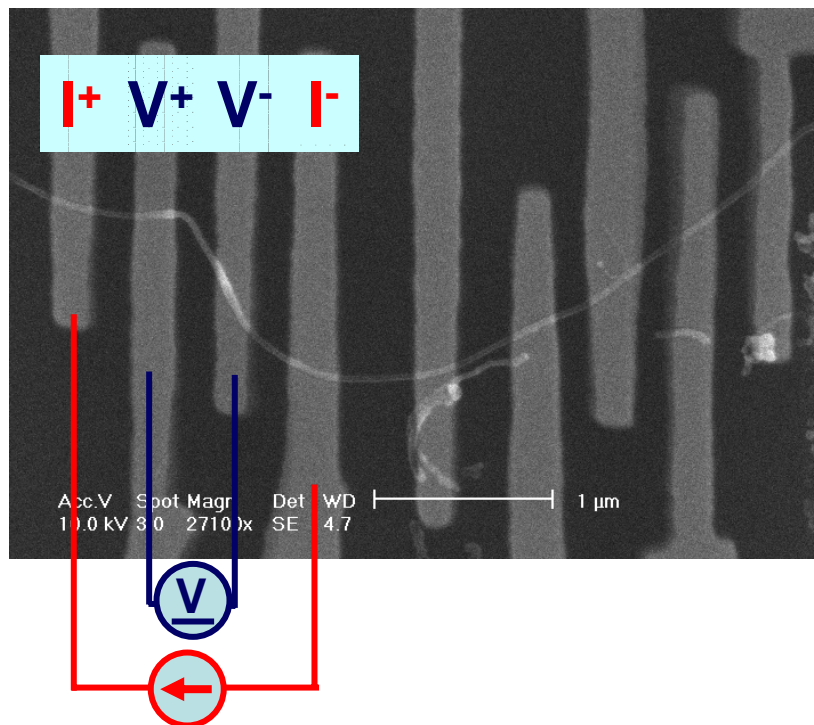
Tunneling rate in orthodox theory

$$\Gamma(\Delta W) = \frac{1}{e^2 R} \frac{\Delta W}{1 - \exp(-\Delta W/k_B T)}$$

Coulomb blockade oscillation in Multiwalled Carbon nanotubes observed by 4-probe techniques

4-probe = measuring the intrinsic properties of wires
contacts are irrelevant

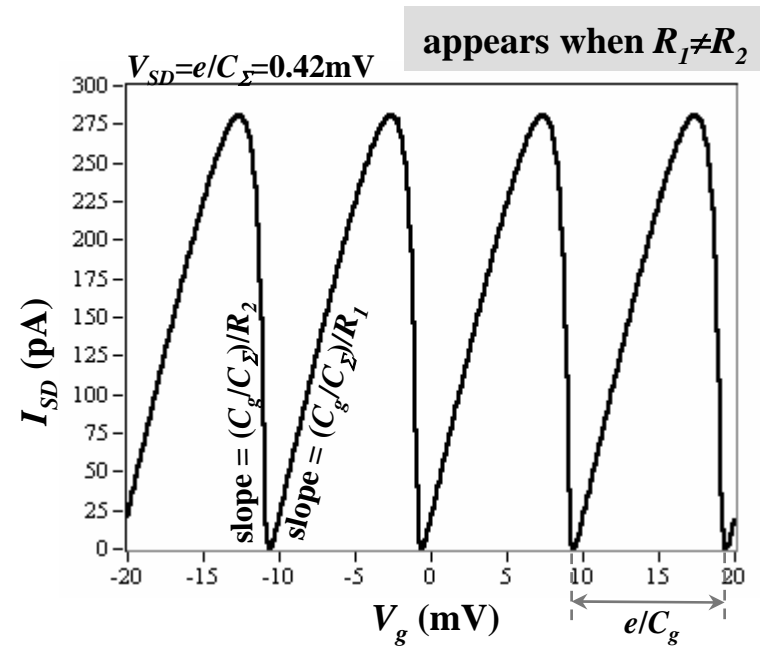
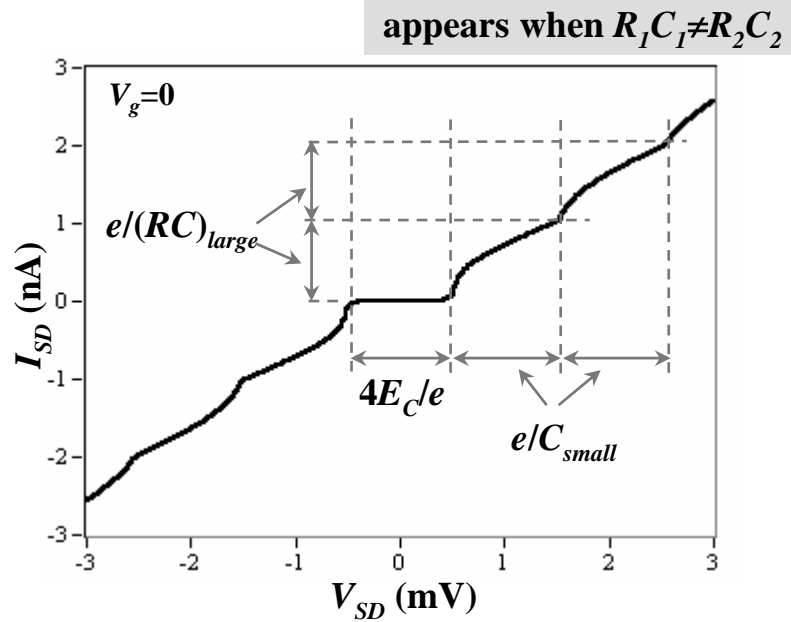
→ tunneling related phenomena such as SET should not appear



Asymmetric SET (simulations)

Coulomb Staircase

$R_1 = 50 \text{ k}\Omega$, $R_2 = 1 \text{ M}\Omega$, $C_1 = 0.2 \text{ fF}$, $C_2 = 0.15 \text{ fF}$, $C_g = 16 \text{ aF}$,
 $E_C = 2.54 \text{ K}$, $T = 20 \text{ mK}$



Charging effect probed by STM

Capacitance of a dielectric disk : $C = 8\epsilon_0\epsilon_r r$

Capacitance of a dielectric sphere : $C = 4\epsilon_0\epsilon_r r$

For a GaAs sphere, $C = 1.47 \times 10^{-18} r$ farad for radius r in nm

VOLUME 63, NUMBER 7 p. 801 PHYSICAL REVIEW LETTERS

14 AUGUST 1989

Scanning-Tunneling-Microscope Observations of Coulomb Blockade and Oxide Polarization in Small Metal Droplets

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⁽²⁾School of Physics and Astronomy, Raymond and Beverly Sackler Faculty of Exact Sciences, Tel Aviv University, 69978 Tel Aviv, Israel

⁽³⁾Scientific Laboratory, Ford Motor Company, Dearborn, Michigan 48121

(Received 20 March 1989)

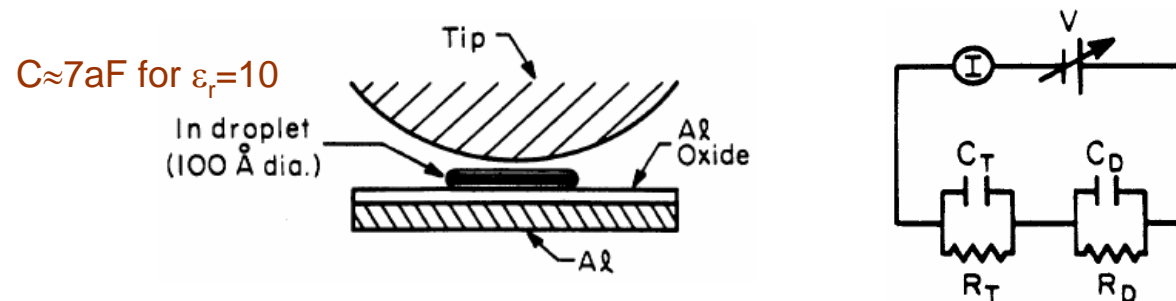


FIG. 1. Schematic showing an In droplet separated from an Al ground plane by a tunneling oxide layer ($\approx 10 \text{ \AA}$ thickness) with an Au STM tip positioned about 10 \AA above it. The equivalent circuit is shown with a voltage source and capacitor C_T for tip to droplet and C_D for droplet to ground plane. The resistors characterize the tunneling probability for each junction and are strictly shot-noise devices.

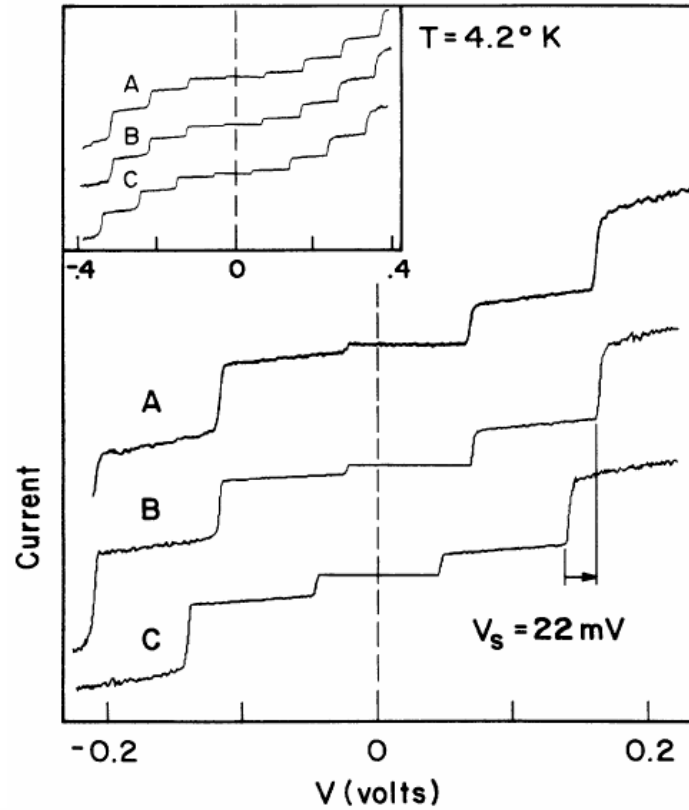
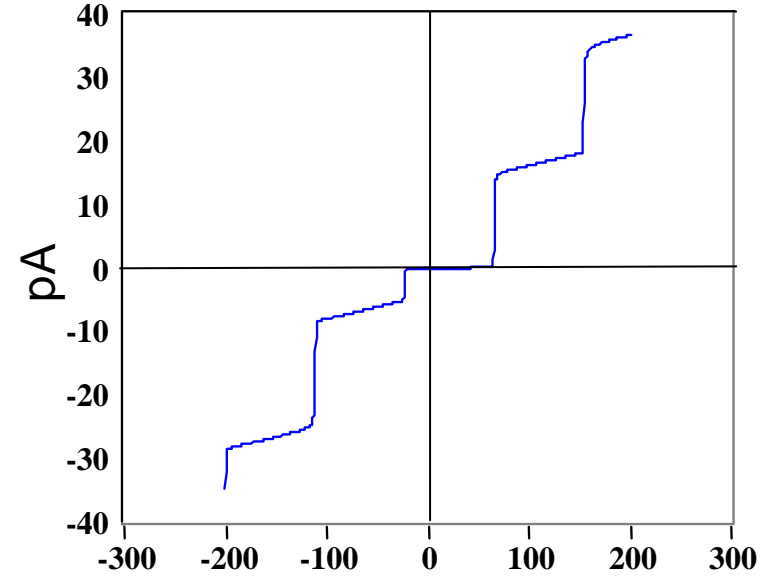


FIG. 2. Curve *A* is an experimental *I-V* characteristic from an In droplet in a sample with average droplet size of 300 Å. The peak-peak current is 1.8 nA. Curve *B* is a theoretical fit to the data for $C_D = 3.5 \times 10^{-19}$ F, $C_T = 1.8 \times 10^{-18}$ F, $R_D = 7.2 \times 10^6$ Ω, and $R_T = 4.4 \times 10^9$ Ω. The obvious asymmetric features in curve *A* require a voltage shift $V_s = 22$ mV ($V_p = 18$ mV). Curve *C*, calculated for $V_s = 0$, shows the (seldom observed) symmetric case. As explained in the text, a small quadratic term was added to the computed tunneling rate for each junction. Inset: A wider voltage scan for this same droplet; again, the topmost curve is experimental data.



$$\begin{array}{l}
 C_t \boxed{350.0} \text{ zF} \quad C_d \boxed{1.8} \text{ aF} \quad C_g \boxed{0.01} \text{ aF} \quad V_t \\
 R_t \boxed{4.40} \text{ GOhm} \quad R_d \boxed{7.20} \text{ MOhm} \quad V_g \boxed{3.87} \text{ V} \quad \frac{V_p}{V_g} = \frac{C_s}{C_i + C_d}
 \end{array}$$

$$V_D = \frac{VC_T}{C_D + C_T} + \frac{Ne}{C_D + C_T} + V_p$$

by scanning the substrate.¹⁴ However, we do not understand at this time why C_T is greater than C_D . At $T = 4$

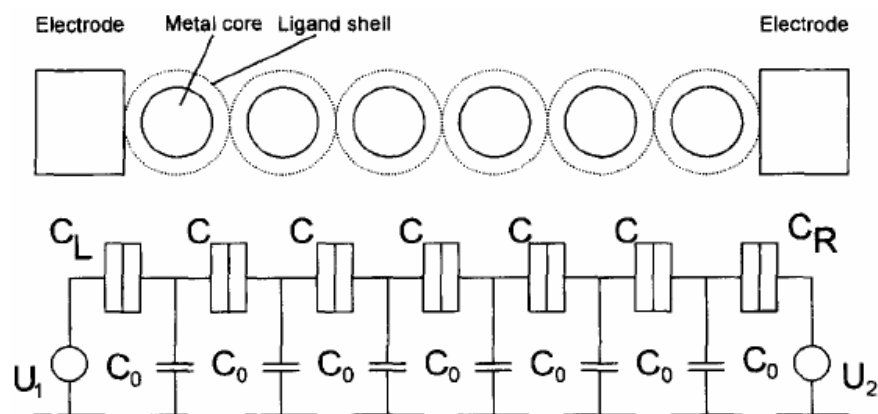
Potential distribution in a finite 1-D array of arbitrary mesoscopic tunnel junctions

V. Gasparian¹, U. Simon*

Institut für Anorganische Chemie, Festkörperchemie, Universität GH Essen, Schützenbahn 70, Essen 45127, Germany

Received 2 September 1997

Physica B 240 (1997) 289–297



$$-\varphi_i = \frac{e}{C_{eff}} \Pi^{|i-k|}$$

$$\Pi \equiv x - \sqrt{x^2 - 1} \quad C_{eff} = \sqrt{C_0^2 + 4CC_0}$$

$$x = 1 + C_0/2C$$

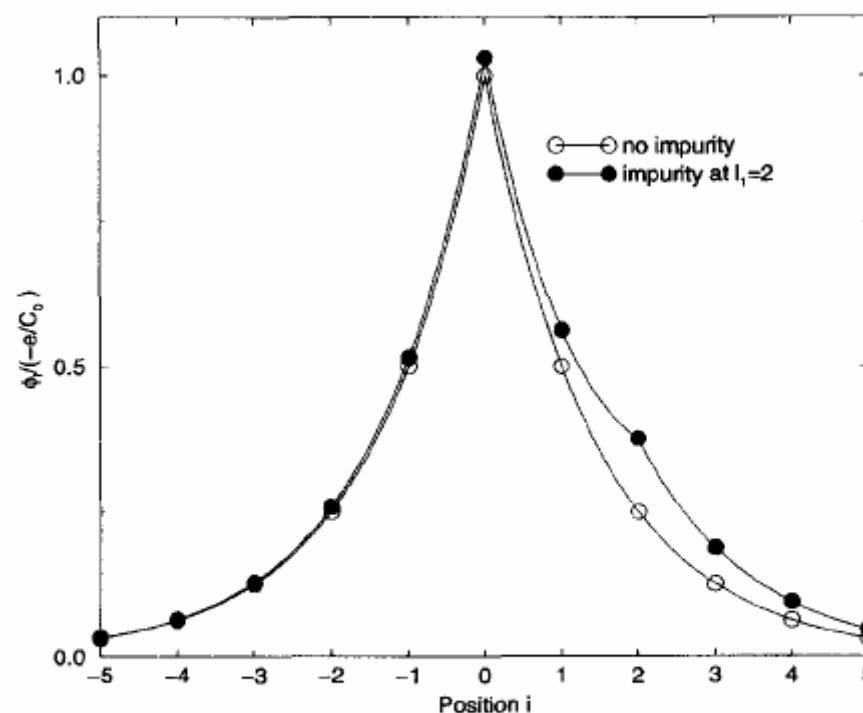
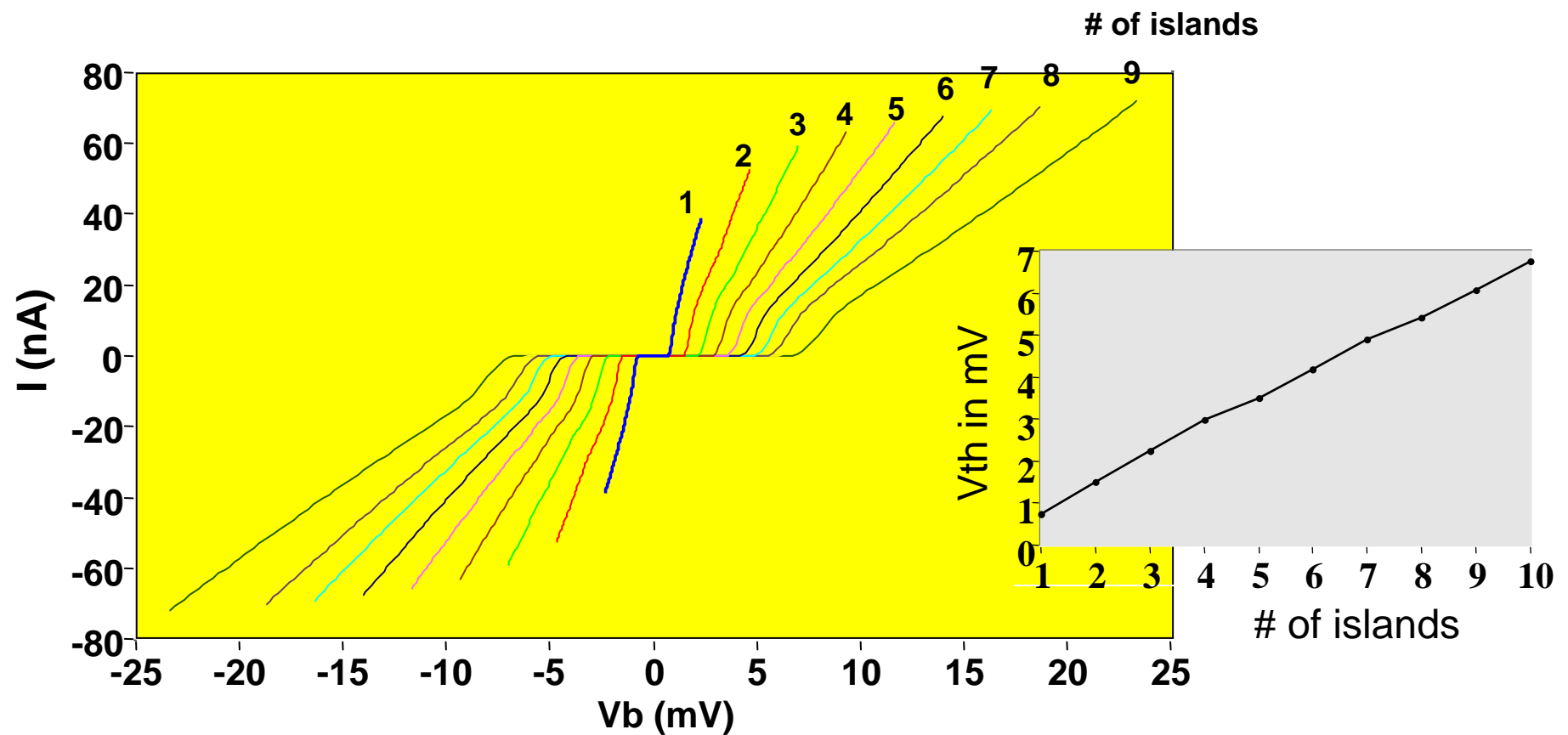


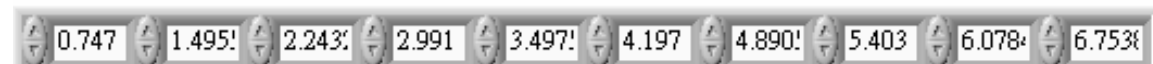
Fig. 2. Potential distribution of the infinite array in the units of $-e/C_0$ and for one impurity at $l_1 = 2$ with $C_1/C_0 = -0.999$.

IV characteristics for 1D array with $C=100\text{aF}$, $R=20\text{k}\Omega$

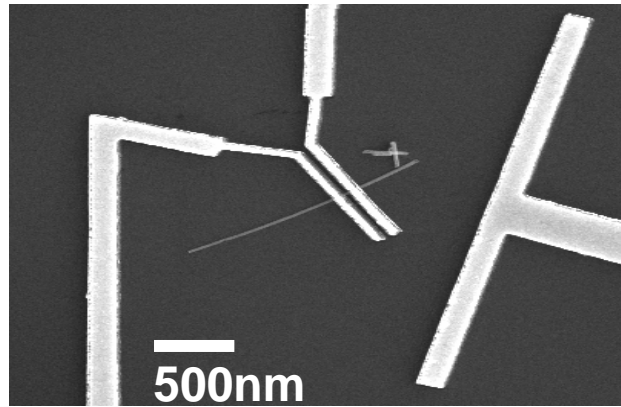


$$e/C=0.8\text{mV}$$

V_{th} in mV

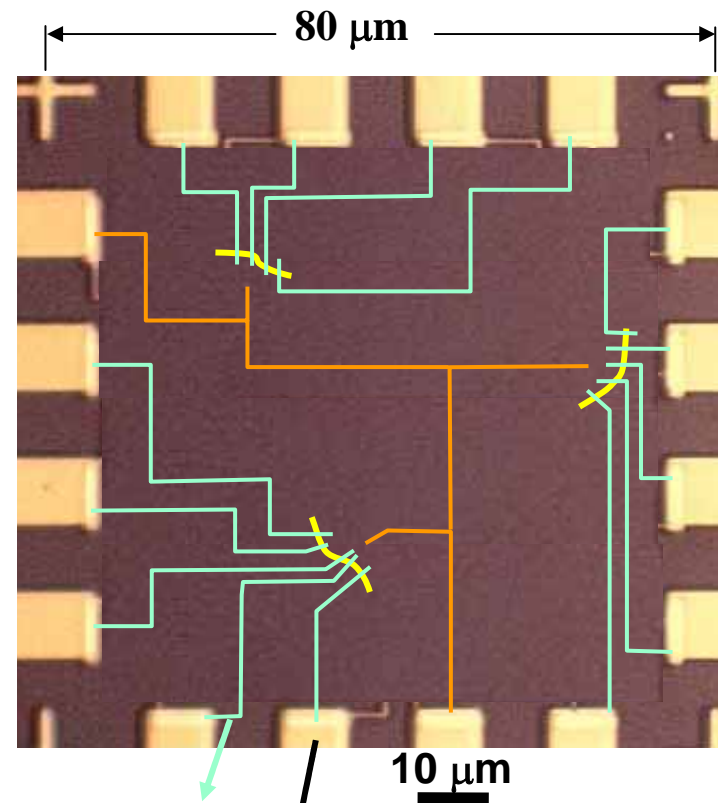
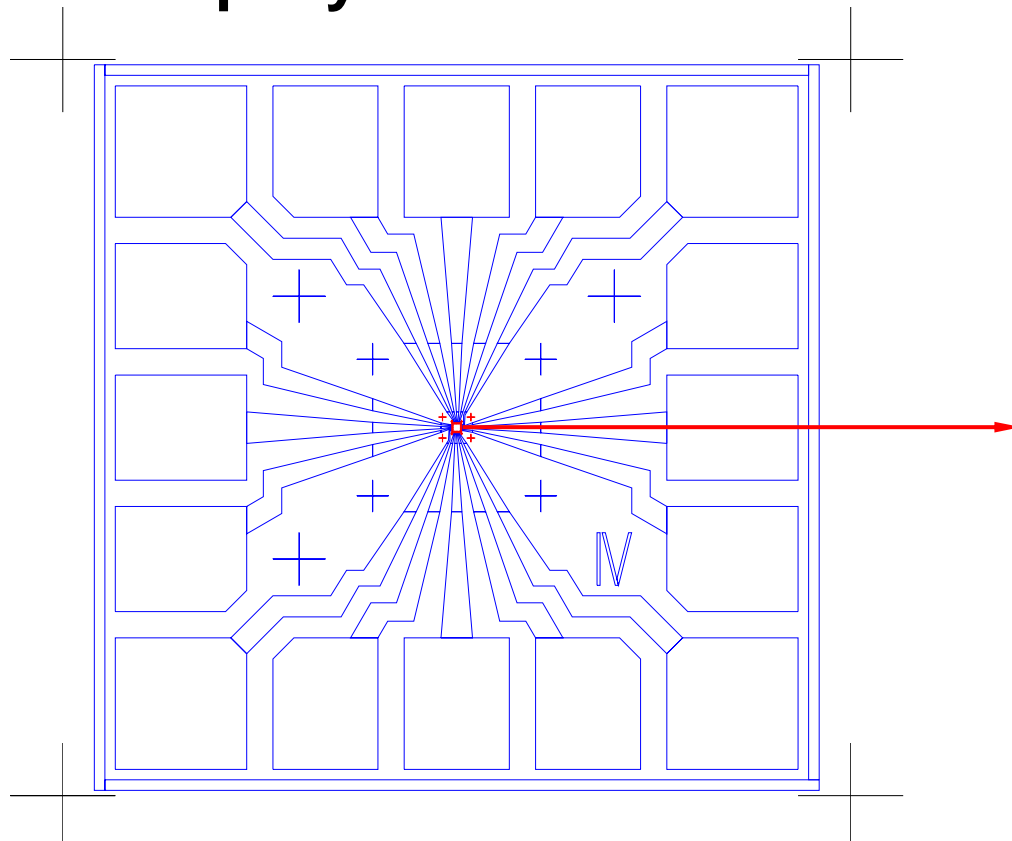


Device fabrication



Wires from
台大凝態中心林麗瓊教授
中研院原分所陳貴賢教授

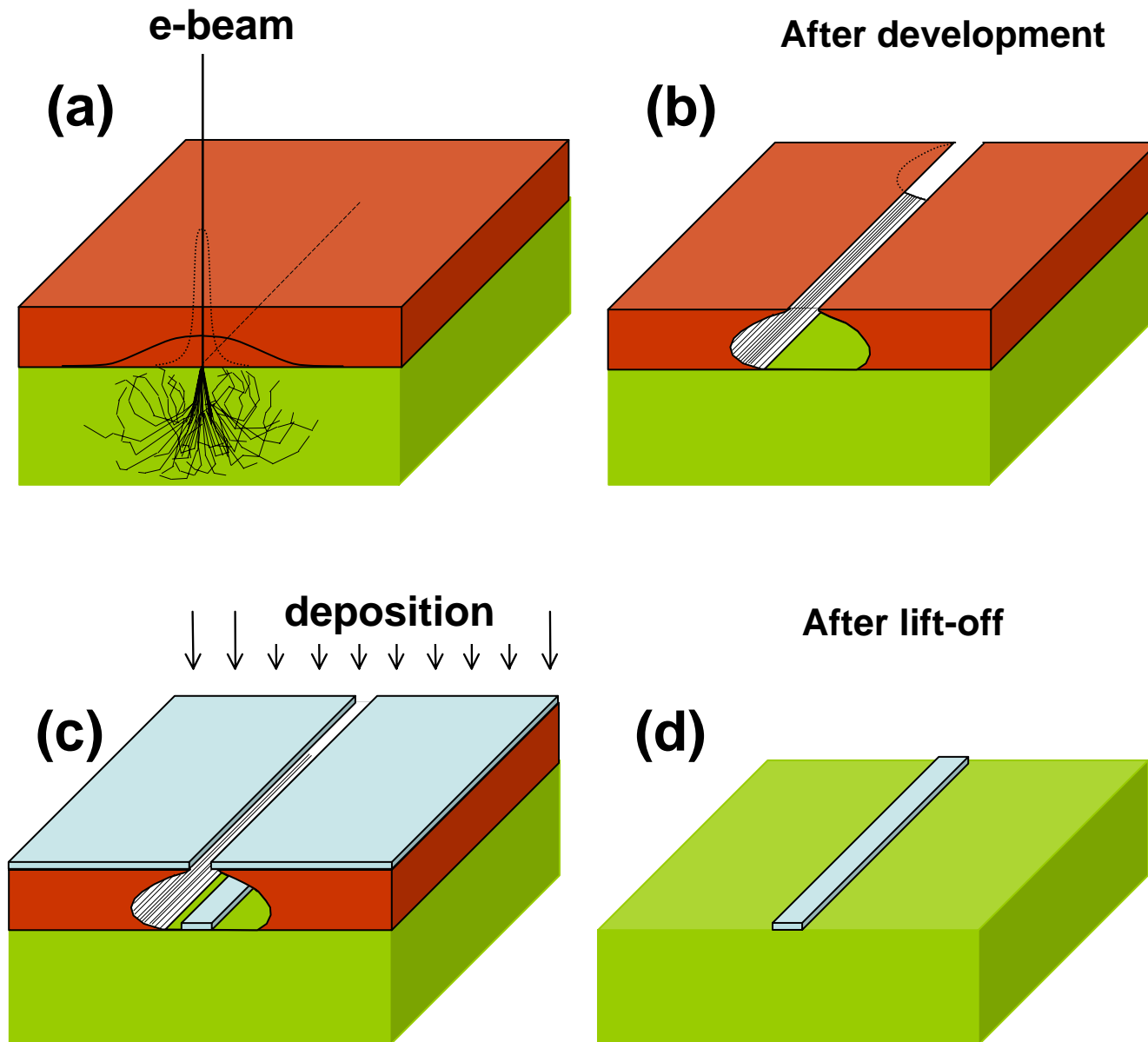
Chip layout



Ni/Au leads

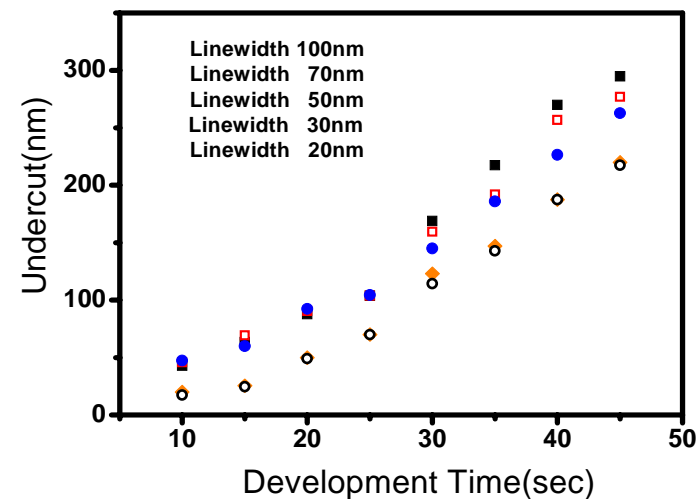
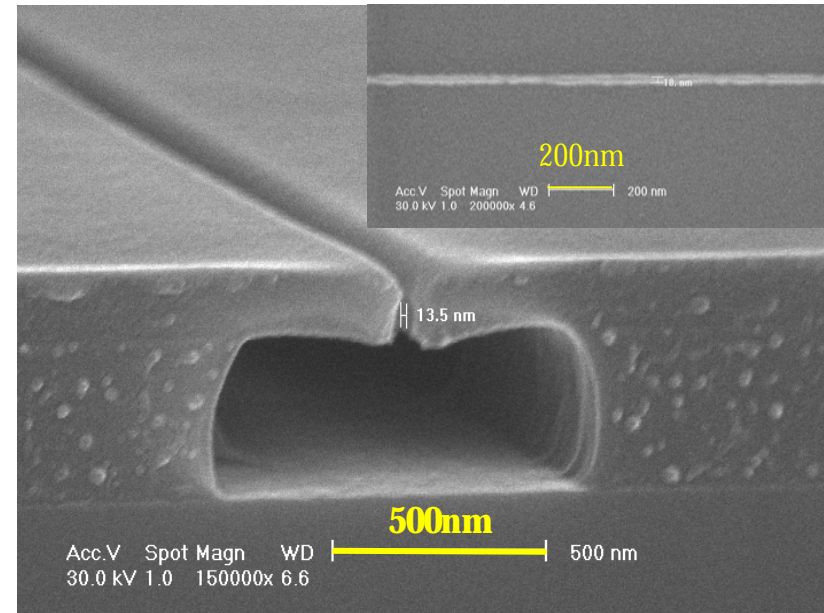
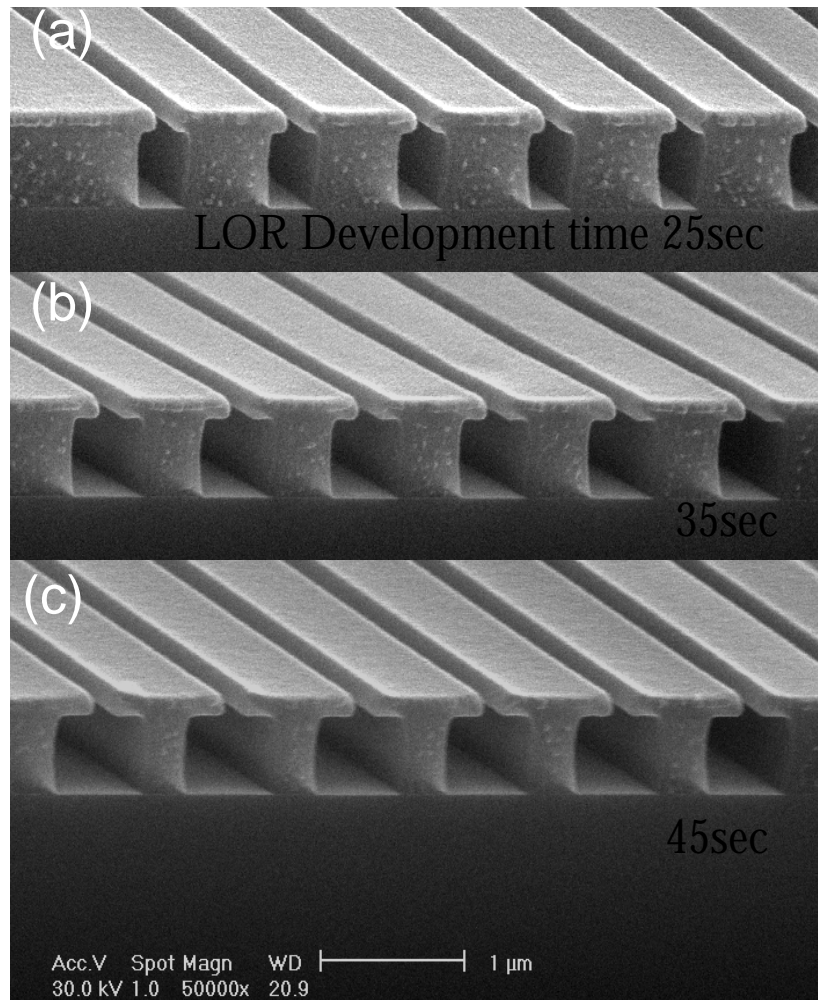
Au leads made by photo-lithography

E-beam lithography for nano-scaled electrodes



Bi-layer polymer system (PMMA/LOR)

We chose to use a PMMA/LOR bi-layer polymer system because LOR is a polydimethylglutarimide (PMGI)-based polymer and has a much higher charge-sensitivity than that of top PMMA layer.



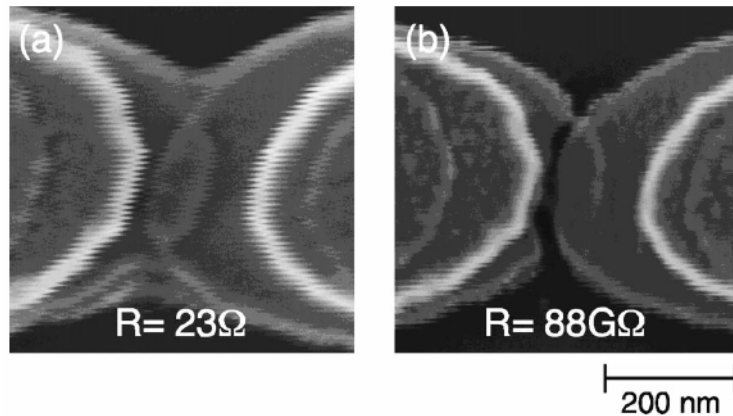
Sample fabrication facilities



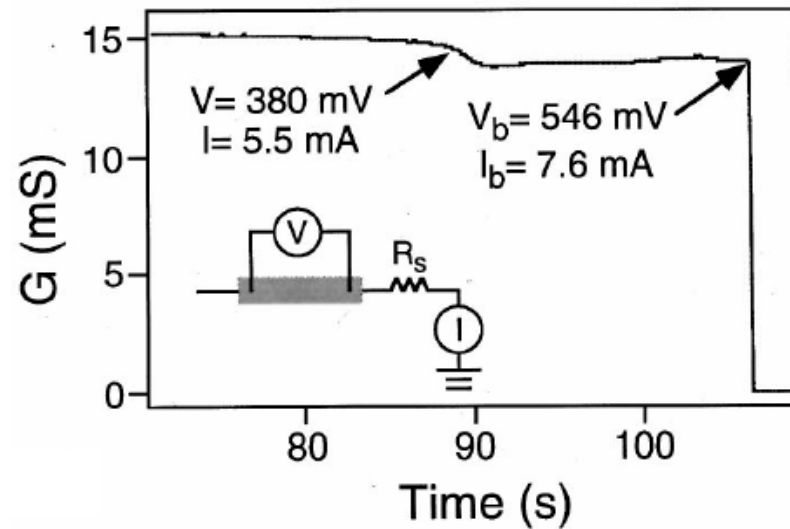
Measurement equipment



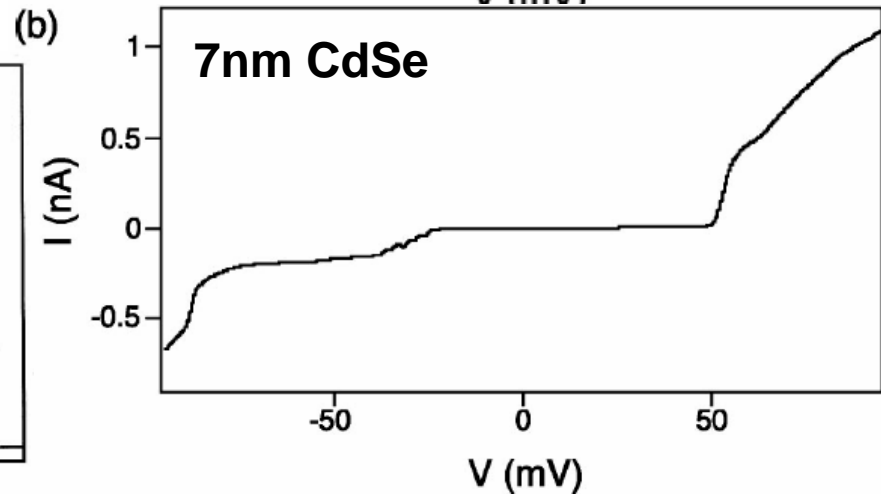
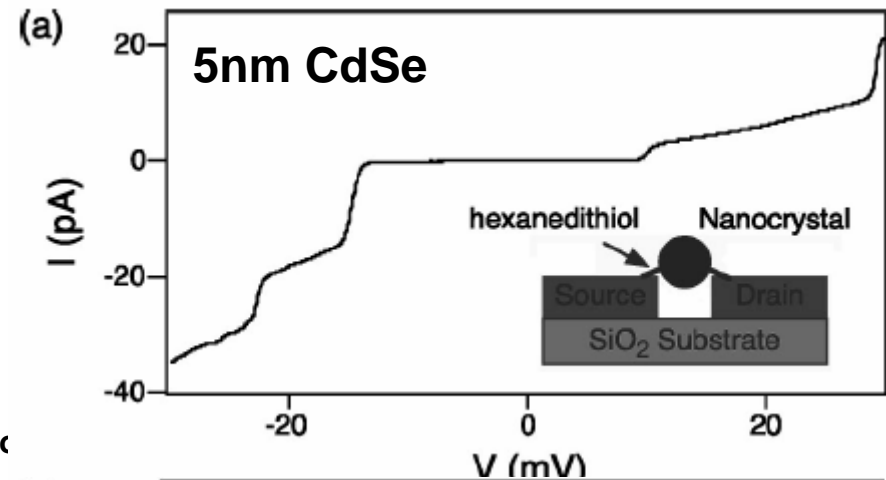
Break junction approach for electrical measurement of a single nano-particle



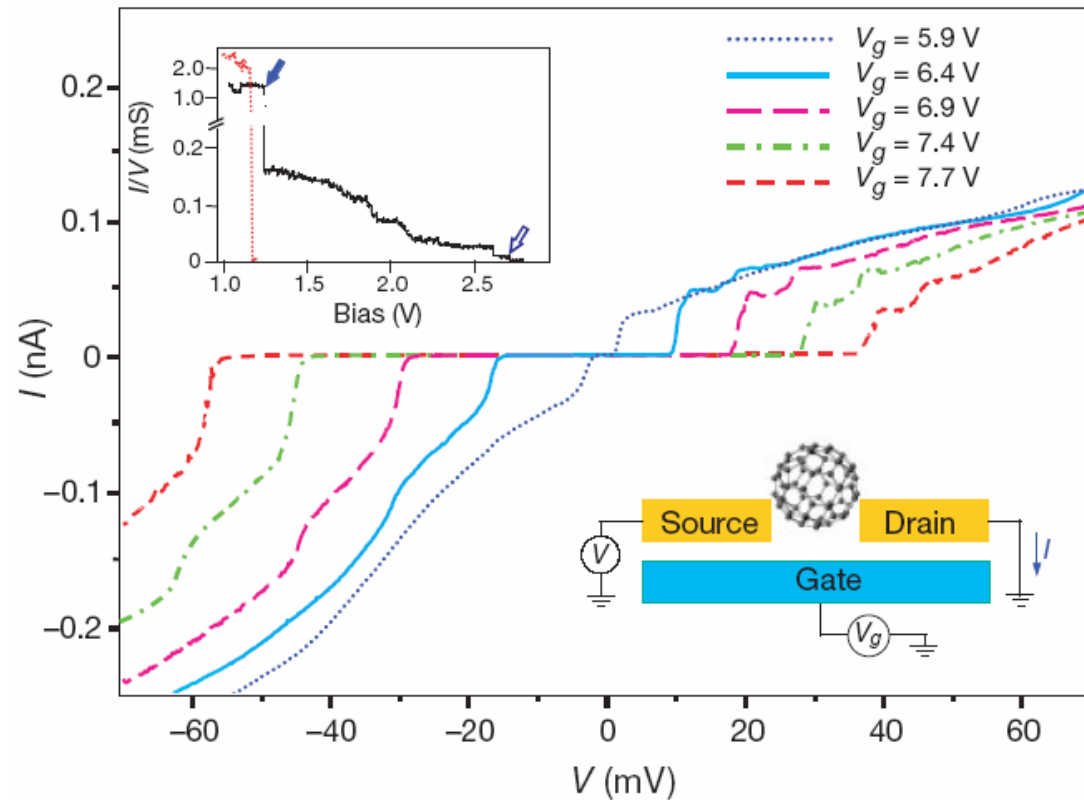
Fabrication of metallic electrodes with nanometer separation by electromigration [APL, 75, 301 (1999)]



IV at 1.5K



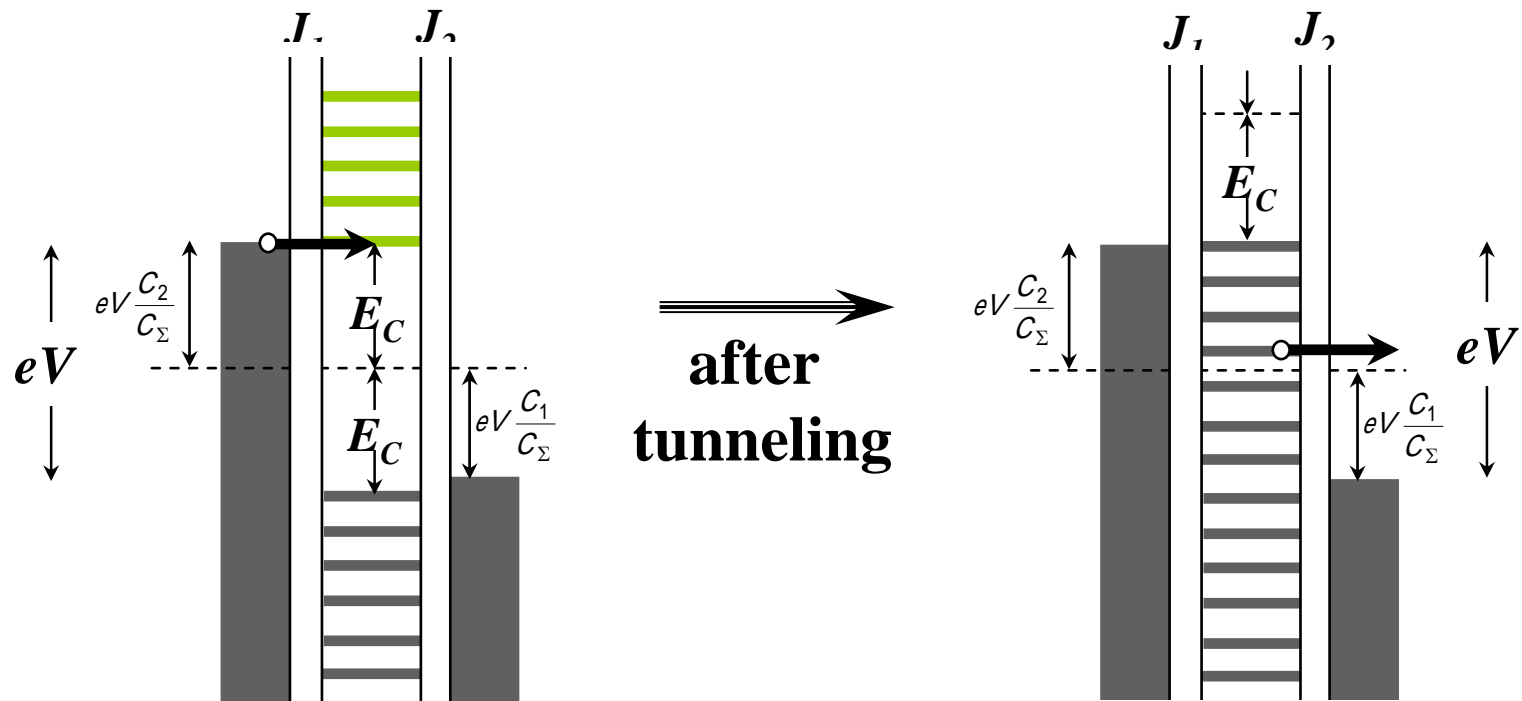
IV characteristics of a single C₆₀ transistor at 1.5K



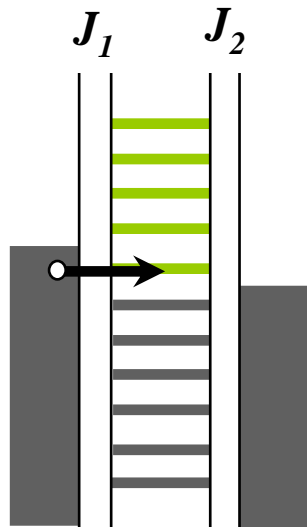
Nature 407, 58 (2000)

Figure 1 Current–voltage (I–V) curves obtained from a single-C₆₀ transistor at T= 1.5 K. Five I–V curves taken at different gate voltages (V_g) are shown. Single-C₆₀ transistors were prepared by first depositing a **dilute toluene** solution of C₆₀ onto a pair of connected gold electrodes. A gap of,1 nm was then created using electromigration-induced breaking of the electrodes. Upper inset, a large bias was applied between the electrodes while the current through the connected electrode was monitored (black solid curve). After the initial rapid decrease (solid arrow), the conductance stayed above ,0.05 mS up to ,2.0 V . This behaviour was observed in most single-C₆₀ transistors, but it was not observed when no C₆₀ solution was deposited (red dotted curve). The bias voltage was increased until the conductance fell low enough to ensure that the current through the junction was in the tunnelling regime (open arrow). The low bias measurements shown in the main panel were taken after the breaking procedure. Lower inset, an idealized diagram of a single C₆₀-transistor formed by this method.

Electron transfer through a quantum dot with charging energy E_C



What makes electrons flow?



biasing: $\mu_1 - \mu_2 = qV_D$

Fermi function:

$$f_1(E) \equiv \frac{1}{1 + \exp[(E - \mu_1)/k_B T]} = f_0(E - \mu_1)$$

$$f_2(E) \equiv \frac{1}{1 + \exp[(E - \mu_2)/k_B T]} = f_0(E - \mu_2)$$

Coupling strength for J_1 and $J_2 = \gamma_1$ and γ_2

$$I_1 = e \left(\frac{\gamma_1}{\hbar} \right) (f_1(\varepsilon_1) - p) \quad I_2 = e \left(\frac{\gamma_2}{\hbar} \right) (f_2(\varepsilon_2) - p)$$

p = average number of electron in the QD $0 \leq p \leq 1$

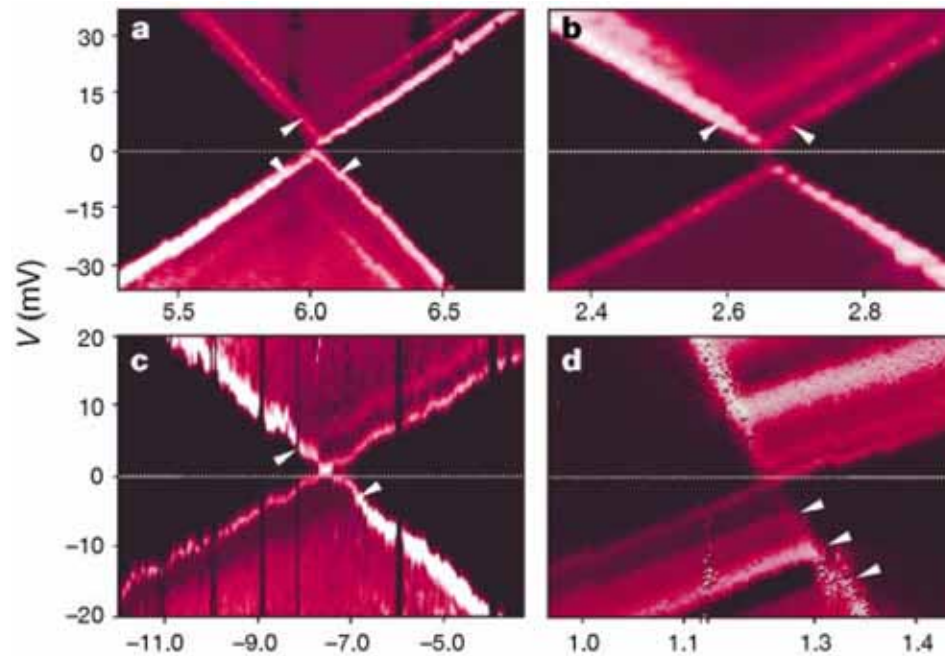
At steady state, $I_1 = I_2$

$$p = \frac{\gamma_1 f_1 + \gamma_2 f_2}{\gamma_1 + \gamma_2}$$

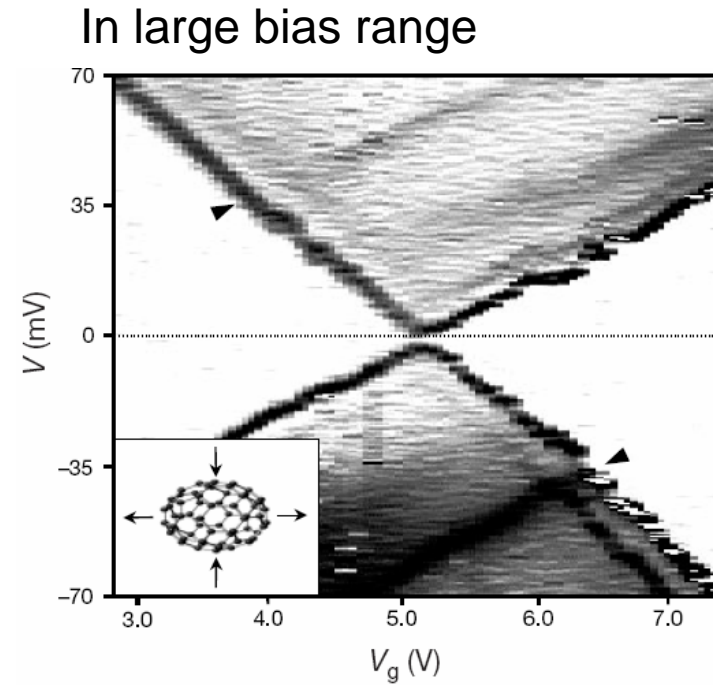
$$\Rightarrow I = I_1 = -I_2 = \frac{e}{\hbar} \frac{\gamma_1 \gamma_2}{\gamma_1 + \gamma_2} [f_1(\varepsilon) - f_2(\varepsilon)]$$

Handbook of Nanoscience, Engineering, and Technology
Section: 12.2.2 Current flow as a balancing act
By Magnus Paulsson, Ferdows Zahid and Supriyo Datta,
Edited by William A. Goddard, III et al. CRC Press, 2003

Quantum transport: atom to transistor
Sec. 1.2 What makes electrons flow



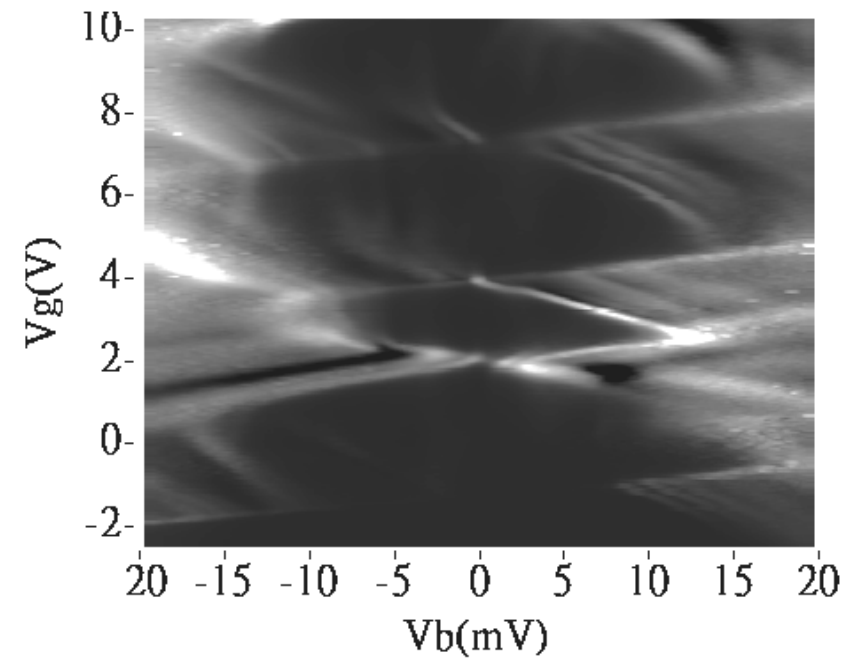
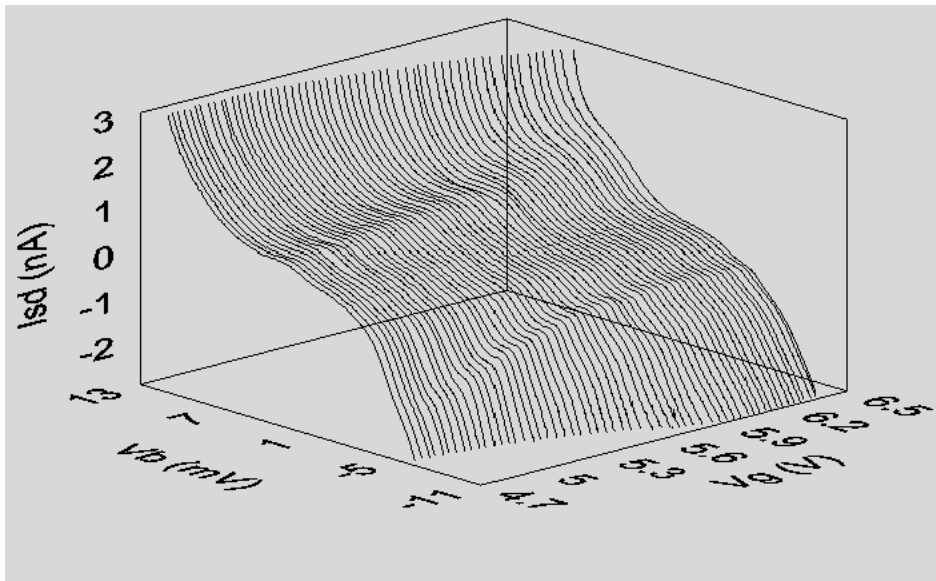
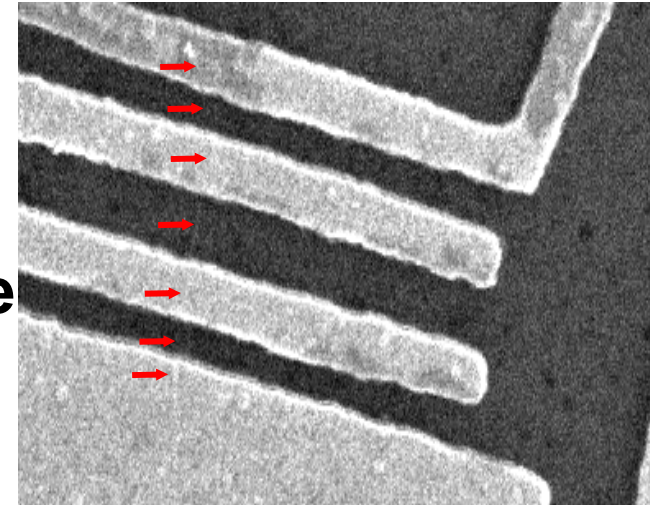
Charging energy $> 270\text{meV}$ V_g (V)



Two-dimensional differential conductance ($\partial I / \partial V$) plots as a function of the bias voltage (V) and the gate voltage (V_g). Data were obtained from four different devices prepared from separate fabrication runs. The dark triangular regions correspond to the conductance gap, and the bright lines represent peaks in the differential conductance. a–d, The differential conductance values are represented by the colour scale, which changes from black (0 nS) through pink to white (white representing 30 nS in a, b and c and 5 nS in d). The white arrows mark the point where $\partial I / \partial V$ lines intercept the conductance gap. During the acquisition of data in d, one ‘switch’ where the entire $\partial I / \partial V$ characteristics shift along the V_g axis occurred at $V_g = 1.15$ V. The right portion of the plot d is shifted along the V_g axis to preserve the continuity of the lines.

Single-walled carbon nanotube with quantum dot behavior

SEM image

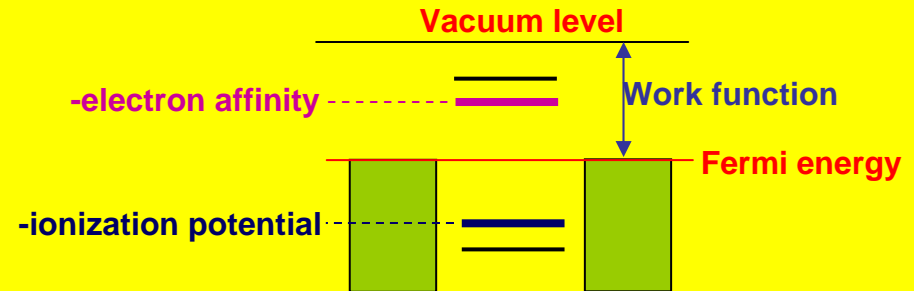
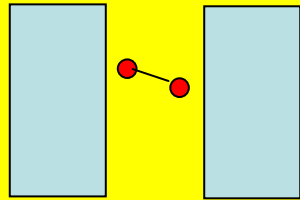


$$\Delta E = 2.5 \sim 3.0 \text{ meV}$$

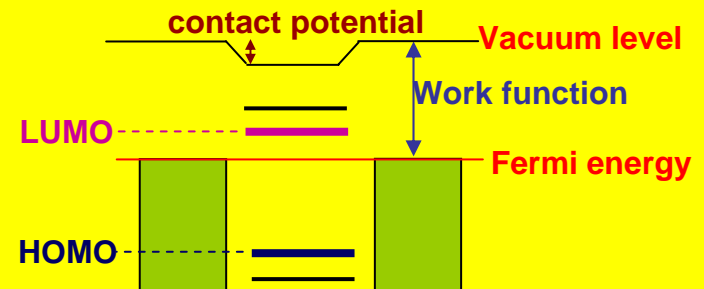
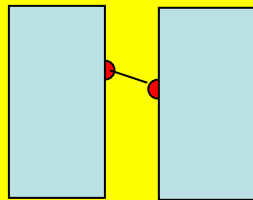
$$E_C = 6 \sim 9 \text{ meV (70} \sim 100 \text{ K)}$$

Where is the Fermi Energy?

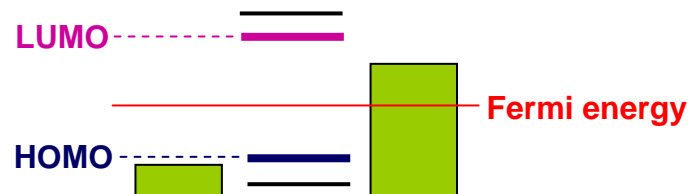
Weakly coupled molecule



Strongly coupled molecule



in positive bias



in negative bias

